

# 1 Investigating the Role of Gravity Waves on Mesosphere-Lower-Thermosphere

## 2 (MLT) Inversions at Low Latitudes

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### 7 **Abstract**

8 The Mesosphere and Lower-Thermosphere (MLT) transitional region, encompassing a height  
9 range of 60-100 km, is a distinct and highly turbulent zone within the Earth's atmosphere. The  
10 region is significant owing to dynamics of atmospheric processes like planetary, tidal, and  
11 particularly gravity waves, which contribute to the formation of the Mesospheric Inversion Layer  
12 (MIL). Investigating the inversion phenomena is crucial for understanding the dynamics of the  
13 middle and upper atmosphere, especially regarding stability and energy transfer. These  
14 phenomena are associated with energy transfer processes, vital for understanding the overall  
15 dynamics of the atmosphere. Despite extensive study on an inversion, the formation mechanisms  
16 of mesospheric inversions remain poorly understood. Hereunder, the upper and lower inversion  
17 phenomena and their causative mechanisms are explored. It utilizes long-term SABER  
18 observations during 2005-2020 over the latitude, 3-15° N and longitude, 33-48° E ranges. The  
19 results show that the upper inversion occurs more frequently, with a frequency below 40%,  
20 compared to the lower inversion, which occurs below 20%. The upper inversion occurs within  
21 the height range of 78-91 km, with an inversion amplitude of approximately 20-80 K and a  
22 thickness of around 3-12 km. In contrast, the lower inversion is confined to the height range of  
23 70-80 km, with an inversion amplitude of about 10-60 K and a thickness of around 4-10 km.  
24 Moreover, the gravity wave indicator potential energy shows high energy (below 100 J/kg) in the  
25 upper MLT region (85-90 km) compared to the lower MLT region (70-75 km) with less than 50  
26 J/kg. Considering gravity waves, the Brunt-Väisälä frequency ( $N^2$ ) stability criteria indicate  
27 instability in the upper MLT region with very low values compared to the lower MLT region.  
28 This suggests that the high amount of gravity wave potential energy is a consequence of the  
29 higher instability in the upper inversion compared to the lower inversion.

30 **Keywords:** Mesosphere and Lower Thermosphere (MLT), Upper and Lower Inversions,  
31 Perturbed Temperature, Causative Gravity Waves, Potential Energy, Brunt-Väisälä Frequency,  
32 Atmospheric Instability.

## 33 **Introduction**

34 The Mesosphere and Lower-Thermosphere (MLT) region serves as a transitional zone for  
35 atmospheric wave processes from the lower and upper atmospheres, including tidal, planetary,  
36 and gravity waves. Gravity waves (GWs) originating in the lower atmosphere propagate into the  
37 upper mesosphere, where they break and dissipate, releasing energy and momentum. This  
38 process influences the thermal structure, global atmospheric circulation, and mesospheric  
39 inversion layers (MILs), which are associated with increased temperature variability in  
40 mesosphere. MILs indicate wave saturation when the lapse rate falls below the dry adiabatic  
41 lapse rate (Sica et al., 2007). Temperature inversions in the mesosphere have been widely  
42 observed and studied using various techniques, including lidar, radar, rocket sondes, and  
43 satellites, across different geographic locations. Sivakandan et al. (2014) utilized  
44 TIMED/SABER kinetic temperature data to examine the occurrence and characteristics of  
45 mesospheric inversions over the equatorial Indian region (0 to 10° N and 70 to 90° E) for the  
46 years 2002 and 2008. However, they did not explore the causative factors. This study aims to  
47 investigate the causes of these inversions, focusing specifically on the role of atmospheric  
48 gravity waves.

49 Gravity waves and mesospheric inversion layers (MILs) are interconnected phenomena within  
50 the Earth's atmosphere, particularly in the mesosphere and lower thermosphere (MLT).  
51 Inversions are layers within the mesosphere where the temperature profile exhibits an increment.  
52 As a result, the temperature increases with altitude, contrary to the typical decrease. These  
53 inversion layers often form because of this atmospheric wave dynamic processes, including the  
54 breaking and dissipation of gravity waves. As gravity waves propagate upwards, they can grow  
55 in amplitude since the atmospheric density decreases with altitude. While these waves reach a  
56 critical amplitude, they become broken. This breaking process releases energy and momentum  
57 into the surrounding atmosphere, leading to localized heating that creates or enhances  
58 mesospheric inversion layers by increasing the temperature with altitude. The breaking of gravity  
59 waves contributes to momentum and energy deposition can also generate turbulence, which  
60 further influences the structure and instability of an inversion layers.

61 The deposition of momentum and energy from gravity waves (GWs) is considered a major factor  
62 driving large-scale atmospheric circulation, the coupling between atmospheric layers, and the  
63 occurrence of inversion phenomena (Fritts and Alexander, 2003; Lindzen, 1981; Smith, 2012).  
64 Researchers are also investigating the effects of gravity wave breaking on MLT dynamics to

65 better understand its role in inversion phenomena, especially in mid- and high-latitude regions  
66 (Gan et al., 2012; Walterscheid and Hickey, 2009; Collins et al., 2011; Szewczyk et al., 2013).  
67 Observational and modeling studies have examined gravity waves (GWs) as a contributing factor  
68 to these inversions (Fritts, 2018; Collins et al., 2014; Sridharan et al., 2008; Ramesh and  
69 Sridharan, 2012; Ramesh et al., 2013, 2014, 2017). Despite extensive research, our  
70 understanding of how gravity waves influence mesosphere inversions-particularly regarding  
71 temperature variability-remains incomplete, even in mid- and high-latitude regions (Singh and  
72 Pallamraju, 2018; Fritts et al., 2018). Consequently, studying inversion phenomena and their  
73 underlying causes continues to be a crucial area of investigation, especially within MLT  
74 dynamics over low-latitude.

75 Research on the temporal and spatial variability of the mesosphere inversion phenomenon,  
76 regarding atmospheric waves, particularly gravity wave activity, is notably lacking in low  
77 latitudes. To address this gap, our study investigates the mesosphere inversion phenomenon and  
78 its association with gravity wave activity and instability criteria. We use Brunt-Vaisala frequency  
79 ( $N^2$ ) over the low latitudinal band ( $3^\circ$ – $15^\circ$  N) with long-term SABER observations from 2005 to  
80 2020. The study is organized as follows: Section 2 details the data and methodology used to  
81 analyze the mesosphere inversion phenomenon and their causative gravity waves via the  
82 potential energy. Section 3 presents the results, and Section 4 concludes the findings.

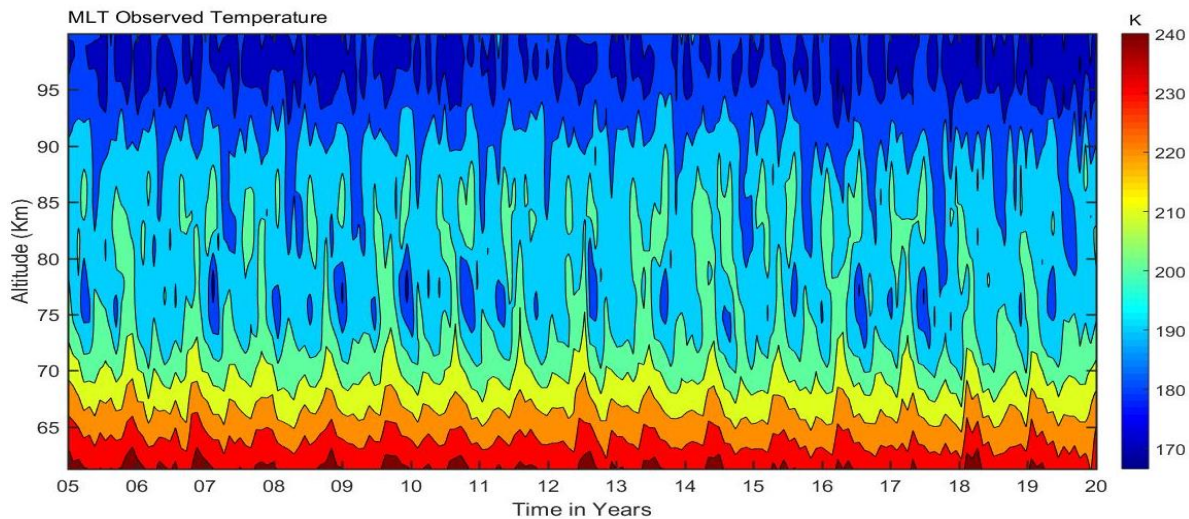
## 83 **2. Observation and Data analysis**

