



27 response of water vapour to the solar variability. ~~The solar cycle response is reduced in the late~~  
28 ~~solar cycle due to a smaller amplitude of Lyman- $\alpha$  variability in the second period.~~ -Applying  
29 the geometry of satellite observations, we find a missing response when averaging over altitudes  
30 of 80 to 85 km, where H<sub>2</sub>O has a positive and a negative response (depending on altitude) which  
31 largely cancel out. One main finding is that during NLCs the solar cycle response of H<sub>2</sub>O  
32 strongly depends on altitude. ~~A negative correlation between H<sub>2</sub>O and Lyman- $\alpha$  is found in the~~  
33 ~~NLC sublimation zone below an altitude of about 83 km, but a positive response is present at~~  
34 ~~the altitudes above 83 km where NLCs form.~~

35

## 36 **1. Introduction**

37

38 The 11-year solar cycle significantly influences the upper atmosphere's temperature and water  
39 vapour (H<sub>2</sub>O) concentration. H<sub>2</sub>O is one of the essential minor constituents in the mesosphere  
40 as it is the primary source of chemically active hydrogen radicals, influencing the chemistry of  
41 all other chemically active minor constituents (Brasseur and Solomon, 2005, Hartogh et al.,  
42 2010). H<sub>2</sub>O concentration plays an essential role in the noctilucent cloud's (NLC) formation.  
43 NLCs are located at about 83 km altitude, consist of water ice particles, and owe their existence  
44 to the cold summer mesopause region (~130K) at mid and high latitudes. NLCs, also called  
45 polar mesospheric clouds, are formed in an environment where small changes in background  
46 H<sub>2</sub>O and temperature can lead to significant changes in NLC properties (e.g., Thomas, 1996;  
47 DeLand et al., 2006; Shettle et al., 2009, Lübken et al., 2009).

48 In comparison to the lower atmosphere, little is known about the upper mesosphere/lower  
49 thermosphere (MLT, 75-110 km) due to a lack of observations at these altitudes. NLCs have  
50 been proposed as indicators of trends in background temperature and H<sub>2</sub>O concentrations  
51 (Thomas & Olivero, 2001). Studying NLC properties provide insight in phenomena occurring  
52 at the altitude of NLC. The 11-year solar cycle has been considered to cause quasi decadal

53 oscillation observed in NLCs (DeLand et al., 2003). NLCs are predicted to decrease during  
54 solar maximum due to increased heating and photolysis of H<sub>2</sub>O (Garcia, 1989). However, some  
55 recent studies strongly suggest that the response of NLCs to the solar cycle has been absent  
56 from 2002 to the present (Fiedler et al., 2011; DeLand & Thomas, 2015; Hervig et al., 2016;  
57 Siskind et al., 2013). Hervig et al. (2019), using satellite observations, found that NLC had a  
58 clear anti-correlation with the solar cycle before 2002, and that response has been absent in  
59 recent years. The leading cause of this absence appears to be the suppression of the solar cycle  
60 response of H<sub>2</sub>O. Lyman- $\alpha$  (Ly $\alpha$ ) radiation is the primary cause of H<sub>2</sub>O photolysis and varies  
61 by a factor of two between solar minimum and maximum (Woods et al., 2000). Understanding  
62 the effects of the solar cycle on H<sub>2</sub>O is more complicated at NLC altitudes because of the  
63 interaction between NLCs and background H<sub>2</sub>O.

64 NLC growth leads to dehydration at higher altitudes (83-89 km) as ice particles are formed by  
65 consuming background H<sub>2</sub>O, and sublimation of ice particles leads to hydration at lower  
66 altitudes as H<sub>2</sub>O is released here (about 78-83 km) (Lübken et al., 2009, Hervig et al., 2003).  
67 Investigating the effects of NLC on the background H<sub>2</sub>O requires an estimate of the H<sub>2</sub>O profile  
68 without NLCs. Investigations using satellite observations are limited due to uncertainty in the  
69 inferred background H<sub>2</sub>O without NLC and vertical resolutions in the order of a few 100~m.  
70 Therefore, using satellite observations to study H<sub>2</sub>O at NLC altitudes could yield misleading  
71 results due to biases in the estimated H<sub>2</sub>O profiles without NLC (Hervig et al., 2015). Hervig  
72 et al. (2015) suggest that in future studies, one approach to investigate the effects of NLC on  
73 H<sub>2</sub>O would be to use a detailed microphysical NLC model. Therefore, for this study, NLC  
74 ~~model~~-simulations are performed with and without microphysics using the same background  
75 conditions, resulting in a H<sub>2</sub>O profile with and without NLC. This allows us to investigate how  
76 NLC formation changes the H<sub>2</sub>O background profile in detail.

