



- 1 Small- and meso-scale field-aligned auroral current structures, their spatial and temporal
- 2 characteristics deduced by Swarm constellation
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10 Abstract Magnetic field recordings by the Swarm A and C spacecraft during the Counter Rotation Orbit phase are used for checking the stationarity of auroral region small-and meso-scale field-11 aligned currents (FAC). The varying separation between the spacecraft in along- and cross-track 12 direction during this constellation phase allow for determining the spatial and temporal correlation 13 lengths for FAC structures of different along-track wavelengths. We make use of the cross-14 correlation analysis to check the agreement of the magnetic signatures at the two spacecraft. When 15 the cross-correlation coefficient exceeds 0.75 at a time lag that equals the along-track time 16 17 difference, the event is identified as stationary. It is found that meso-scale FACs of along-track 18 wavelength >100 km are primarily stable for more than 40 s and over cross-track separations of 20 km. An important reason for their deselection is the latitudinal motion of the current system. 19 20 Conversely, stable small-scale FACs (10-75 km wavelength) are found primarily only in a very 21 limited space, up to about 12 km in cross-track and ~18 s in along-track time difference. This class 22 of small-scale FACs is the typical one found commonly in the cusp region and near the midnight 23 sector. Not all the FACs within this limited spatial and temporal regime are stable. In particular 24 for those with high current density occurring during enhanced solar wind input we do not find equivalent signatures at the accompanying satellite. They seem to represent narrow solitary Alfvén 25 26 wave features.

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30 1. Introduction

31 Field-aligned currents (FAC) are commonly observed in magnetized plasmas due to the high electric conductivity along magnetic field lines. In near-Earth space they can be found in the 32 ionosphere and magnetosphere, with particularly strong currents flowing within the auroral 33 34 regions. FACs appear at a wide range of horizontal scales from less than a kilometer (e.g., Neubert and Christiansen, 2003; Rother et al., 2007) up to 1000 km (e.g., Iijima and Potemra, 1976; 35 Anderson et al., 2014). Commonly, the observed current densities become larger the smaller the 36 scales are. Pairs of upward and downward currents are mostly close together. Rother et al. (2007) 37 38 reported strongest current densities for horizontal scales of about 1 km wavelength (including both 39 current directions). Smaller scales are progressively damped at F-region altitudes (Lotko and 40 Zhang, 2018).

41 Large-scale FACs at auroral latitudes can be treated as quasi-static circuits, having their source in the magnetosphere, and the closure current in the ionosphere acts as a dissipative load. For smaller 42 scales of FACs, approaching 10s of km, the dynamic characteristics of the circuit become 43 important, and reflection at the ionosphere play a role (e.g. Lysak, 1990; Vogt and Haerendel, 44





1998). Thus, Alfvén waves, carried by FACs, are becoming the dominate feature in the kilometerscale range. Ishii et al. (1992) made use of magnetic and electric field measurements by the Dynamic Explorer 2 satellite for distinguishing between waves and static current circuits. From the amplitude relation between the two fields, they could discriminate the two types of FAC observations, showing that a cutoff exists for stationary FACs exists at the small-periods end around 4 s to 10 s. When considering the satellite velocity of about 7.5 km/s, the apparent periods convert to latitudinal wavelengths of 30 km to 75 km.

52 When deriving FAC estimates from satellites, one important assumption is the stationary of the 53 magnetic signal over the time of measurement. This is not always satisfied. As reported by 54 Stasiewicz et al. (2000) in their review article, small-scale FAC structures are commonly 55 associated with kinetic Alfvén waves. Therefore, the current strength exhibits a significant 56 temporal variation. From single-satellite magnetic field measurements it is not possible to 57 distinguish between temporal and spatial variations. To overcome this problem Gjerloev et al. 58 (2011) made use of the three ST5 satellites in pearls-on-a-string formation. They performed a large 59 statistical study of the FAC temporal stability depending on their along-track scale size. On the nightside FAC structures larger than 100 km had been found to be quasi-stationary, while on the 60 dayside this was only true for scales above 200 km. Due to their orbital geometry, the ST5 61 spacecraft were cycling the Earth in a sun-synchronous mode, which provided only little local time 62 coverage. Furthermore, the orbits of the three spacecraft were well lined up. Therefore, no 63 64 information on the longitudinal correlation length of FAC structures could be achieved. A similar study of FAC spatial and temporal scales was performed by Lühr et al. (2015), making use of the 65 three Swarm spacecraft soon after launch when the satellites where slowly separating from each 66 other. These authors generally confirmed the results of Gjerloev et al. (2011). FACs of scales up 67 68 to some 10s of kilometer are highly variable with typical persistent periods of about 10 s. Larger-69 scale FACs (>150 km) can be regarded as quasi-stationary, being stable over more than 60 s. Concerning the longitudinal extension of the small FAC sheets, it was found to be 4 times large 70 than the latitudinal width on the nightside, but on the dayside both scales were found to be 71 72 comparable. In spite of these valuable results, that study had a number of limitations. The data 73 were taken during December 2013 - January 2014, over less than 50 days. Thereafter, dedicated 74 orbit maneuvers were performed. This means, it was not possible to investigate any seasonal 75 dependences nor get a good local time coverage.

76 One of the standard Swarm Level-2 data products is the FAC density estimate derived from the magnetic field measurements of the Swarm A and Swarm C spacecraft flying side-by-side. An 77 78 important assumption for the reliability of the product is that both satellites record magnetic field 79 variation caused by the same FAC structure. In a dedicated study Forsyth et al. (2017) compared 80 the recordings at the two Swarm spacecraft by means of a cross-correlation analysis. They 81 generally find large cross-correlation coefficients when the magnetic field data are low-pass filtered with a cutoff of about 20 s. In addition, these authors request a similarity in signal 82 amplitude at the two satellites, not exceeding a difference of 10%. This request significantly 83 reduces the number of suitable data pairs for FAC estimates. We do not consider the amplitude 84 request as justified because the FAC estimate derived from the dual-spacecraft approach is the 85 86 mean value of the estimates at the two sites, and therefore the dominating linear part of the spatial, 87 temporal gradient is taken care of by the averaging.