### 84 **2.1 SABER Observation**

85 The TIMED/SABER satellite, launched on December 7, 2001, operates in an elliptical orbit at  
86 approximately 625 km altitude with a  $74^\circ$  inclination relative to the equator. Since its launch,  
87 SABER has been a crucial tool for atmospheric research, providing extensive data on the middle  
88 and upper atmosphere. SABER is a limb-viewing radiometer working in the infrared region  
89 (1.27–17 microns) and can measure radiative emissions across a wide range of altitudes. It offers  
90 nearly global coverage and continuous 24-hour data over 60 days. The instrument completes 15  
91 orbits daily, each taking about 97 minutes, and collects around 1400 data files per day, with each  
92 profile taking 58 seconds. SABER's high-resolution temperature profiles are essential for  
93 studying the dynamics and wave processes in the MLT. It provides temperature measurements  
94 with an accuracy of 1 to 2 K between 15 and 60 km. The accuracy decreases to 5 K below 85 km  
95 and increases to 6.7 K to 10 K near 100 km. This data has been crucial in enhancing our  
96 understanding of the thermal structure and dynamic processes in the mesospheric region, as

97 emphasized by several studies (Garcia et al., 2008; Gan et al., 2012, 2014; Bizuneh et al., 2022;  
98 Lingerew et al., 2023; Rezac et al., 2015; Meriwether and Gerrard, 2004; Fechine et al., 2008;  
99 Dou et al., 2009; France et al., 2015).

100 Owing to this, we used the SABER vertical temperature profiles collected within the 60–100 km  
101 altitude range. These profiles encompass the period from 2005 to 2020, covering latitudes from  
102 3°N to 15°N and longitudes from 33°E to 48°E. Figure 1 illustrates the monthly mean SABER  
103 temperature data for the mesosphere and lower thermosphere. The data aim to illustrate the MLT  
104 temperature variability, which helps us identify the inversion layers (MIL). The monthly mean  
105 temperatures in the MLT region show a maximum of 200-240 K at altitudes of 60-70 km. Then it  
106 decreases to around 160-180 K at 95-100 km throughout the entire period. While the temperature  
107 patterns in the 70-90 km altitude range suggest an inversion, these inversions are not visible.



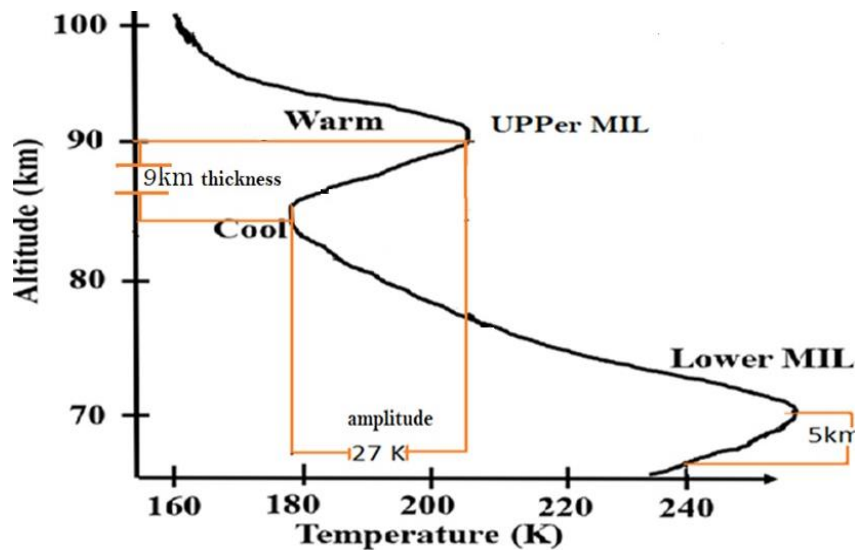
108 **Figure 1.** The monthly mean of MLT temperature variability in the height range of 60-100 km  
109 during 2005-2020 over the low-latitude.  
110

## 111 **2.2 Analysis Technique**

112 The Earth's middle atmosphere typically has a negative temperature gradient, but some reports  
113 have shown positive temperature gradients in the mesosphere (Meriwether and Gardner, 2000;  
114 Gan et al., 2012). This phenomenon, known as the "mesospheric inversion layer (MIL)," is  
115 identified using the method described by Leblanc and Hauchecorne (1997) and Fechine et al.  
116 (2008). Mesospheric inversions are defined by their thickness (the altitude difference between  
117 the maximum warming and cooling) and their amplitude (the temperature difference between

118 these points) (Meriwether and Gardner, 2000). The following are the criteria for identifying these  
119 inversions:

- 120 1. The bottom level of the lower inversion is above 70 km, and the top level is below 80 km.  
121 For the upper inversion, the bottom level is above 80 km, and the top level is below 92  
122 km.
- 123 2. The amplitude is considered larger than 5 K.
- 124 3. The thickness is greater than or equal to 3 km.



125  
126 Figure 2. Schematic of upper and lower mesospheric inversion layers shown in the temperature  
127 profile for the MLT regions (Adapted from Meriwether and Gerrard, 2004).

128 Figure 2 illustrates this concept, highlighting the positive temperature difference between the top  
129 and bottom levels of the inversion. This method has been widely applied in numerous studies  
130 investigating mesospheric inversions (Leblanc et al., 1998; Meriwether and Gardner, 2000; Duck  
131 et al., 2001; Duck and Greene, 2004; Cutler et al., 2001; Siva Kumar et al., 2001; Ratnam et al.,  
132 2003; Gan et al., 2012). In addition, the frequency of mesospheric inversion layer (MIL)  
133 occurrences is determined for the period 2005–2020 in both the upper and lower MLT regions.  
134 This frequency is calculated by dividing the number of inversion days in each month by the total  
135 number of days in that month over the 16-year observation period (2005–2020).

136 Mesospheric temperature inversions are linked to MLT instabilities driven by the dynamics of  
137 atmospheric wave processes. To identify the causative, short-period atmospheric gravity waves,  
138 a high-pass filter with a one-hour interval cutoff frequency is applied using the Brunt-Väisälä

139 frequency (N2). Another important concept to estimate the Brunt-Vaisala frequency is the  
 140 potential temperature ( $\theta$ ). It represents the air parcel's temperature when it is displaced  
 141 adiabatically to a standard pressure level,  $p_0$ , from the current pressure level,  $p$ . This is based on  
 142 the first law of thermodynamics.

$$143 \quad \frac{dT}{T} = \frac{R}{c_p} \frac{dp}{p} \Rightarrow \int_T^0 \frac{dT}{T} = \int_p^{p_0} \frac{R}{c_p} \frac{dp}{p} \quad (1) \text{ it yields}$$

$$144 \quad \theta = T \left( \frac{p_0}{p} \right)^{R/c_p} \quad (2)$$

145 The vertical motion of an atmospheric air parcel can thus be described by (Liu, 2011; Vadas and  
 146 Fritts, 2005) as shown in equation (2). This equation calculates the Brunt-Väisälä frequency of  
 147 the parcel, accounting for the buoyant and gravitational forces acting upon it.

$$148 \quad \frac{d^2s}{dt^2} = -g \frac{\rho - \rho_0}{\rho} \sin a \quad (3)$$

149 Based on the hydrostatic equation,  $\rho = \rho_0$ , and  $p = p_0 \Rightarrow \frac{\partial p}{\partial z} = \frac{\partial p_0}{\partial z} = -g\rho_0$  (4) and the ideal gas  
 150 law,  $\rho = p/RT = p_0/RT$  gives the parcels motion of an equation:

$$151 \quad \frac{d^2s}{dt^2} = -\frac{g}{\rho} \left( \frac{d\rho}{dp} \frac{\partial p_0}{\partial z} - \frac{\partial \rho_0}{\partial z} \right) z \quad (5)$$

152 Following the same approach using the hydrostatic equation (4) and adiabatic equation (6)

$$153 \quad d \ln p = \frac{d \ln p}{\gamma}, \gamma = c_p/c_v \quad (6) \text{ yields}$$

$$154 \quad \frac{d^2s}{dt^2} = -\frac{g}{\rho} \left( \frac{\rho}{\gamma p_0} \frac{\partial p_0}{\partial z} - \frac{\partial \rho_0}{\partial z} \right) z = g \left( \frac{\partial \ln \rho_0}{\partial z} - \frac{1}{\gamma} \frac{\partial \ln p_0}{\partial z} \right) z \quad (7)$$

156 For the ideal gas law of  $p = \rho RT$ , the natural logarithm is taken for altitude,  $z$  on both sides,  
 157 yielding

$$158 \quad \frac{\partial \ln \rho}{\partial z} = \frac{\partial \ln p}{\partial z} - \frac{\partial \ln T}{\partial z} \quad (8)$$