77 We compare the model result to satellite observations published by Hervig et al. (2019) to  
78 investigate the mechanism behind the solar cycle response of NLC and H<sub>2</sub>O. We also focus on

79 the missing solar cycle response of H<sub>2</sub>O during recent years. This paper aims to answer a  
80 number of questions: How does the formation of NLCs affect the H<sub>2</sub>O profile and the variation  
81 of water vapour with the solar cycle? How do the solar cycle-induced temperature and  
82 photolysis changes affect the H<sub>2</sub>O response? Why is the response of water vapour to solar cycle  
83 nearly absent in satellite observations after 2005 (Hervig et al., 2019)? Our study is focused on  
84 the core NLC period, i.e., July at 68±5°N. The following section describes the modelling  
85 framework of this study and discusses the various model simulations performed. The third  
86 section discusses the mechanisms behind the solar cycle H<sub>2</sub>O response, such as the separation  
87 of the solar cycle-induced temperature and photolysis effects on H<sub>2</sub>O. Sections four and five  
88 explore the possible reasons behind the missing solar cycle response. Concluding remarks and  
89 a summary are given in the last section.

90

## 91 **2. Model description and numerical experiments**

92

### 93 **2.1. Model**

94

95 The modelling framework used in this study consists mainly of two components: the Leibniz  
96 Institute Middle Atmosphere (LIMA) model and the Mesospheric Ice Microphysics And  
97 tranSport (MIMAS) model (see Fig. 1). LIMA is a non-linear, global, 3D Eulerian grid-point  
98 model reaching from the troposphere to the lower thermosphere, which calculates winds and  
99 temperature and is well described in a number of papers (Berger, 2008; Lübken et al., 2013).  
100 The LIMA model in this study is nudged to reanalysis data NOAA-CIRES (National Oceanic  
101 and Atmospheric Administration-Cooperative Institute for Research in Environmental Sciences  
102 20CR; Compo et al., 2011) ~~in the lower atmosphere up to an altitude of 45 km.~~ The resulting  
103 winds and temperatures in the mesosphere and lower thermosphere (MLT) are then used in

104 MIMAS. The MIMAS model run was performed for all years with background wind conditions  
105 and gravity wave forcing from a representative year (1976).

106 MIMAS is a 3D Lagrangian transport model specifically designed for modelling ice particles  
107 in MLT region (Berger and Lübken, 2015). MIMAS calculates NLC parameters from 10 May  
108 to 31 August, and it is constrained from mid-latitudes to high latitudes (37°-90°N) with a  
109 horizontal grid resolution of 1° in latitude and 3° in longitude and a vertical resolution of 100  
110 m from 77.8 to 94.1 km (163 levels). In this study, the dynamics calculated by LIMA, solar  
111 Ly $\alpha$ , and the initial H<sub>2</sub>O distribution are the input for MIMAS as sketched in Figure 1. Below  
112 the MIMAS lower boundary two~~Two~~ effects determine the mixing ratio of H<sub>2</sub>O in the  
113 mesospherestratosphere: (i) transport of H<sub>2</sub>O from ~~lower altitudes~~ the troposphere and (ii)  
114 oxidation of methane (CH<sub>4</sub>). The oxidation of each CH<sub>4</sub> molecule produces two H<sub>2</sub>O molecules.  
115 Methane is nearly completely converted to H<sub>2</sub>O in the mesosphere by photochemical processes  
116 (e.g., Lübken et al., 2018). MIMAS assumes that ~~the H<sub>2</sub>O mixing ratio transported from lower~~  
117 ~~altitudes~~ transport from the troposphere is constant. The increase in H<sub>2</sub>O is primarily through  
118 (ii) i.e. due to the increase in CH<sub>4</sub> concentration ~~Therefore, H<sub>2</sub>O concentration varies only~~  
119 ~~according to CH<sub>4</sub> concentration~~ (Lübken et al., 2018). Then, mesospheric. Mesospheric H<sub>2</sub>O in  
120 MIMAS is transported by background winds, dispersed by turbulent diffusion, and reduced by  
121 photolysis. Hence, we parametrize H<sub>2</sub>O as a function of CH<sub>4</sub> following by Lübken et al., 2018  
122 (see Section 2). MIMAS makes use of 40 million dust particles, which can act as condensation  
123 nuclei. Dust particles are formed from meteors evaporating in the atmosphere (for more details,  
124 see Berger and von Zahn, 2002; von Zahn and Berger, 2003, Killiani, 2014). These are then  
125 coated with ice in H<sub>2</sub>O supersaturated regions and transported according to three-dimensional  
126 and time-dependent background winds, eddy diffusion, and sedimentation. In MIMAS,  
127 standard microphysical processes such as the Kelvin effect determine the nucleation and growth  
128 of ice particles (Berger & Lübken, 2015; Gadsden & Schröder, 1989). For the comparison with  
129 satellites, we used model run A, which includes CO<sub>2</sub> and CH<sub>4</sub> variations (Lübken et al., 2018;