What is missing after all these studies is a detailed investigation of the smaller-scale FACs. What
are their spatial, temporal correlation lengths? This information is needed, e.g., for determining
the range dual-spacecraft FAC estimates. An opportunity for this kind FAC scale-size analysis
erose during the Swarm Counter-Rotating Orbit Phase in 2021. During that campaign the orbits of
Swarm A and C were brought close together, and Swarm B cycled the Earth in opposite direction.
Early October 2021 all three orbital planes were quasi-coplanar. Thereafter, the orbits slowly
separated again. For the study presented here we make use of the magnetic field recordings from





95 1 May 2021 to 28 February 2022. The dataset covers both June and December solstices and for 96 both seasons all local times are visited. We make use of the cross-correlation analysis, which is 97 applied to the transverse magnetic field component of the Swarm A and C satellites. As a result, 98 we obtain the variation of cross-correlation coefficient and its dependence on the spacecraft 99 separations in along- and cross-track directions. This allows us to address a number of remaining 100 questions concerning the temporal and spatial scales of small- and meso-scale FAC structures.

101 In the section to follow we present the Swarm mission and magnetic field data, in addition, the 102 preprocessing is described. Section 3 introduces the cross-correlation analysis, our prime tool for 103 determining the similarity between the magnetic field signatures at the two spacecraft. This is 104 followed by a statistical analysis, in Section 4, of the derived correlation results during the 10 105 months of our study period. It includes both the characteristics of the identified stable and unstable 106 FAC events separately for 6 spatial scale ranges. The obtained results are discussed in Section 5. 107 Here we try to explain the behavior of the current structures and compare the findings with earlier 108 publications. The Summary and Conclusions give a wrap-up of the derived auroral zone FAC 109 stability characteristics.

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112 2. Data and processing approach

113 ESA's Swarm satellite mission was launched on 22 November 2013. It comprises three identical 114 satellites in near-polar orbits at different altitudes (Friis-Christensen et al., 2008). During the first 115 mission phase, starting on 17 April 2014, the lower pair, Swarm A and C, was flying at an altitude 116 of about 450 km side-by-side with a longitudinal separation of 1.4°, while the third, Swarm B, cruised about 60 km higher. Due to their orbital inclination of 87.3°, Swarm A/C need about 133 117 118 days to cover all 24 local time (LT) hours, when considering both ascending and descending orbital 119 arcs. Swarm B, with a slightly larger inclination (88°) needs about 141 days for a full local time 120 coverage. As a consequence of the difference in inclination, the angle between the orbital planes of Swarm A/C and Swarm B slowly increased. 121

In preparation for a second mission phase, the counter-rotating orbits campaign, the longitudinal separation between Swarm A and C was slowly reduced starting in October 2019. Around 1 October 2021 the orbital planes of Swarm A/C and Swarm B were quasi coplanar. All three crossed the equator at similar longitudes, with Swarm B flying in opposite direction (e.g., Xiong and Lühr, 2023). Thereafter, the separation between Swarm A and C orbital planes increased again at a rate of 0.7° in longitude per year. The months of small Swarm A/C separation around the epoch of coplanarity are of special interest for this study.

129 Here we make use of the Swarm Level-1b 1 Hz magnetic field data with product identifier 130 "MAGx LR", where lower-case "x" in the product names represents a placeholder for the spacecraft, A, B, C. The magnetic vector data are given in the North-East-Center (NEC) frame. 131 132 For our purpose we remove from the Swarm A and C magnetic field observations the contributions 133 of core, crustal, and magnetospheric fields by subtracting the geomagnetic field model CHAOS-134 7.11 (Finlay et al., 2020). The residuals of the horizontal components, Bx and By, are used for 135 studying the magnetic signatures caused by the FACs. From these two components we calculate 136 the deflections of, *B*_{trans}, transverse to the flight direction.

$$B_{trans} = B_{\gamma} \cos(\gamma) - B_{x} \sin(\gamma) \tag{1}$$

138 where $sin(\gamma) = cos(incl)/cos(lat)$ with *incl* as orbital inclination and *lat* as latitude of 139 measurement point. For application in Eq. (1) $\gamma = \gamma$ has to be used on the ascending part of the 140 orbit and $\gamma = \pi - \gamma$ on the descending part.





141 We actually used in our further investigations the difference between two adjacent values of B_{trans} , 142 separated by 1 s. These differences, ΔB_{trans} , help to remove remaining large-scale biases after 143 model subtraction; furthermore, they better represent the characteristics of FACs. A related single-144 satellite FAC estimate would read

$$j_z = \frac{1}{\mu_0} \frac{\Delta B_{trans}}{v_{SC}} \tag{2}$$

146 where μ_0 is the vacuum permeability and v_{SC} is the spacecraft velocity. When inserting, for 147 example, $\Delta B_{trans} = 10$ nT/s and the typical orbital speed of Swarm, 7.55 km/s, we obtain for the 148 FAC density, $j_z = 1.1 \ \mu A/m^2$. This means, the FAC intensities amount approximately to a tenth 149 of the B-field change rate.

150 Several phenomena are associated with FACs of different scale sizes. In order to identify these 151 FAC properties we subdivide the signal of ΔB_{trans} into six quasi-logarithmically spaced period 152 bands. The chosen -3 dB pass-band filter limits are 1-3 s, 3-7 s, 7-13 s, 13-23 s, 23-39 s, and 39-153 60 s.

The studied interval lasts from 1 May 2021 to 28 February 2022. During these months the Swarm A/C cross-track separation is small and stays below 20 km at auroral latitudes. The along-track separation between the spacecraft is deliberately modified around the time of coplanarity between 2 s and 41 s, as shown in Figure 1 of Zhou et al. (2024). The same figure displays also the variation of longitudinal (cross-track) separation at the equator over the study period. Here we are interested in the small and meso-scale FAC characteristics at auroral latitudes beyond $\pm 60^{\circ}$ magnetic latitude (MLat).

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162 **3.** Correlation analysis

A common assumption, when estimating FAC density from magnetic field satellite observations,
is that the signal is quasi-stationary. This condition can be tested, when Swarm A and C observe
the same current structure. We check the stationarity with the help of a cross-correlation analysis.

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$$Cc = \frac{\sum[(X-X_m)\cdot(Y-Y_m)]}{\sqrt{\sum(X-X_m)^2\cdot\sum(Y-Y_m)^2}}$$
(3)

167 where, X represents the signal amplitude of ΔB_{trans} from Swarm A, Y represents the signal of 168 ΔB_{trans} from Swarm C, and Xm and Ym represent the mean values of ΔB_{trans} over the correlation intervals of the two satellites, respectively. We further determine the maxima of Cc and the 169 170 corresponding time lags (T-lag) between the two satellite data series. It has to be noted that in this 171 way T-lag is derived only with a one-second resolution, as the Swarm magnetic measurements 172 have a 1-s resolution. For improving this situation, we consider the five cross-correlation results 173 centered on the maximum Cc value and apply a fourth-order polynomial fitting. From the peak 174 location of the fitted polynomial, we get the T-lag in fractions of a second. According to the 175 definition of Cc in Eq. (3), a positive T-lag means, Swarm A samples the magnetic signal before 176 Swarm C.