159 Then after, the potential temperature ( $\theta$ ) of the parcel is calculated as follows based on the  
 160 equation (2):

$$161 \quad \frac{\partial \ln \theta}{\partial z} = \frac{\partial \ln T}{\partial z} - \frac{R}{c_p} \frac{\partial \ln p}{\partial z} = \frac{1}{T} \left( \frac{\partial T}{\partial z} + \frac{g}{c_p} \right) = \left( 1 - \frac{R}{c_p} \right) \frac{\partial \ln p}{\partial z} - \frac{\partial \ln \rho}{\partial z} \quad (9) \text{ to derive the}$$

162 Parcels acceleration based on equations (7) to become:

$$163 \quad \frac{d^2s}{dt^2} = -g \frac{\partial \ln \theta_0}{\partial z} z \sin a = -g \frac{\partial \ln \theta_0}{\partial z} ds \cdot \sin^2 a \quad (10)$$

164 Whereas by introducing the frequency, N, with  $N^2 = g \frac{\partial \ln \theta_0}{\partial z}$

165 The Brunt-Vaisala frequency,  $N^2$  is calculated based on the following mathematical formulation  
 166 used to characterize atmospheric stability/instability.

$$167 \quad N^2(z) = \frac{g(z)}{T_0(z)} \left( \frac{\partial T_0(z)}{\partial z} + \Gamma_d \right) \quad (11)$$

168 Where  $g$  is the acceleration due to gravity,  $N$  is the Vaisala frequency,  $T_0$  is the background  
 169 temperature (estimated based on third-order polynomial fitting),  $\Gamma_d = g/c_p$  is the adiabatic lapse  
 170 rate, and  $c_p = 1004 J K^{-1} kg^{-1}$  is the specific heat capacity of the atmosphere at constant  
 171 pressure. When the Väisälä frequency,  $N^2$ , is positive, the atmosphere is stable whereas is  
 172 negative, the atmosphere is unstable. When the Brunt-Väisälä frequency,  $N^2$ , is positive the  
 173 atmosphere is stable, whereas a negative  $N^2$  indicates atmospheric instability. In this regard, the  
 174 atmospheric lapse rate,  $\Gamma = -\frac{\partial T}{\partial z}$ , is larger than the adiabatic lapse rate,  $g/c_p \approx 9.5 K km^{-1}$ . A  
 175 third-order least squares polynomial fit was applied to the SABER-observed temperature ( $T$ )  
 176 profile to determine the background temperature ( $T_0$ ), following the method outlined by Leblanc  
 177 and Hauchecorne (1997). subsequently the perturbed temperature ( $T_p$ ) is computed by  
 178 subtracting the background temperature from the observed temperature data ( $T$ ) in equation (12),

$$179 \quad T_p' = T - T_0 \quad (12)$$

180 After estimating the perturbed temperature ( $T_p$ ), a high-pass band filter is applied. This filter  
 181 removes low-frequency components associated with planetary and tidal waves, retaining the  
 182 high-frequency components related to short-period gravity waves (John and Kumar, 2012). This  
 183 process isolates the influence of gravity waves, enabling accurate calculation of their potential  
 184 energy. The high-pass filter operates within known frequency ranges, typically below a one-hour  
 185 period.

$$186 \quad E_p(z) = \frac{1}{2} \left( \frac{g(z)}{N(z)} \right)^2 \left( \frac{T_p'(z)}{T_0(z)} \right)^2 \quad (13)$$

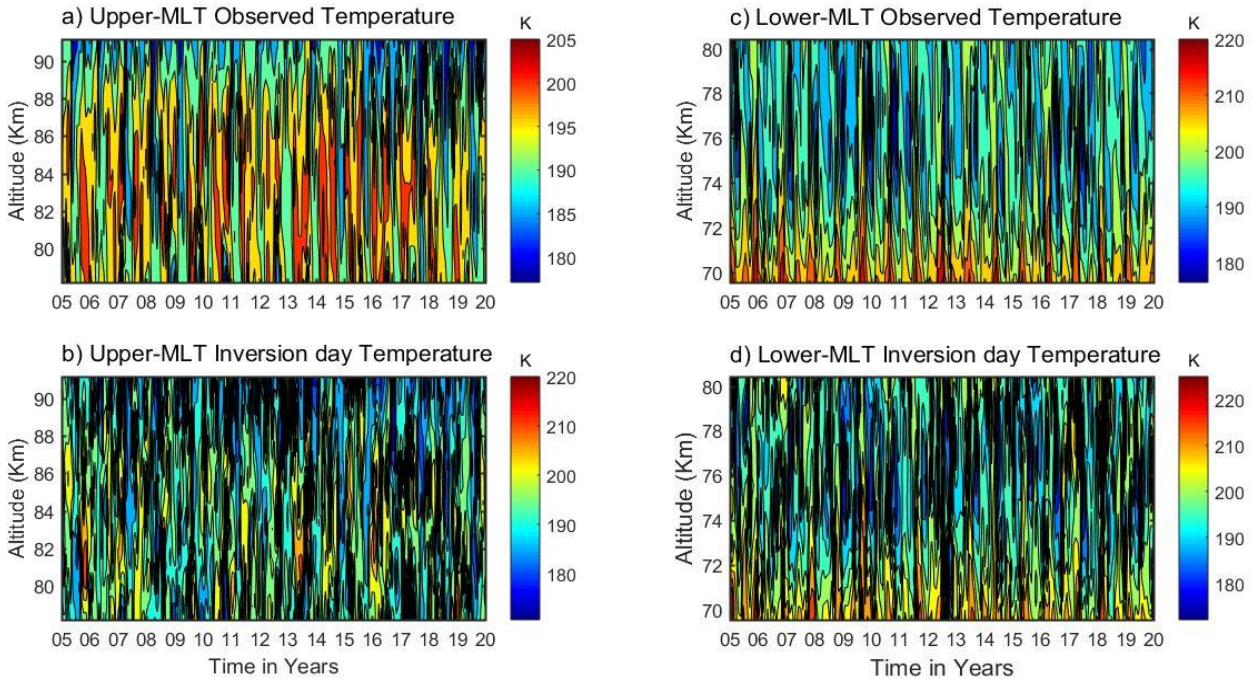
187 The potential energy of the waves, a function of altitude ( $z$ ), is used to determine the  
 188 impact of atmospheric gravity waves on atmospheric inversions.

### 189 **3. Results and discussion**

#### 190 **3.1 Identification and Characteristics of the Lower and Upper MLT Inversion**

191 The SABER temperature profiles, covering altitudes of 60–100 km during 2005 to 2020, are  
 192 shown in the contour plots of Figure 3. Figures 3(a and b) depict the upper MLT (mesospheric  
 193 and lower thermospheric), while Figures 3(c and d) show the lower MLT region. The first

194 horizontal panel of Figures 3(a) and 3(c) show observed temperatures ranging from  
 195 approximately 180–220 K, before accounting for inversion layers. Whereas, the second  
 196 horizontal panel of Figures 3(b) and 3(d) show inversion day temperatures, ranging from 180–  
 197 225 K.

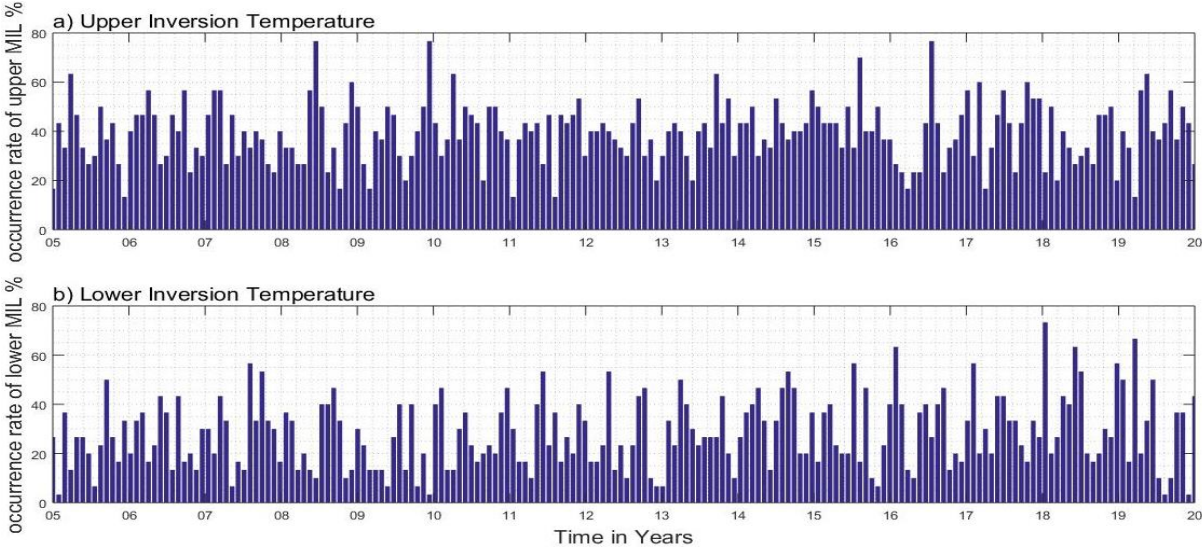


198 **Figure 3.** The upper and lower mesosphere observed temperatures in the first horizontal panel at  
 199 (a and c) with their inversions in the second horizontal panel at (b and d).  
 200