Fig. 1

130 Lübken et al., 2021). We performed MIMAS model simulations with ice formation turned off  
131 and on respectively to investigate the effects of ice formation on background H<sub>2</sub>O. In both runs,  
132 the background conditions and model inputs are the same. The main outputs of the model are  
133 the microphysical properties of the NLC ice particles, such as radius, backscatter value, and the  
134 number density of the ice/dust particles. More detailed descriptions of the MIMAS model and  
135 its precursors are available in the literature (Berger and von Zahn, 2002; Berger, 2008; Berger  
136 and Lübken, 2011; Lübken et al., 2018; Lübken et al., 2021).

137

## 138 **2.2. Model simulations**

139

140 LIMA and MIMAS use daily Ly $\alpha$  fluxes [taken from the LASP Interactive Solar Irradiance Data](#)  
141 [Center \(LISIRD\)](#) as a proxy for solar activity from 1961 to 2019 ([Machol et al. 2019](#)). Ly $\alpha$  (and  
142 other [spectral](#) bands) variations in LIMA cause atmospheric temperature variations, while Ly $\alpha$   
143 variations in MIMAS cause photolysis of H<sub>2</sub>O. In LIMA, variation of other bands, namely,  
144 Chappius band, Huggins band, Hartley band, Schumann-Runge band, and both Schumann-  
145 Runge continuuums are [taken into account parametrised following the Hamburg Model of the](#)  
146 [Neutral and Ionized Atmosphere \(Schmidt et al., 2006\)](#). [The parametrization schemes are](#)  
147 [discussed in more detail in Berger, 2008 \(see Section 2.2\)](#). [Variations of these bands are](#)  
148 [parametrized based on Ly \$\alpha\$  values](#) according to Lean et al. (1997). Therefore, it is possible to  
149 study the effects of solar cycle on H<sub>2</sub>O due to temperature changes and photolysis separately  
150 by performing model simulations with constant and varying Ly $\alpha$  in MIMAS and LIMA. We  
151 conducted four model runs, as described in Table 1. We also performed LIMA model  
152 simulations with constant CO<sub>2</sub> for runs E, F, and G to filter out their effects on temperature  
153 changes. For these runs we use a constant CH<sub>4</sub> concentration in MIMAS to avoid its influence  
154 on the H<sub>2</sub>O profile.

155 In LIMA, the mixing ratios of CO<sub>2</sub> (28-150 km) vary as function of time (years), while all other  
156 trace gases are kept constant. An increase in CO<sub>2</sub> leads to a decrease in temperature in the  
157 stratosphere mainly due to enhanced cooling by CO<sub>2</sub> (e.g., Roble and Dickinson, 1989; Garcia  
158 et al., 2007; Berger & Lübken, 2011; Marsh et al., 2013; Lübken et al., 2013). At NLC altitudes,  
159 this cooling leads to an altitude decrease of pressure levels, referred to as the ‘shrinking effect’  
160 (Lübken et al., 2009). For LIMA we use the long-term increase of CO<sub>2</sub> concentration according  
161 to observations at Mauna Loa (19°N, 155°W).