Figure 1 presents examples of passes over the north pole. The upper frame is from 1 September 2021, along a midnight - noon orbit. For these example-passes we consider an interval of 60 s (corresponding to a distance of 450 km) of ΔB_{trans} from Swarm A and C to calculate the root mean square (RMS) amplitude and to perform the cross-correlation analysis. The ΔB_{trans} in the top panel shows two groups of signal bursts at auroral latitudes in the recordings of the two satellites. These represent intense FAC at a wide range of scales, the earlier one occurring around midnight and the other at noon. For getting an impression of the related FAC densities, a ΔB_{trans} amplitude of 100





184 nT/s corresponds, as shown above, to approximately $10 \ \mu A/m^2$. The RMS values are given in the 185 second panel. In the third panel optimal T-lag values are plotted. Over large parts of the pass it 186 stays close to the actual time shift, $\Delta t = 4$ s, between the spacecraft, as expected for a static 187 structure. The Cc values at the bottom herald good correlation (Cc > 0.75) for most of the signal. 188 However, that is not always the case for the bursts. In this example, the equatorward part of the 189 nightside burst exhibits an almost perfect correlation at expected lag time, while the correlation on 190 the poleward side is poor. At noon the cross-correlation coefficient is also below the threshold on 191 the poleward parts but is high during the stronger later burst.

192 In the lower frame of Figure 1 an example from 1 January 2022 is displayed. The recordings, in 193 the same format as above, are again from a noon - midnight pass. Bursts of small-scale FACs appear around 78° MLat. An obvious difference to the example shown above is the generally low 194 195 cross-correlation coefficient. Here again, the more equatorward parts of the bursts tend to show larger Cc. The major difference of this example compared to the former is the much longer along-196 track separation between the two satellites of $\Delta t = 22.3$ s. In spite of that enlarged distance, best 197 198 correlation between the two recordings is achieved at time lags close to the actual time difference, 199 but the quality of correlation has decreased considerably.

200 These examples show that small-scale FACs do reach large amplitudes at auroral latitudes, but 201 they cannot generally be considered as static structures.

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203 4. Statistical analysis

The aim of our study is to identify the typical properties of small- and meso-scale FAC structures at auroral latitudes. Such information can best be achieved when simultaneous observations from multiple points are available. For the statistical analysis we considered the time 1 May 2021 - 28 February 2022, when the orbital planes of Swarm A and C differed only by small angle (<0.3°). This promises meaningful multipoint measurements down to the smallest resolvable scales of <10 km.

The study interval overlaps with the trailing part of the solar minimum. Therefore, the mean solar flux level increases slowly from F10.7 = 80 sfu in the beginning to about F10.7 = 110 sfu around the end. The magnetic activity stays generally below Kp = 4. There were just some stormy days, quite outstanding is 4 November 2021, followed at decreasing activity levels by 12 October 2021 and 4 February 2022. Due to the generally calm conditions, the derived mean FAC characteristics represent normal solar wind driving states.

216 It is known from earlier publications that FACs of different scales exhibit different dynamic 217 properties, e.g. Stasiewicz, K., et al. (2000) and references therein. For that reason, we divided the 218 magnetic variations ΔB_{trans} into the six period bands as outlined in Section 2. The cross-correlation 219 analysis, as defined in Eq. (3), was applied separately to all the period bands and all the days. A 220 current structure is considered as stationary when the magnetic field measurements from the two 221 Swarm spacecraft achieve a peak cross-correlation coefficient, Cc > 0.75, at a time shift that fits 222 the along-track separation, $\Delta t \pm 1.5$ s. In order to avoid fault results from too small signals, in 223 addition a minimum RMS amplitude is requested. For the decision on a suitable threshold, we 224 considered the following facts. In their study on low-latitude FAC properties Zhou et al. (2024) had chosen a rather small amplitude limit, RMS > 0.03 nT/s. As can be seen from their Figure 2, 225 FAC signatures frequently surpass 1 nT/s. Those FACs are generally driven by atmospheric 226 processes, as winds and waves. Here we are interested in FACs driven by magnetospheric 227 228 processes, therefore those from atmospheric dynamics, also active at auroral latitudes, should be suppressed. For that reason, we decided to choose a threshold, RMS > 2 nT/s, in this study. For 229 230 assessing the consequence of this selection, Figure 2 shows the occurrence distribution of wavy magnetic signals at auroral latitudes beyond 60° MLat and within the period range of 1-60 s 231





- (corresponding to wavelengths up to 450 km). As expected, the occurrence rate declines, the larger
 the threshold. For RMS > 2 nT/s about 60% of the satellite recordings provide significant FAC
 signal. This percentage fits also approximately the typical latitude range of lijima and Potemra
 (1976) FAC coverage of the auroral oval.
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237 4.1 Variation of identified static structures over the study period

Here we start with a look at the wavy magnetic field signal that could be caused by FACs and how 238 239 this activity varies over the cause of our study period. Figure 3 presents the occurrence frequency, 240 on a daily base, of wavy signals with amplitudes RMS > 2 nT/s, as recorded by Swarm A, 241 separately for the six period ranges and the two hemispheres. No comparison with Swarm C 242 measurements has been performed. Across the bottom of the figure additional information is 243 added, the typical orbital local time (taken at 70° MLat) separately for up- and downleg arcs, the 244 along-track time difference between Swarm A and C, and the east-west separation of the spacecraft 245 at the equator in degree.

246 The curves for the different periods are quite similar. They reach up to rates of 50%. This is 247 expected as a result of our chosen RMS threshold. General features to be noticed; higher rates are 248 observed at local summer conditions. This is the case during the first half of the study period in 249 the northern hemisphere and the second half in the southern. More events are found during daytime 250 than at night. This holds for both hemisphere but is more pronounced for shorter periods (smaller 251 structures) and for local summer time. The occurrence rate varies partly at regular period. This is 252 close to 13.5 days, reflecting halve the rotation period of the sun. There appears a particularly deep 253 dip in occurrence rate at 27-28 October 2021. This is a time of very low magnetic activity, having 254 obviously direct influence on the FAC activity at auroral latitudes. Only one week later, 3-4 255 November, a major magnetic storm occurred, giving rise to enhanced signal rates, in particular, in 256 the southern hemisphere.