201 The upper left panel of Figure 3(a) shows the observed temperature in the upper mesosphere  
 202 ranges from approximately 180–205 K at altitudes of around 80–90 km. The upper right panel of  
 203 Figure 3(c) shows the lower mesosphere, with temperatures ranging from about 180–220 K at  
 204 altitudes of approximately 70–80 km. In contrast, the lower left panel of Figure 3(b) shows an  
 205 upper-mesosphere inversion day temperature ranges from 180–220 K at an altitude of  
 206 approximately 80–90 km. The lower right panel of Figure 3(d) shows a lower-mesosphere  
 207 inversion day temperature, with the temperature ranging from 180–225 K at an altitude of  
 208 approximately 70–80 km. These inversion day temperatures are higher than those shown in  
 209 Figures 3(a) and 3(c). This indicates that maximum temperatures occur on inversion days in both  
 210 the upper and lower MLT regions relative to the observed temperature. These inversion day  
 211 temperatures in Figures 3(b) and 3(d) suggest a temperature gradient shifting from negative to  
 212 positive. This could be due to factors such as atmospheric waves (planetary, tidal and particularly  
 213 gravity waves), chemical reactions, or solar radiation. Our temperature observations for the  
 214 lower MLT region on an inversion day, within the altitudinal range of 70–80 km, align with



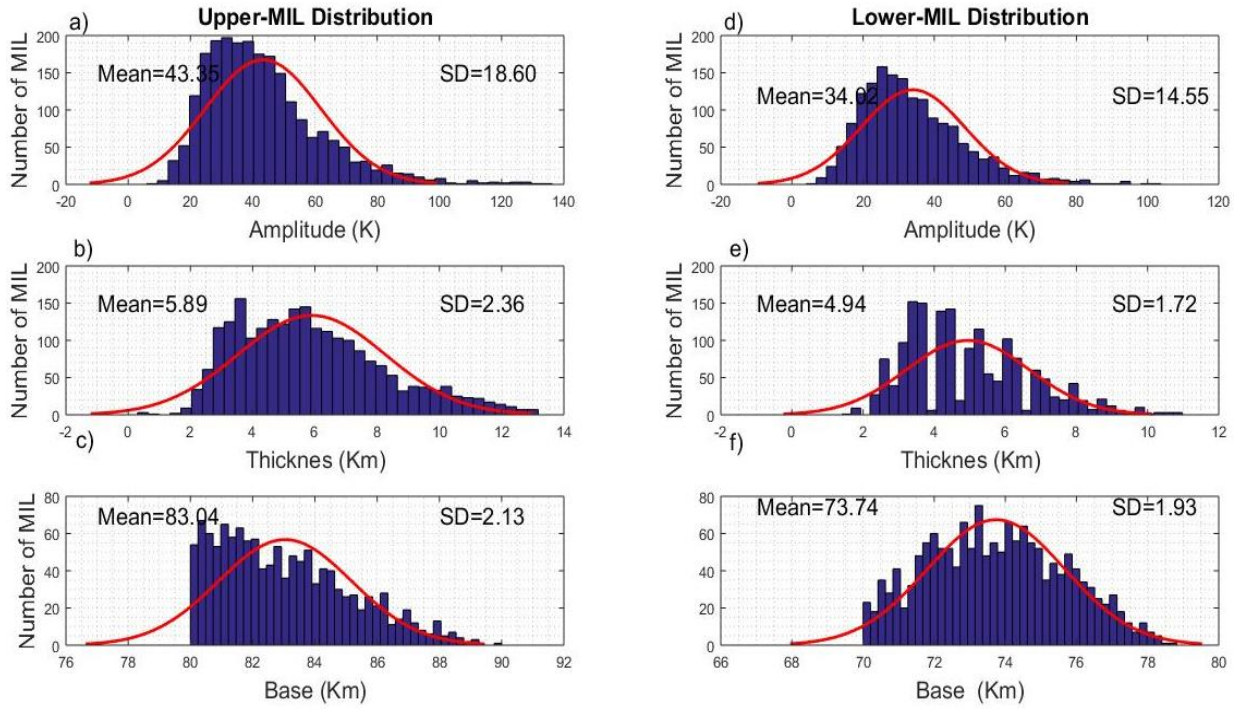
215 those reported by Sivakumar et al. (2001), who identified inversion day temperature variability  
 216 in the altitudinal range of 73–79 km. Additionally, Sivakandan et al. (2014) examined  
 217 mesospheric inversions in the 60–105 km altitude range over low-latitude regions; their findings  
 218 closely match our results.



219 **Figure 4.** The frequency occurrence rate (percentage) of the (a) upper and (b) lower inversion  
 220 temperatures during 2005-2020 over low latitudes.  
 221

222 Figure 4 shows the frequency occurrence rate of mesospheric inversion layers (MILs) in  
 223 histograms. Figure 4(a) shows the occurrence rate for upper MILs, while Figure 4(b) shows the  
 224 rate for lower MILs. The mean frequency occurrence rate of upper inversions is approximately  
 225 below 40%. Peak rates range from 60% to 78%, notably in the years 2008, 2010, and mid-2016.  
 226 In contrast, the mean occurrence rate for lower inversions of Figure 4(b) is approximately below  
 227 20%. The overall occurrence rate for upper inversions is relatively higher compared to lower  
 228 inversions; this may be related to atmospheric wave activities, mainly gravity waves.  
 229 Hauchecorne et al. (1987) and France et al. (2015) discuss the effects of gravity waves on  
 230 inversion variability in the upper and lower mesosphere. Regarding these findings, Figure 5  
 231 examines the characteristics of the inversion day temperature variability, based on their  
 232 amplitude and thickness. It focuses on base height, amplitude, and thickness, before examining  
 233 the effects of gravity waves on an inversion. Histograms display the frequency distribution of  
 234 amplitude, thickness, and base height for MLT temperature variability on inversion days, with  
 235 best-fit Gaussian distribution curves shown in red. The observed distributions align with  
 236 Gaussian curves, indicating that the number of mesospheric inversion layers (MILs) follows a  
 237 normal distribution. This suggests that the attributes are real-valued random variables. The left  
 238 column has three rows showing histograms of (a) amplitude, (b) thickness, and (c) base height of

239 the inversion day temperature variability for the upper MLT. These histograms also include the  
 240 statistical metrics mean and standard deviations (SD) with their values. The corresponding right  
 241 column has three rows representing (d) amplitude, (e) thickness, and (f) base height of the  
 242 inversion day temperature variability for the lower MLT region.