257 For comparison, Figure 4 shows for the two hemispheres the temporal variation of the mean 258 occurrence rates of positively detected static structures by means of cross-correlation between 259 Swarm A and C recordings, separately for the up- and downleg orbital arcs within auroral latitudes 260 (beyond 60° MLat), in the same format as before. Particularly low event occurrences are found at 261 the smallest scales (1-3 s). Only during the 2 weeks around coplanarity (1 October) somewhat 262 enhanced rates appear. Within that fortnight the along-track separation was reduced to $\Delta t = 2$ s. 263 This indicates rather small correlation length of these narrow FACs. For larger current structures 264 the curves start to agree more and more with the before shown activity curves, at least for the first half of the study period. During the second half we find a prominent dip in event occurrences 265 266 centered on December 2021. In that month the along-track separation is particularly large between 267 Swarm A and C.

268 Finally, Figure 5 presents the ratio of magnetic signatures that can be interpreted as stable FAC 269 structures. Plotted are the daily detected static structures divided by the counts of activity intervals. 270 These normalized curves reveal which fraction of the high latitudes is covered by stable FACs of 271 different scale-lengths. Rather similar curves are derived from the two hemispheres. Towards 272 longer periods (larger scales) fairly constant ratios between 80% and 100% are derived. This 273 confirms that almost all FACs of these scales at auroral latitudes can be considered as quasi 274 stationary. The rates start to drop for shorter periods shorter than ~ 10 s (<75 km along-track). 275 Particularly low percentages are returned for the very short period (1-3 s). Also here, the rates are 276 somewhat outstanding in the 2 weeks around orbit coplanarity, when the spacecraft separation was 277 reduced to 2 s. Peak occurrence ratios can also be found for other period bands during those two 278 weeks.





279 When normalizing the identified FAC structures by the number of wavy signal events above a certain amplitude threshold (RMS > 2 nT/s), the influences of local time and season are largely 280 281 removed, but the effects of the Swarm constellation prevail. Therefore, these plots are more 282 suitable for evaluating the properties of the detected currents. Larger ratios of stable FACs are 283 obtained in the afternoon sector, compared to the morning side, in particular in the northern 284 hemisphere summer. For large structures, >150 km (>20s), we find ratios of almost 100% everywhere. This means, practically all magnetic variations can be regarded as caused by quasi-285 286 stationary FAC structures. Exceptions appear during November, December 2021. According to 287 Figure 1 of Zhou et al. (2024), during those months the along-track separation between the satellites was above 20 s, increased up to 41 s on 16 December, after which it linearly decreased. 288

289 These ratio plots provide the opportunity to distinguish between the temporal and spatial 290 correlation-lengths of the different FAC scales. From Figure 1 of Zhou et al. (2024) it can be seen 291 that the along-track separation does not vary too much about its mean of ~ 5 s in the beginning of 292 our study period up to mid-September. However, the longitudinal separation becomes 293 progressively shorter. During these early months the occurrence ratios, for example of the 3-7 s 294 and 7-13 s period structures exhibit at both hemispheres linear increases from the beginning of the 295 study interval up to the date of coplanarity (1 October). We may attribute this behavior to the 296 change of cross-track separation, and as expected, its influence is stronger the shorter the signal period. Quite differently, during the time after closest approach, here we find first a rapid decrease 297 298 and after 16 December a recovery of the occurrence ratios. This variation of the ratio we like to 299 attribute to the combination of growing along- and cross-track separations. It is worth noting that 300 the ratios for all periods show approximately the same values at the beginning and end of the study time, when along- and cross-track separations have attained again similar values. For obtaining an 301 302 estimate of the along-track separation on the ratio, one should first remove the cross-track effect, 303 as observed during the time before coplanarity. Without going into details, it is obvious to see that also the along-track separation has a stronger effect on the shorter periods. The number of 304 305 commonly observed stable structures, for example of 3-7 s period signals, is reduced by this effect to more than a factor of 2 at $\Delta t = 40$ s. Smaller reductions are derived for longer periods (larger 306 scales). These observations provide a first idea of the FAC structure's temporal correlation lengths. 307

Somewhat special are the occurrence ratios obtained for the 1-3 s periods. Outside the true coplanarity period with $\Delta t = 2$ s the ratios are quite low, ranging between 10% and 20%. Interestingly, they do not approach zero. This indicates that there is another class of small (some 10 km) FACs that is stable over the range of along- and cross-track separations considered here. We will revisit this issue at the end of Section 5.2.

Another feature, worth to be noted, is the local minimum in FAC ratios on 4 November 2021. This coincides with the strongest magnetic storm during our study period. Suggested reasons for the reduced number of stable FACs are the more dynamic character of the currents with smaller scales and/or the expansion of the auroral oval beyond our limit of 60° MLat for the larger scales. No such effect on the ratios has been found during the other, weaker storm times. Interestingly, while the storm-related reduction of stable FACs is more prominent at shorter scales, there is a deeper dip in ratio at longest periods coinciding with the very quiet days on 27-29 October 2021.

320 An even more detailed view on the occurrence of selected FAC events is provided by Figure 6. It 321 shows the latitudinal distribution of the ratio of stable FAC structures. The prime features of Figure 322 5 are also visible here, but this presentation offers another dimension. When starting with the first 323 half of the study period, before the time of orbital coplanarity, we find for periods larger than ~ 20 324 s at all the latitudes, for all local times and both hemispheres high occurrence ratios. For shorter 325 period, say <15 s, higher ratios tend to concentrate in the northern hemisphere at very polar 326 latitudes. This could be due to some current properties, but it has to be taken into account that the 327 orbital cross-over takes place close to the pole and with that, small cross-track separations are 328 found there. However, when turning to the southern hemisphere (lower frame of Figure 6) we





329 conversely find, e.g. for the 7-13 s period band, largest ratios at the low-latitude end. This means, 330 our expectation of large ratios near the cross-over is obviously not confirmed by the southern 331 hemisphere observations. Rather, it seems that the season, summer in the norther hemisphere, is 332 responsible for the preferred appearance of stable small-scale FACs at high latitudes. Conversely, 333 in the southern, winter hemisphere these small-scale high-latitude FACs are missing, as is obvious 334 from the blank parts in the plot near the pole. The effect of cross-track distance on this period band is stronger at lower latitudes. The larger distances, cause more reduced occurrence ratios in the 335 morning and prenoon sectors than at afternoon to evening. 336

337 The second part of the study period exhibits more dynamic variations. Most prominent is the 338 obvious drop in ratio during the time of extended along-track separation, Δt . As mentioned already 339 above, the reduction is most pronounced for the small periods (short scales) and becomes less 340 prominent the longer the period. On the upleg passes, covering the noon, afternoon sector, the 341 reductions in ratio are quite evenly distributed over the latitudes. Conversely for the downleg, sampling night and early morning hours, there is a clear preference of ratio reduction at higher 342 343 latitudes. Differently, the lower-latitude auroral FAC structures (e.g. R2 FACs) in this sector seem 344 to be stable for more than 40 s.