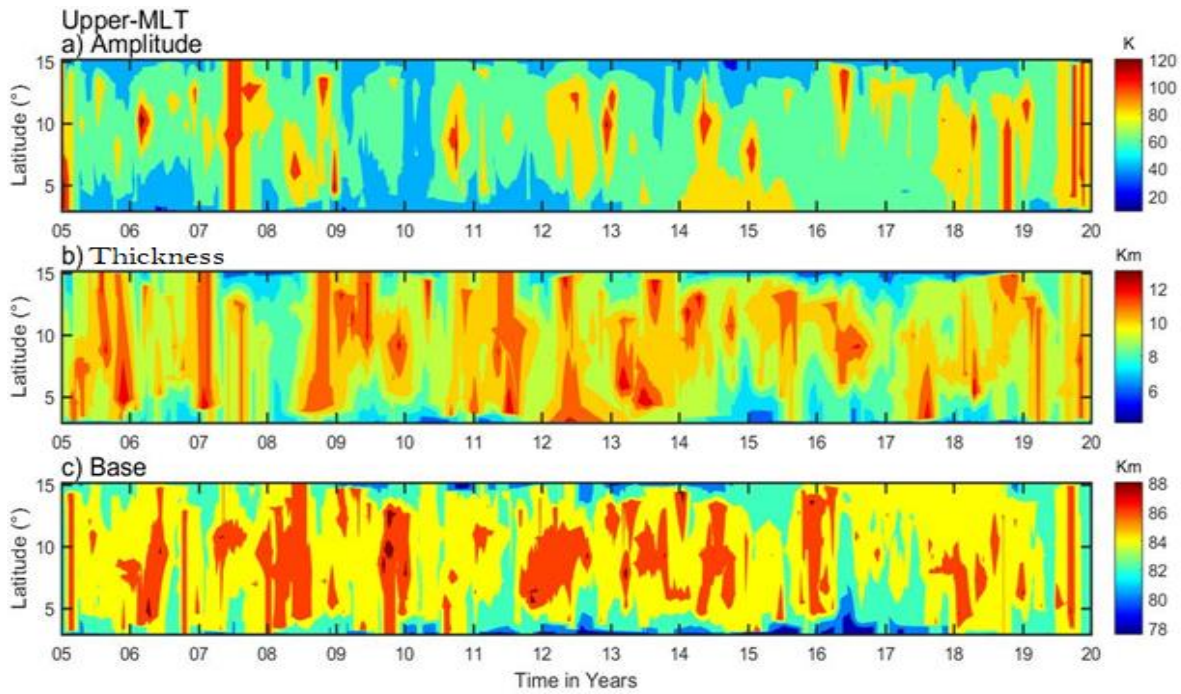
243  
 244 **Figure 5.** The histograms depict the occurrence of MLT inversion day temperature variability.  
 245 The first vertical panel shows the distribution of (a) amplitude, (b) thickness, and (c) base height  
 246 for the upper inversion day. The second vertical panel presents the corresponding distribution for  
 247 the lower inversion, including (d) amplitude, (e) thickness, and (f) base height.

248 The amplitude of upper inversion day temperature variability in Figure 5(a) ranges between 20  
 249 and 80 K, with a peak value of 38 K. This follows a Gaussian distribution with a large standard  
 250 deviation (SD) of 18.6, which indicates high inversions. The thickness of the inversion layer for  
 251 upper MILs, shown in Figure 5(b), ranges from 3 to 9 K, and their most probable value is 5.5 K,  
 252 with a low SD of 2.3. The base height of the upper MIL in Figure 5(c) spans from ~80 to 90 km,  
 253 with a peak value of around 83 km. This indicates a large number of upper MLT inversions, with  
 254 an SD of 2.13. The highest number of upper inversions between 2005 and 2020 is observed at 82  
 255 km. This may be attributed to gravity wave breaking and energy dissipation, influenced by waves  
 256 generated from lower atmospheric regions and solar flux impacts.

257 The lower inversion amplitude, shown in Figure 5(d), ranges from 10 to 60 K, with a peak value  
258 of 25 K and a standard deviation (SD) of 14.5. The inversion thickness, as illustrated in Figure  
259 5(e), spans 3 to 8 km, with the most likely value at 3.8 km and a low SD of 1.72. The base height  
260 of the lower inversion in Figure 5(f) ranges from 70 to 80 km, with a peak value of around 74 km  
261 and a lower SD of 1.93. Previous investigations by Sivakandan et al. (2014) from the Indian  
262 sector reported amplitudes ranging from 14–39 K in 2002 and 15–42 K in 2008. The thicknesses  
263 ranged between 2.7–7.5 km in 2002 and 2.8–7.3 km in 2008, under the influence of solar flux.  
264 These findings align well with the present study, indicating no significant variation in  
265 characterizing mesospheric inversion based on amplitude and thickness in the low-latitude region  
266 within the altitude range of 60 to 90 km.

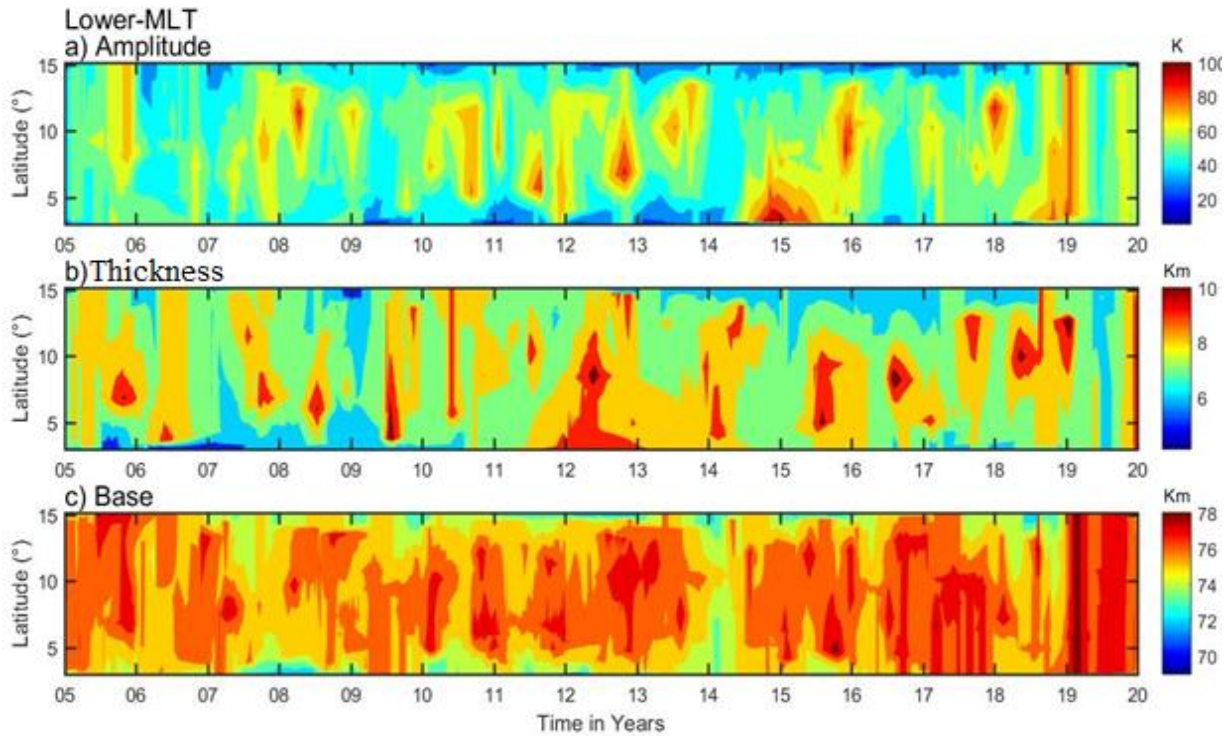
### 267 **3.2 Latitudinal Variations of Mesospheric Inversion Layers (MILs)**

268 This section examines the spatiotemporal variability of upper and lower mesosphere inversion  
269 phenomena. Contour plots of time vs. latitude in Figures 6 and 7 show the variability of upper  
270 and lower MLT inversion amplitude, thickness, and base height over the low-latitude band (3-  
271 15°) during 2005–2020. The upper inversion is observed around 80–90 km; their maximum  
272 amplitude, in the range of 90–120 K, occurs over latitude bands (5-12°) during 2005, 2007, mid-  
273 2011, 2013, 2015, 2016, mid-2019, and 2020 (Figure 6(a)). The second horizontal panel of  
274 Figure 6(b) shows the thickness, with a maximum range of ~(8–12 km) across the entire  
275 latitudinal region (3-15° N). The third horizontal panel of Figure 6(c) shows the base height, with  
276 relative maximum values around ~(84-88 km), in the latitudinal range between (4° & 14°) N  
277 during 2006, 2008, 2010, 2012, 2016, and 2018.



278  
 279 **Figure 6.** The daily upper inversions (~80-90 km) of (a) amplitude, (b) thickness, and (c) base  
 280 height during 2005-2020 over latitudinal variation.

281 Whereas, the contour plots of Figure 7(a, b, and c) show the lower inversions (MILs) amplitude,  
 282 thickness, and base height over an altitudinal range of ~70-80 km, respectively. Overall,  
 283 amplitude values ranged from approximately 30-60 K across all latitudinal bands, reaching peak  
 284 ranges of about 80-100 K during 2013, 2015, 2016, and 2019 in various latitude regions between  
 285 5° and 14°N. The second horizontal panel of Figure 7(b) indicates thickness values ranging from  
 286 5-7 km across the entire latitude band, with maximum thickness reaching 8-10 km in 2012-2013,  
 287 2016, and 2019. The final horizontal panel of Figure 7(c) shows a base height of 76–80 km  
 288 across most latitudes and periods, with notable exceptions in 2008, 2014, and mid-2018, where  
 289 the base height reaches its maximum.