345 Worth mentioning are also the two narrow features of low ratios around the beginning of 346 November 2021. They represent contrasting activity conditions, the quietest period (27-29 October 347 2021) and the most intense storm on 3-4 November 2021 of the study period. In an attempt to 348 explain the reasons for the low ratios we had a look at the cross-correlation plots from the two very 349 different days. An example of a pass from the quiet day over the southern auroral region is shown 350 in Figure 7, top frame. We find moderate activity in the morning and night sector at fairly high 351 latitudes, around 70° MLat. For most parts of the orbital arc high correlation coefficients are 352 obtained at the right time lag. Just the burst at night-time makes an exception. In spite of the 353 primarily stable FAC structures on this quiet day, overall low ratios are obtained. This is caused 354 by the relatively small amplitude of the signal, falling over large parts below our threshold of RMS 355 > 2 nT/s. A quite different picture emerges from the stormy day on 4 November 2021. The example 356 in the lower part of Figure 7 shows a northern hemisphere polar pass. Bursts of activity appear in the evening and morning sectors. Their amplitudes are an order of magnitude larger than those on 357 358 the quiet day. The correlation coefficient generally stays below the threshold, Cc = 0.75, and also 359 the optimal time lag is in most cases too short in comparison with the spacecraft separation. 360 Another but minor effect is the equatorward expansion of the auroral activity during the storm to 361 latitudes below 60° MLat and with that a movement out of our range activity monitoring. Large parts of the greatly expanded polar cap are practically free of signal. All this results in very few 362 positive detections of stable FAC structures during this intense storm. 363

All the observations presented so far indicate that there are several processes that influence the
stability of small- and meso-scale FAC structures. When looking at series of satellite observations,
several of the controlling parameters vary at the same time. Therefore, it is not easy to disentangle
the effects.

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369 5. Discussion

The observations presented in the previous sections show that the correlation length of FAC structures have different dependences on spatial and temporal scales. In addition, also local time, season, latitude range and solar wind input may play a role. In this section we try to disentangle the various influences that determine the stability of a FAC structure. No effort has been undertaken so far to achieve this in comparable detail.

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376 5.1 General characteristics of small- and meso-scale FACs





377 In a first step we look at the spatial scales and temporal lengths of stable FACs. Observations presented in Figure 5 indicate that the percentage of detected stable current structures depends on 378 379 both the along- and cross-track separations between the Swarm satellites. For further investigating these dependences, Figure 8 presents the occurrence ratios for all the six band passes in d_{cross} 380 versus Δt frames. The individual ratio percentages have been dropped into bins of 1 km for d_{cross} 381 382 by 3 s for Δt . Bin averages are then shown in color. Results from both hemispheres are combined. 383 Due to the given orbit geometry, the cross-over points are generally close to the geographic poles. 384 Thus, only relatively small transverse separations are experienced in this high-latitude study. 385 Within the range of $d_{cross} = 0 - 15$ km, as shown in Figure 8, we cover almost all the available cases (only a few reach up to $d_{cross} = 20$ km at $\Delta t = 6 - 9$ s). For classifying a current event as 386 387 suitable for dual-SC FAC estimates we request occurrence ratios larger than 50%, as done in the 388 earlier study by Zhou et al. (2024).

In general, Figure 8 confirms the impression gained from Figure 5, in particular for short period
structures (small FACs); the occurrence ratios drop off rapidly away from the epoch of coplanarity.
It is quite clear from Figure 6 that the patterns of occurrence ratios differ significantly between the
upper and lower groups, the three shortest periods and the two longest, respectively, while the 13
- 23 s band contains something from both groups. Here we will go a little more into the details of
the different characteristics observed for the two groups of FAC scale ranges.

395 The zonal correlation length of the meso-scale FACs (23-60 s) is obviously larger than the 396 experienced satellite cross-track separations (0 -15 km). Therefore, the time between samples is 397 more decisive for the occurrence ratio. Also here, larger structures are stable over a longer times. 398 Generally, the 50% demarcation lines in Figure 8 lies for the lower two period bands beyond the covered Δt range. This confirms earlier studies, e.g. Gjerloev et al. (2011); Lühr et al. (2015), who 399 400 reported stability periods of 60 s and more for FACs of these scale sizes. With the given limitations, 401 the present study cannot provide much new information about the character of these meso-scale 402 FACs at auroral latitudes. Still, it should be noticed, these results confirm well the applied 403 preprocessing approach of the magnetic field recordings for the Swarm standard dual-SC FAC 404 processing. For the calculations of those products the horizontal field components are low-pass 405 filtered with a cutoff period of 20 s. This suppresses the variable small-scale signals. The along-406 track separation between Swarm A and C varies during normal operation between 4 and 10 s. In 407 addition, there are the step sizes of 5 s, forming the integration quad (see Ritter et al. (2013) and 408 Lühr et al. (2020) for more details). This means, measurements contributing to a dual-SC FAC 409 density can be separated by up to 15 s. According to our results, the field recordings at the two 410 spacecraft are still well correlated with each other. Figure 8 confirms that around 90% of the FAC 411 structures with periods longer than 20 s are stable over even longer times than the 15 s.

412 The stability characteristics of structures with periods of 3 to 13 s are quite different. Their cross-413 track size seems to be the more limiting factor than the temporal stability. For example, in the case 414 of the 3-7 s band, for small satellite separations, $d_{cross} < 6$ km, we find ratios of about 50% or more 415 up to $\Delta t \sim 16$ s. Conversely, for larger cross-track separation, say $d_{cross} > 8$ km, we find nowhere 416 50% ratio levels for $\Delta t > 4$ s in this period band. Also in the 7-13 s bands, we find a similar tongue 417 of elevated ratio levels for small spacecraft separations, $d_{cross} < 6$ km. This indicates that a majority 418 of small-scale FACs exhibits such a narrow transverse sizes.

419 The FAC structures in the shortest period band seem to belong to another, third category of current 420 features. Outside the $\Delta t = 2$ s separation range only very few common magnetic features are 421 observed by the two Swarm satellites. These small FACs with horizontal scale of less than 10 km 422 seem to be stable only for a second or so. Thus, the longitudinal separation of the spacecraft has 423 no effect on the derived occurrence ratio. This type of small FAC is surely worth a more detailed 424 investigation and may be the topic of a follow-up study.