290

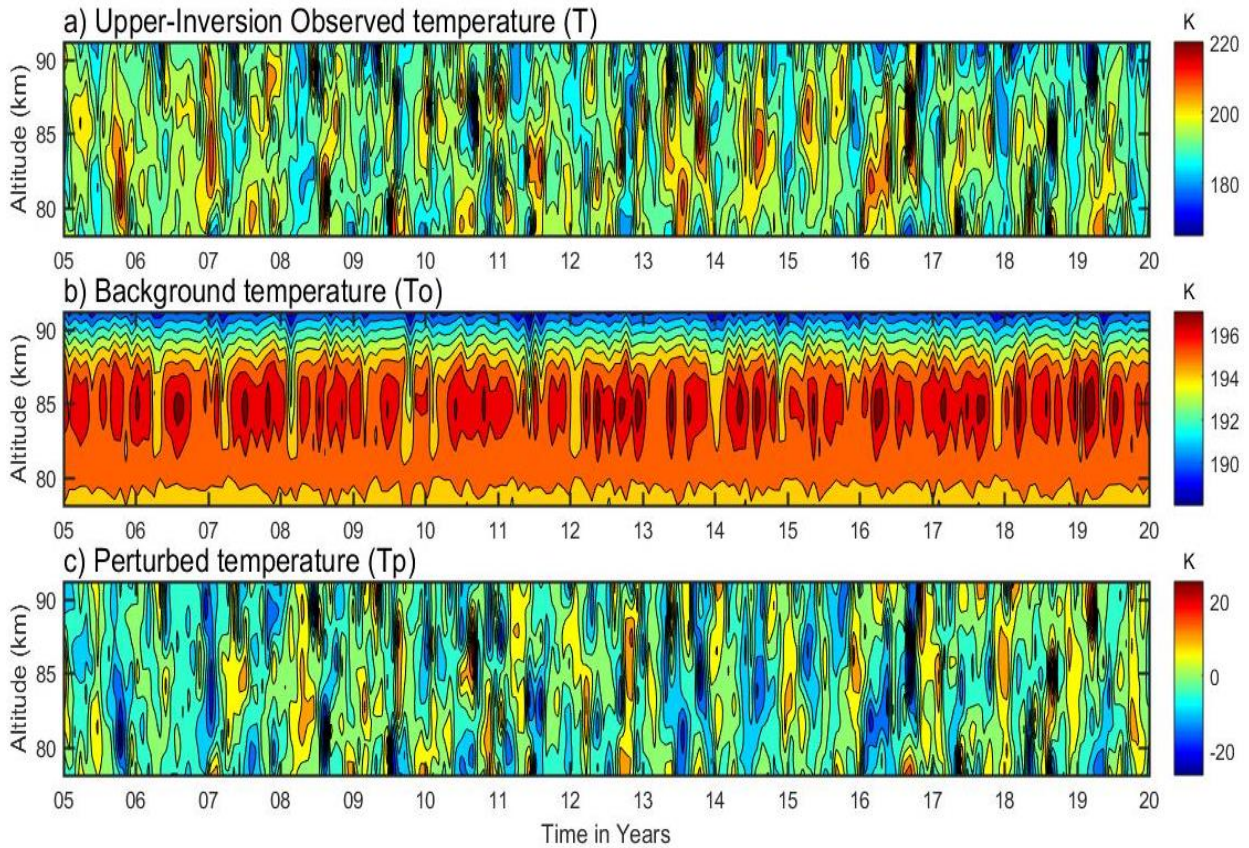
291 **Figure 7.** Same as Figure 5, but for the lower mesosphere inversions (~70- 80 km).

292 The higher amplitude and thickness are demonstrated in the upper and lower inversion to exhibit  
 293 a suggested highly dynamic phenomenon. These valuable investigations are confirmed in Gan et  
 294 al. (2012) based on SABER satellite observation at low latitudes.

### 295 **3.3 Analysis of Perturbed Temperature Variations in the MLT Region**

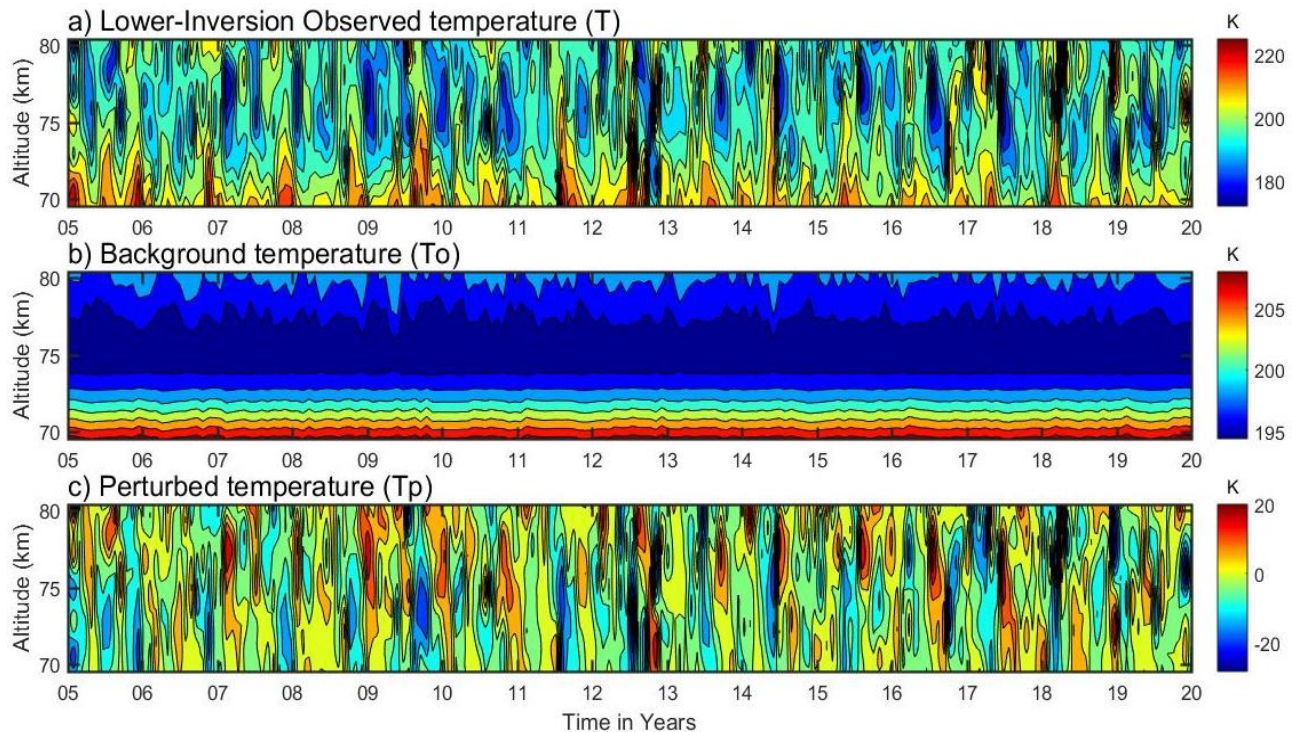
296 The perturbed temperature ( $T_p$ ) of the upper and lower MLT inversions in Figure 8 and 9 can  
 297 further be used to calculate their derived potential energy of gravity waves and the Brunt-Väisälä  
 298 frequencies ( $N^2$ ). First, the upper inversion profiles are identified in the MLT region during the  
 299 entire observational period of 2005-2020, as displayed in the contour plot of Figure 8(a). Based  
 300 on the observed temperature, which ranges from ~170 to 220 K with minimal variability, the  
 301 background temperature is estimated. A 3<sup>rd</sup>-order polynomial fit is applied to calculate the  
 302 background temperature ( $T_0$ ), as shown in the contour plot of Figure 8(b). This maximum  
 303 background temperature exhibits a periodic variability over an altitude of around ~82-87 km,  
 304 ranging from ~195-197 K. The perturbed temperature profiles ( $T_p$ ), determined by subtracting  
 305 the background temperature profiles ( $T_0$ ) from the observed inversion temperature ( $T$ ), are in the  
 306 range between -25 and +25 K, as illustrated in Figure 8(c).

307 The lower-MLT region perturbed temperature ( $T_p$ ) is calculated using the same approach as the  
 308 upper-MLT perturbed temperature, using the observed and background temperatures. Their  
 309 corresponding contours are displayed in Figure 9(a-c).



310 **Figure 8.** The upper mesosphere temperatures in the vertical panel are: (a) inversion day  
 311 observed temperature; (b) background temperature; and (c) perturbed temperature in the upper  
 312 mesosphere region.  
 313

314 In Figure 9(a), the observed temperature of the lower inversion ranges from  $\sim 170$ - $220$  K.  
 315 Whereas, the background temperature of the lower inversion ranges from  $\sim 195$ - $210$  K, with  
 316 maximum values of  $\sim 200$ - $210$  K at a height of  $\sim 70$ - $72$  Km, as shown in Figure 9(b). The derived  
 317 lower-MLT perturbed temperature from the observed and background temperature is presented  
 318 in Figure 9(c) and ranges from  $-25$  to  $20$  K. Notably, the upper-MLT perturbed temperature is at  
 319 its maximum compared to the lower-MLT region, possibly due to a highly dynamic  
 320 phenomenon.