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426 Already Ishii et al. (1992) had reported about two types of FACs deduced from the ratio ΔB_{east} over E_{north} . The one, related to longer periods in satellite recordings, was identified as stable 427 428 current structures. For such FACs it can be assumed that the field lines act as equal potential lines 429 between magnetosphere and ionosphere. In those cases, the ΔB over $\mu_0 E$ ratio reflects the Pedersen 430 conductance. These authors mentioned a short-period limit of 8 s for the stable FACs. Here we have identified a period around 15 s as the limit of stable meso-scale FACs. Although based on 431 432 very different approaches, both studies come to consistent results for the typical scale range of 433 stationary FACs. For shorter period (smaller) current structures Ishii et al. (1992) report a 434 progressive increase of the E-field amplitude above that of ΔB down to transverse scales of a few 435 kilometers. They call the FAC type observed here the transient Alfvénic mode. For even shorter 436 scales the pure Alfvén mode is approached with $E/B = V_A$, where V_A is the Alfvén velocity. In the following we will focus on the FACs in the transient Alfvénic scale range. 437

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439 5.2 Features of small-scale FACS

440 For learning more about the small FACs at auroral latitudes we looked into the distribution of 441 stable events. Figure 9 shows in the upper two rows how the amplitude of the magnetic variations 442 for selected FAC structures is distributed in magnetic latitude over our study period. The 3-7 s period band (~20-50 km wavelength) is used as example for this class of FACs. Only positive 443 444 event detections within the above-described range, $d_{cross} < 6$ km and $\Delta t < 18$ s, are considered. 445 Separate frames display the results from up- and downleg orbital arcs and from the two 446 hemispheres. The blank areas, in particular at lower latitudes and times of large spacecraft 447 separations, reflect the lack of entries due to our spatial/temporal constrains. Prominent stable current structures appear in the northern hemisphere near 80° MLat in the morning to noon sector 448 449 during the first half of the study period. At the same time, much less activity is observed in the 450 southern hemisphere with smaller amplitude and at somewhat lower latitude. We attribute this 451 hemispheric differences primarily to the season, with local summer in the north and winter in the 452 south. During night-time hours events appear at much lower latitudes and significantly smaller 453 amplitudes. The second half of the study period is rather sparsely populated due to our constrains. 454 Here we have southern summer conditions. As expected, more significant FAC activity appears in 455 this hemisphere on upleg orbital arcs during morning and prenoon hours. Although we have winter 456 in the northern hemisphere, it still shows appreciable small-scale FAC amplitudes in this time sector. The downleg arcs exhibit in the south some moderate FAC activity in the auroral region 457 458 during pre-midnight hours. In the northern hemisphere the sampling is too sparse in this time 459 interval for providing additional information.

In comparison to the identified stable small FACs, the two lower rows of Figure 9 shows the 460 amplitude distribution of the deselected current events within the above-described range, d_{cross} < 461 6 km and $\Delta t < 18$ s in the same format. An obvious difference between the two distributions is the 462 463 generally larger amplitude of the deselected events. This difference in amplitude is more 464 pronounced on the nightside than on the dayside. On the nightside obviously only the really large 465 FAC spike are unstable. The latitude distribution of deselected events is clearly narrower than that of the selected. This is in support of the preferred stability of the smaller amplitude FAC structures 466 467 located away from the latitude of peak small-scale FAC activity. These facts suggests that the 468 larger amplitude FACs tend to have smaller temporal and spatial correlation lengths. Here we 469 should add a note. Not all the deselected events represent unstable current structures. In several 470 cases one satellite may sample a stable FAC structure at the fringe while the other samples a 471 different structure. That event will falsely appear under deselected. From the experience collected





so far it probably concerns predominately the small-amplitude events that are falsely classified asdeselected.

474 From the results presented in the upper rows of Figure 9 we suggest that the class of stable smallscale FACs detected in the range, $d_{cross} < 6$ km and $\Delta t < 18$ s, represent the majority of auroral 475 476 zone FACs occurring in this size. In actual numbers, we detected about 230,000 stable FAC events in this very limited range, compared to a total of 340,000 stable events distributed over the whole 477 478 spatial and temporal range. This fact provides additional evidence for the small correlation length 479 of this type of FACs. Earlier studies of small-scale FACs, e.g. Kervalishvili and Lühr (2013), based 480 on 5 years of CHAMP magnetic field data, found mean characteristics that resemble well this class in the 3-7s band. They report a clear preference of summer season for the noon-time amplitudes at 481 482 cusp latitudes. A secondary activity peak appears in the night sector, which is less dependent on 483 season, and lower activities show up on the dawn and dusk sides. Following the probability 484 considerations, as outlined by Xiong and Lühr (2023) and Zhou et al. (2024), the mean correlation 485 length of the current structure is expected to be twice as long as the spacecraft separation at the 486 50% occurrence ratio. In our case this means a distance of 12 km in azimuthal direction. From 487 these results we may conclude that the common small-scale FACs at auroral latitudes are of 488 filamentary character with short correlation lengths of the order of some 10 km.

489 Besides the above-described class of short-scale current structures there seems to be a minority of 490 stable small-scale FACs exhibiting larger correlation lengths. In Figure 8 we find a horizontal band 491 of slightly enhanced occurrence ratios (10%-20%) for $d_{cross} = 12-15$ km in the period band of 3-7 492 s. In this cross-track range the ratios vary only little with the along-track separation between the 493 satellites. Such large cross-track distances are only experienced near the lower border, 60° MLat, 494 of the considered latitude range. This second population of stable small-scale FACs is obviously 495 a subauroral phenomenon, distinctly different from the first that prefers the cusp region and high 496 latitudes. It would require additional investigations to characterize the features of this second 497 population.

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499 5.3 Characteristics of unstable small-scale FACs

From Figure 8 it is evident that even in the preferred small-scale FAC range, $d_{cross} < 6$ km and Δt 500 501 < 18 s, of the 3-7 s period band detection ratios reach at best 60%. That means, about every second 502 magnetic signature even in this favorable range is deselected. This is partly expected for a structure 503 of finite size when one satellite passes close to the edge and the other samples another structure. 504 There may, however, also be other reasons for deselection. For checking that we show in the lower 505 two rows of Figure 9 the amplitude distribution of deselected events from the same spatial-506 temporal range as the selected above. When comparing the upper and lower pairs of rows, panel by panel, distributions are narrower in latitude for the deselected cases. The regions of deselected 507 508 events are mainly populated by larger amplitude variations. Mean amplitudes of less than 4 nT/s 509 are largely missing there. Thus, large-amplitude small-scale FACs are less stable than the smaller 510 amplitude FACs, and this seems to be valid for both day and nightside events.