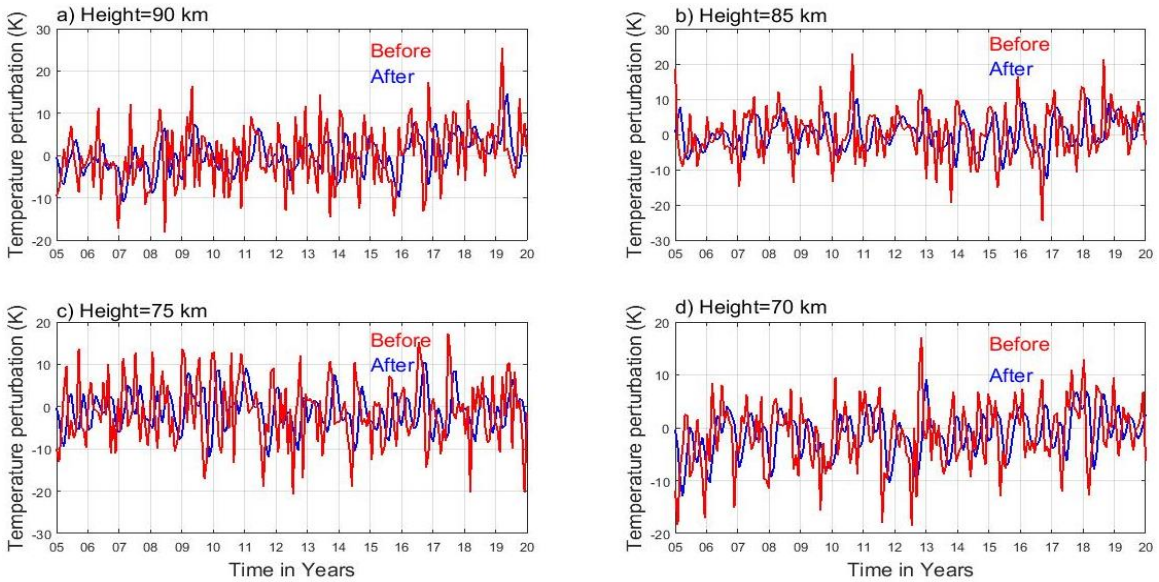
321  
322 **Figure 9.** Same as Figure 7, but for the lower mesosphere atmospheric region.

### 323 **3.4 Effects of Gravity Waves on Mesosphere Inversions and Associated Instability**

324 Atmospheric gravity waves form when air parcels are oscillated due to the restoring force of  
 325 gravity after being transported vertically. Several factors contribute to these waves, including  
 326 airflow over mountains, convection, and wind shear. As the waves propagate vertically, they  
 327 break and dissipate, releasing energy and momentum into the surrounding atmosphere, which  
 328 contributes to the formation of inversion layers. The gravity wave contribution is quantified by  
 329 calculating the potential energy and assessing their impact on MLT instability through the Brunt-  
 330 Väisälä frequency ( $N^2$ ), derived from perturbed temperature ( $T_p'$ ) data spanning 2005–2020.  
 331 Several authors (Tsuda et al., 2000; Wang and Geller, 2003; Liu et al., 2014; Thurairajah et al.,  
 332 2014) suggest that gravity wave activity is represented by potential energy. Further investigation  
 333 is required, focusing on altitudes of 90, 85, 75, and 70 km, to evaluate the impacts of gravity  
 334 waves on an inversion by applying a high-pass filter with a one-hour interval to the  $T_p'$  data (see  
 335 Figure 10 a-d). The high-pass filter attenuates low-frequency components, removing the effects  
 336 of long-period wave oscillations, such as tidal and planetary wave contributions. This effectively  
 337 isolates the gravity waves (Gw), allowing a clearer focus on their impact on MLT inversions, as  
 338 appeared in Figure 10.

339 The blue curve in Figure 10(a to d) appears smoother after applying the high-pass filter to the  
 340 perturbed temperature. However, the filter removes the peaks of low-frequency variations,

341 resulting in retained perturbed temperature values that appear more uniform, creating a smooth  
 342 plateau effect. In the upper mesosphere (90 and 85 km), the filter reduces the amplitude of wave  
 343 oscillations from approximately  $\pm 20$  K to  $\pm 10$  K, as shown from the blue curve in Figure 10a and  
 344 b, compared to the red curve. Similarly, in the lower mesosphere (75 and 70 km) at (Figure 10 c  
 345 & d) the amplitude decreases from  $\sim(-20$  to  $20$  K) to  $\sim(-8$  to  $8$  K) by filtering out higher  
 346 amplitudes.



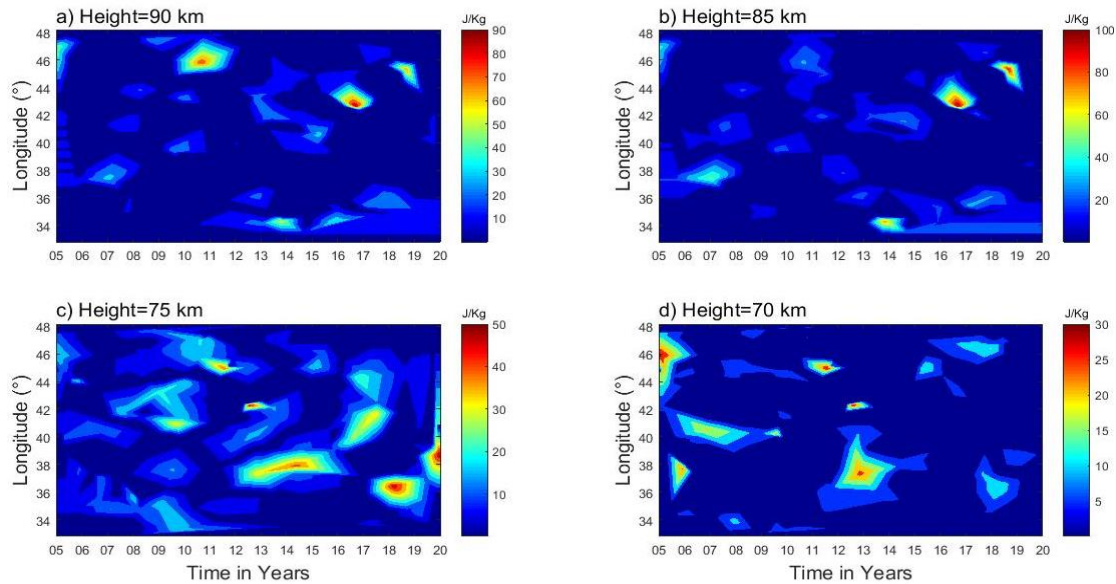
347 **Figure 10.** Perturbed temperature profiles before (red color) and after (blue color) applying the  
 348 high-pass filter for the upper (90 and 85 km) and lower (75 and 70 km) regions.  
 349

350 In the MLT atmospheric region, gravity wave breaking typically dissipates their potential and  
 351 kinetic energy, leading to increased turbulence and mixing. As illustrated, gravity wave  
 352 propagation and dissipation are major forces in the MLT (Lindzen, 1981; Holton, 1983),  
 353 influencing the middle and upper atmospheric regions. This has a substantial impact on the  
 354 overall dynamics as well as the MLT's thermal structure, particularly the increase in temperature  
 355 variability with elevation, known as inversion. Holton et al. (2003) and Holton and Hakim  
 356 (2013) has demonstrated an interaction between the potential energy of gravity waves and  
 357 inversions. This notable upper and lower inversions are observed in Figure 4 during the period  
 358 2005–2020 over the low-latitude regions. During this period, particularly for the upper-MLT  
 359 region above 80 km altitude, high-resolution SABER satellite temperature data revealed the  
 360 presence of a strong mesospheric inversion layer (MIL), with peak occurrence rates ranging  
 361 between 60% and 78%, especially during 2010, 2014, 2016/17, and 2018/19, Figure 4a.  
 362 Correspondingly, at the same time and in the same region (upper-MLT), at altitudes of 85 and 90  
 363 km, there is a noticeable increase in gravity wave potential energy ( $E_p$ ), shown in Figure 11. The



364 maximum potential energy (PE) for the upper-MLT region corresponds to the breaking or  
365 dissipation of gravity waves as they propagate upward. This spike in potential energy coincides  
366 with the occurrence of the inversion layer, suggesting that the breaking or dissipation of gravity  
367 waves releases energy into the atmosphere, contributing to localize heating in the mesosphere  
368 and leading to the formation of the inversion. The sudden transfer of momentum and energy  
369 from the breaking GWs to the surrounding atmosphere disrupts the thermal structure, causing the  
370 temperature inversion. In this case, the temporal and spatial coincidence between the peak in  
371 gravity wave potential energy and the formation of the inversion demonstrates a clear physical  
372 connection. The energy released from breaking GWs plays a direct role in the creation of the  
373 inversion layer, as shown in Figures 4 and 11. Similarly, the statistical distributions of upper-  
374 MLT inversions in Figure 5(a) show maximum amplitudes, which correspond to the maximum  
375 potential energy of gravity waves in Figure 11(a & b). This provides a straightforward  
376 demonstration of how gravity wave dynamics-specifically, the dissipation of their potential  
377 energy-are linked to the formation of mesospheric inversion layers (MILs).