511 In the quest for reasons that favor unstable current structures we checked prevailing IMF and solar 512 wind conditions during deselected events. As a parameter for representing the amount of solar 513 wind input, we choose the merging electric field, E_m , as defined by Newell et al. (2007) (they call 514 it coupling parameter)

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$$E_m = \frac{1}{3000} V_{SW}^{\frac{4}{3}} (\sqrt{B_y^2 + B_z^2})^{\frac{2}{3}} \sin^{\frac{8}{3}} (\frac{\theta}{2})$$
(3)





516 where V_{SW} is the solar wind velocity in km/s, B_y and B_z , both in nT, are the IMF components in 517 GSM coordinates, θ is the clock angle of the IMF. With these units the value of the merging 518 electric field will result in mV/m. Here solar wind and IMF data, averaged over 1 min and 519 propagated to the bow shock, are used to obtain a weighted time-integrated merging electric field 520 which accounts for the memory effect of the magnetosphere-ionosphere system. Details of the 521 approach can be found in the publication of Zhou et al. (2018).

For investigating the conditions accompanying our events we perform a superposed epoch analysis of the smoothed merging electric field. For the epoch of each event in the 3-7 s period band, occurring in the $d_{cross} < 6$ km and $\Delta t < 18$ s range, the E_m value evolution is considered from 30 min before the event to 30 min after it. The E_m time series from all the events are stacked and averages are calculated of each 1-minute bin. This provides the mean E_m , evolution around a certain class of small-scale FAC events.

Figure 10 shows the superposed epoch analysis results of the mean E_m evolution over the hour 528 around the key time, $\delta T = 0$, for selected and deselected current structures as blue and red dotted 529 530 lines, respectively. The two columns in Figure 10 represent results from the northern hemisphere on the left and from the southern on the right side. The curves in the top two panels are from the 531 dayside, the ones in the two bottom panels are from the nightside. Seasonal effects may show up 532 533 in a comparison between observations from days before and after 1 October 2021 (the epoch of 534 coplanarity). Generally, we find larger E_m values accompanying deselected events than selected. 535 Independent of season, hemisphere and local time, the level of E_m value ranges between 0.8 and 1 mV/m for the positively selected events. This represents well the mean merging electric field of 536 537 quiet times during the considered months. Conversely, we obtain $E_m > 1$ mV/m for deselected events. Just during the months after 1 October 2021 E_m values below 1 mV/m are found in the 538 539 northern hemisphere dayside for deselected events. This is cause by a generally low solar activity 540 during the related days. Overall, the superposed epoch analysis confirms our earlier finding that 541 large amplitude FAC structures are prone to instability, and the intensity of current density is 542 largely dependent on the amount of solar wind driving. The evolution of the E_m values around the 543 key time, $\delta T = 0$, does not provide a clear picture of the preconditions causing unstable FACs. The eight panels show various kinds of rising and falling trends of the red lines. Common to all is the 544 elevated E_m value. 545

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547 5.4 The roles of the deselection criteria

548 There is another question that we may address here. What kind of current variation causes the 549 deselection of an event? The criteria for a stable current structure are, an amplitude RMS > 2 nT/s550 and a peak cross-correlation coefficient above 0.75 at a time lag agreeing with the along-track separation of the satellites within 1.5 s. By checking, which of the criteria causes the deselection, 551 552 one may obtain information about the kind of variation. A general finding is that a low Cc is 553 involved in the large majority of deselected events. In the cases of a very low Cc, it cannot be 554 expected that the right T-lag is found. Therefore, a dropout of both criteria can also frequently be found. From this observation we infer that the large percentage of deselected small-scale events 555 556 (up to period of 15 s) (see Fig. 8) is caused by a temporal or spatial change of the current signature between the visits by the two satellites. 557

558 In the case of meso-scale FACs high occurrence ratios, partly close to 100% (see Fig. 8), are 559 observed. For large along-track separations the percentage drops somewhat. In order to find an 560 explanation for that behavior we looked at the percentage of deselected events based solely on the 561 The mitation for that behavior the metated action distribution in formers of d

561 T-lag criterion. Figure 11 shows the related ratio distribution in frames of d_{cross} over Δt , separately





562 for the six period bands. The color scale represents the percentage of deselection by T-lag relative to the total number of deselected events. As mentioned before, the T-lag criterion plays hardly any 563 564 role for the small-scale FACs. Corresponding ratios vary for periods up to ~ 15 s around 10%. We 565 find a quite different picture for the meso-scale FACs, in particular for the longest period. The white gaps represent bins that contain no deselected events. This generally indicates the low count 566 numbers of deselections on which these ratios are based. The fairly high percentage caused by T-567 lag, varying for the long periods, around 80%, indicates that a motion of the current system along 568 569 the orbit track is the main cause here. This means, the FAC pattern does not vary in time, but the 570 system is moving equator or pole ward. Such motions are well-know from auroral observations. 571 The ratios for deselects by T-lag become progressively smaller for shorter periods, but the absolute number and distribution of these deselected events is quite similar over the period range 13 s - 60 572 573 s. That means, the latitudinal propagation of the current system is similar for the meso-scale FAC 574 structures of different sizes. However, the change in dominant color from period to period indicates 575 a clear reduction of stability, also for meso-scales, in space and time increasing towards smaller 576 sizes.

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578 6. Summary and Conclusions

579 In this study we made use of magnetic field data from the closely spaced Swarm A and C spacecraft 580 during the counter rotating orbit phase. They allowed us to investigate the spatial and temporal 581 correlation lengths of small-scale and meso-scale FAC structures at auroral latitudes. A number of 582 detailed characteristics of small-scale FACs (10 - 70 km wavelength) have been identified for the 583 first time. In particular, their stability features could be determined. Furthermore, the stationarity 584 of meso-scale FACs (100 - 300 km wavelength) has been confirmed. The main findings of the 585 study are:

- 1. The meso-scale FACs exhibit transverse (azimuthal) correlation lengths larger than our 586 maximum spacecraft of 20 km. Therefore, no limit value can be provided here. Generally, 587 588 large fractions of the current structures are identified as stable. Figure 5 shows for this class of FACs detection ratios close to 100% over most parts of the study period. Larger deflections 589 590 from this level are confined to times of long along-track separations (up to 41 s) between the spacecraft. We identified the latitudinal motion of the current system as the main cause for the 591 deselection of the events, not the temporal variation. All these results are consistent with 592 593 earlier findings.
- 594 2. More interesting and new results have been obtained for the small-scale auroral FACs. For most of them an azimuthal correlation length of less than 15 km seems to be an important 595 limit. A characteristic time period of less than 18 s has been obtained for their temporal 596 stability. The distribution of current features identified in this limited spatial and temporal 597 598 range agrees well with those of earlier studies about small-scale FACs. This includes their 599 distribution in local time, magnetic latitude, and season. From these similarities we infer that 600 our limits are typical for small-scale auroral FACs. However, during enhanced solar wind driving unstable small-scale FAC structures appear even in this limited spatial and temporal 601 602 range.
- 6033. An investigation of deselected small-scale current structures reveals that in particular the604large-amplitude FACs are unstable. This is valid for the dayside and even more obvious on605the nightside. In support of this fact, we could show that deselected events generally follow606more intense solar wind driving. Stable current structures are accompanied by merging607electric fields, E_m , with the value ranging between 0.8 and 1 mV/m, typical for low solar wind608driving.
- 609 4. There exists another population of small-scale FACs that is stable over longer time, $\Delta t > 40$ s. 610 Also, its azimuthal correlation width seems to be longer than 15 km. The preferred appearance





611 is at the equatorial end of our latitude range, near 60° MLat. No detailed investigations of
612 these subauroral FACs have been performed, but they seem to be worth considering in a
613 follow-up study.

From the results listed here we can conclude, the large-scale FACs are stable well over 40 s. 614 Reasons for deselection are mainly the latitudinal motion of the current system, while its shape 615 616 remains constant. Meso-scale FACs (100-300 km wavelength) also show the effect of latitudinal 617 motion, but in addition, the shape starts to vary within 40 s. The smaller structures exhibit more 618 variation. The small-scale auroral FACs exhibit quite different characteristics. They behave more 619 like waves. Stable events are only found in a limited range of spacecraft separations both in time 620 and spatial differences. Their stability is shorter than the Alfvén transition time from the magnetosphere to the ionosphere. The small-scale FACs with large amplitudes, e.g. during active 621 622 periods (see Fig. 7 lower frame), have very short correlation lengths, falling below the resolution 623 of this study. They seem to represent narrow solidary current structures that propagate along the field line in form of an Alfvén wave. Such intense narrow FACs can be found on the day and 624 625 nightside. We may speculate that they are caused on the dayside by patchy reconnection and on 626 the nightside by the bursty bulk flow. Multi-point measurements with higher spatial and temporal 627 resolution would be needed to resolve the details of these intense solidary FACs.

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634 Data Availability Statement

The authors thank the European Space Agency for openly providing the Swarm data. The data
products used in this study are Level-1b MAGx_LR with version number 0602, which are
available at the European Space Agency website: https://earth.esa.int/web/guest/swarm/data-
access.

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640 Author contributions

The study concept and the first draft of the manuscript were contributed by HL. Data collectionand analysis were performed by YLZ. Both authors finalized the manuscript.

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644 **Competing interests**

645 The corresponding author declares that none of the authors has a competing interests.

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- 732 Geophysical Research: Space Physics, 123, 5159–5181. https://doi.org/10.1029/2018JA025526







735 Figure 1. Examples of magnetic variations in the transverse, ΔB_{trans} , component within the period 736 range of 1-60 s. The top panels of the two frames show the recordings of Swarm A and C along 737 their orbits, crossing the polar region of the northern hemisphere. The second panel reflects the RMS value of the signal amplitude. The third panel contains the peak cross-correlation coefficient, 738 739 Cc. Significant correlations between Swarm A and C should have values, Cc > 0.75. The dashed 740 green lines mark the thresholds of the three detection criteria. The bottom panel shows the lag 741 time, T-lag, between the signals at optimal correlation. T-lag values between the two green dashed lines indicate delays equal to the time separation between the spacecraft. The top frame is from a 742 743 time with small along-track separation, $\Delta t = 4$ s and the bottom frame from a larger separation, Δt 744 = 22.3 s





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Figure 2. Occurrence distribution of transverse magnetic signature variations in the auroral region with amplitudes above a certain RMS value. Here, signals in the period range 1- 60 s (7.5 - 450

750 km along-track scale size) have been considered. For our FAC study amplitudes, RMS > 2 nT/s,

751 are required. With that, about 50% of the signals are taken into account.







Figure 3. Temporal variation of all wave structures in the ΔB_{trans} component with amplitudes of RMS > 2 nT/s over the full study period, separately for the 6 period ranges and the two hemispheres. Results from the orbital upleg arcs at high latitudes (>60° MLat) are shown as red curves, those from downleg arcs are in blue. Across the bottom, besides the date, the local times at 70° MLat of the orbital arcs are listed, as well as the along-track time difference, Δt in seconds, between the spacecraft and the varying longitudinal separation at the equator, $\Delta Long$, in degree.







Figure 4. The same format as Figure 3, but for the occurrence frequency of positively detected static FAC structures.







Figure 5. The same format as Figure 3, but for the ratio of detected static FAC structures (Fig. 4)

divided by all the number of wave structures presented in Figure 3.







all wavy signals above the threshold of RMS > 2 n1/s, separately for all the six period bands, upand downleg orbit arcs, and the two hemispheres. The labels along the horizontal axis are as described for Figure 3.







Figure 7. Magnetic variations in the transverse, ΔB_{trans} , components within the period range of 1-60 s in the same format as Figure 1. The top frame is from a very quiet day, while the bottom frame presents variations of the storm period on 4 November 2021. Peak amplitudes differ by a factor of 10 between the two days.



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Figure 8. Distributions of stable FAC occurrence ratios in frames of cross-track distance, d_{cross} ,versus along-track time separation, Δt , separately for the six period bands. There are hardly anymeasurements for $d_{cross} > 15$ km.







801 $d_{cross} < 6$ km and $\Delta t < 18$ s range are considered. The top two rows present the distribution of stable 802 current structures and the bottom two row the deselected events. White areas mark ranges without 803 entries.







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Figure 10. Superposed epoch analysis of the mean merging electric field, E_m, evolution accompanying the stable small-scale structures (blue lines) and deselected events (red lines). The relative time indicates the behavior before and after the FAC event detection. Individual analyses have been performed for dayside and nightside, for months around June and December solstices and for the auroral regions in the two hemispheres. The irregular curve of deselected events in the northern hemisphere after October is cause by the really small number of cases (see also Fig. 9).



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814Figure 11. The same format as Figure 8, but for deselected events in the $d_{cross} < 6$ km and $\Delta t < 18$ 815s range based solely on the T-lag criterium. White patches represent bins without any deselected816events. The T-lag criterium is of significance only for the meso- and large-scale FACs. Low cross-817correlation coefficients clearly dominate the deselection of small-scale current structures.