378 Figure 11 (a-d) demonstrates that the spatiotemporal variability of gravity wave potential energy,  
379 showing over the upper-MLT at (90 and 85 km) and the lower-MLT at (75 and 70 km). Figure  
380 11(a) of upper-MLT inversions at 90 km shows maximum gravity wave potential energies,  
381 ranging from ~70 to 90 J/kg, over the longitudinal regions of 45-47° E, as well as at 43° E, and  
382 44° E during 2011, 2017, and 2019. In contrast, potential energies being the least in amount  
383 around ~10 to 60 J/kg are present across the entire longitudinal region from 33-48° E. Figure  
384 11(b) of upper-MLT region shows maximum potential energies of ~70 to 100 J/kg over the  
385 longitudinal regions of 34°, 44°, and 46° E during 2014, 2016, and 2018 at 85 km. Its minimum  
386 potential energies between 20 and 70 J/kg appear over the longitude regions from 33-48° E.  
387 Whereas, Figures 11(c) and 11(d) depict the gravity wave potential energy in the lower-MLT  
388 regions at 75 and 70 km, respectively. At 75 km, Figure 11(c) shows a relative maximum  
389 potential energy of 40-50 J/kg over the longitudinal regions of 46°, 42°, 40°, 37°, 36°, and 38° E  
390 during 2011, 2012, 2017, 2013–2015, 2018, and 2020. Similarly, Figure 11(d) illustrates gravity  
391 wave potential energy ranging from 2-30 J/kg at 70 km across the longitudinal region of 33-48°  
392 E. Their maximum potential energy of 25-30 J/kg is observed in certain longitudinal regions over  
393 time.

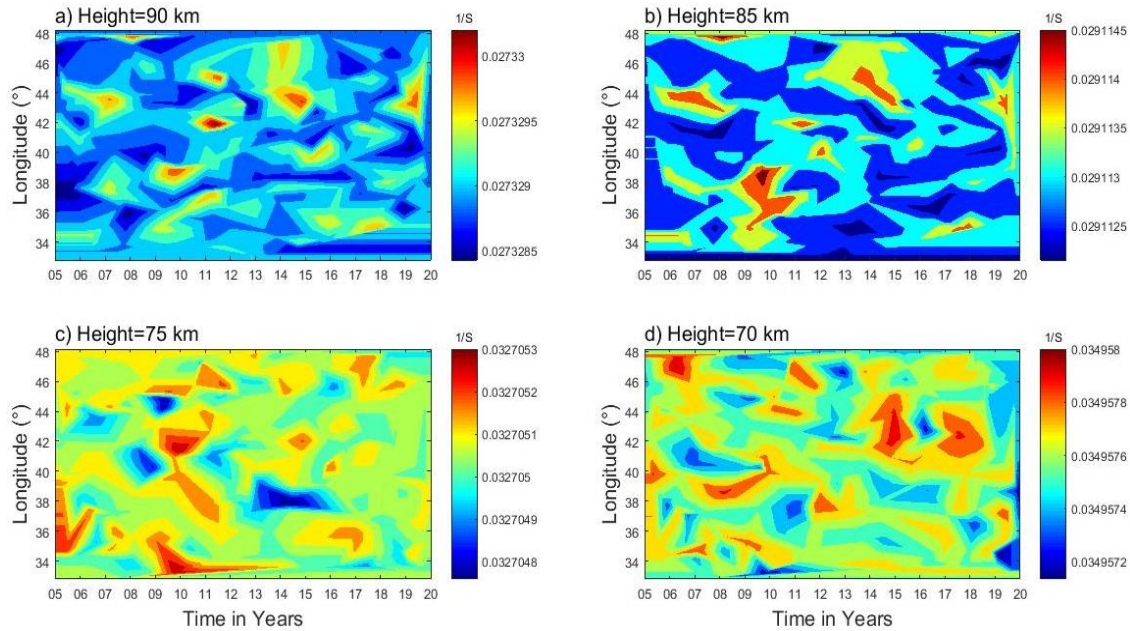


394

395 **Figure 11.** Gravity wave potential energy for the upper (90 and 85 km) and lower (75 and 70  
 396 km) MLT regions.

397 The role of gravity waves in the MLT region instability, the Brunt-Väisälä frequency is  
 398 analysed. Contour plots in Figure 12 (a-d) show the spatiotemporal variability of the Brunt-  
 399 Väisälä frequency, with Figures 12(a and b) representing the upper-MLT (90 and 85 km) and  
 400 Figures 12(c and d) representing the lower-MLT (75 and 70 km). The Brunt-Väisälä frequency  
 401 ( $N^2$ ) shows that the upper MLT region is more unstable ( $\sim 0.027$  at 90 km and  $\sim 0.029$  at 85 km)  
 402 relative to that of the lower MLT region ( $\sim 0.033$  at 75 km and  $\sim 0.035$  at 70 km). These different  
 403 values of brunt-vaiala frequency under a quencyquence of gravity waves generation in different  
 404 sizes, with smaller waves being the main drivers of instability and turbulence in the MLT  
 405 (mesosphere and lower thermosphere) region (Liu and Meriwether, 2004; Szewczyk et al.,  
 406 2013). Hauchecorne et al. (1987) proposed a model where a series of breaking gravity waves  
 407 leads to the formation of MILs through the gradual accumulation of heat, which contributes to  
 408 instability. Conducting mesospheric inversion layer (MIL) phenomena is crucial for  
 409 understanding MLT atmospheric dynamics, especially when it comes to stability and energy  
 410 transfer.

411



412

413 Figure 12. Brunt-Väisälä frequency ( $N^2$ ) variability for the upper (90 and 85 km) and lower (75  
414 and 70 km) MLT regions.

#### 415 **4 Summary**

416 In this article, 16 years of SABER MLT temperature profiles are utilized to investigate the MIL  
417 phenomenon and its causative mechanism through gravity wave potential energy ( $P_E$ ) and  
418 instability criteria of Brunt-Väisälä frequency ( $N^2$ ) over low-latitudes. The following conclusions  
419 are drawn from the observations in this article:

- 420 ✓ The upper mesosphere inversion frequency occurs more often than the lower mesosphere  
421 inversion.
- 422 ✓ Analysis of the MIL characteristic features reveals the most probable values for the upper  
423 inversion: amplitude of 38 k, thickness of 5.5 km, and base height of 78 km. The lower  
424 inversion has an amplitude of 25 K, a thickness of 3.8 km, and a base height of 73 km.
- 425 ✓ The upper mesosphere region has higher gravity wave potential energy compared to the  
426 lower mesosphere region.
- 427 ✓ The high potential energy in the upper mesosphere region is likely due to the deposition of  
428 energy and momentum by gravity wave breaking. This could influence the dynamics of the  
429 inversion phenomenon.
- 430 ✓ The Brunt-Väisälä frequency ( $N^2$ ) indicates that the upper mesosphere region is less stable  
431 than the lower mesosphere region. This lower stability contributes to the high potential  
432 energy in the upper mesosphere, which leads to larger inversion phenomena.

433 ✓ Atmospheric processes vary significantly from region to region, with altitude, and over time.

434 **Data availability.** The SABER data are freely available via the link at [http://saber.gats-inc.com/](http://saber.gats-inc.com/index.php)  
435 [index.php](http://saber.gats-inc.com/index.php).

436 **Author contribution.** Chalachew Lingerew: data curation, investigation, software,  
437 visualization, writing the original draft, and writing review. U. Jaya Prakash Raju; supervision,  
438 and editing.

439 **Competing interest.** The authors declare that they have no conflict of interest relevant to this  
440 study.

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443 <http://saber.gats-inc.com/index.php>.

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