162 This study focuses mainly on the recent two solar cycles from 1992 to 2018. Figure 2 shows  
163 the time series of Ly $\alpha$ , CO<sub>2</sub>, and CH<sub>4</sub> for 1992-2018. The corresponding values of Ly $\alpha$ , CH<sub>4</sub>,  
164 and CO<sub>2</sub> for the years considered for this study are highlighted. We classify 1992-2005 as period  
165 1 (“early”) and 2005-2018 as period 2 (“late”). Satellite observations of H<sub>2</sub>O showed a clear  
166 anti-correlation with the solar cycle in the early period, which was absent in the late period  
167 (Hervig et al. 2019). Certainly, at low and middle latitudes, without NLCs one can detect only  
168 anticorrelation. For example, in H<sub>2</sub>O satellite data averaged over the tropics (30° N-30° S), anti-  
169 correlation is observed for the "late" period (Karagodin-Doyennel et al. 2021). To investigate  
170 ~~this~~ the missing response reported in Hervig et al. 2019, we first examined the early period solar  
171 minimum (1997) and maximum (2002) in more detail. The solar cycle affects the H<sub>2</sub>O  
172 concentration in two main ways. (i) through the photolysis of H<sub>2</sub>O by Ly $\alpha$ , and (ii) through the  
173 temperature effect. We distinguish these effects by performing model simulations with different  
174 background conditions (see Table 1). Namely in section 3.3, we discuss the individual role of  
175 solar cycle-induced photolysis and temperature change on the H<sub>2</sub>O-solar cycle response. Figure  
176 2 shows that the intensity of Ly $\alpha$  radiation during the late period has decreased compared to  
177 the early period, and the concentrations of increased greenhouse gases (GHGs) have increased  
178 in the late period. The effects of reduced Ly $\alpha$  intensity and increased greenhouse gas (GHG)  
179 concentration on long-term H<sub>2</sub>O-solar cycle response are discussed in section 4.

Fig. 2

181 **3. Results and Discussions**

182

183 **3.1. Solar cycle response in ice water content (IWC)**

184

185 To determine if the model agrees with satellite observations, we compared the ice water content  
186 (IWC) anomaly from the model with the satellite observations (see Fig. 3). IWC anomalies are  
187 calculated as follows:

188 
$$IWC_{anom} = 100\% \cdot \frac{\overline{IWC_{July}} - \overline{IWC_{1981-2018}}}{\overline{IWC_{1981-2018}}}, \quad (1)$$

189 Where  $\overline{IWC_{July}}$  represent monthly zonal averages at 68°N, and  $\overline{IWC_{1981-2018}}$  are the averages  
190 of  $\overline{IWC_{July}}$  over the years 1981-2018. The IWC anomaly for satellite measurements are from  
191 the Solar Backscatter Ultraviolet (SBUV), Halogen Occultation Experiment (HALOE), Cloud  
192 Imaging and Particle Size (CIPS), and Solar Occultation For Ice Experiment (SOFIE)  
193 instruments. The time series of SBUV and HALOE data as shown in Figure 3 represent three  
194 years sliding averaged values. For more details on the satellite datasets, see Hervig et al. (2019).  
195 For this comparison, we used the MIMAS run A, in which the simulations are performed with  
196 increasing concentrations of CO<sub>2</sub> and CH<sub>4</sub>. For the comparison, we applied the same calculation  
197 method to our model data as Hervig et al. (2019) did on satellite observations, namely, we used  
198 a threshold of 50 g/km<sup>3</sup> for integrated water content because the PMC detection threshold for  
199 SBUV is 50 g/km<sup>3</sup> (DeLand and Thomas, 2015, 2019).

200 We find an anti-correlation between MIMAS IWC anomaly and Ly $\alpha$  flux throughout the entire  
201 period (1981-2018), with a weaker response in the late period. In satellite observations, SBUV  
202 measurements also show an anti-correlation with Ly $\alpha$  flux until 2005, after which the response  
203 becomes weaker in agreement with MIMAS. The magnitude of the solar cycle IWC anomaly  
204 in SBUV and HALOE is of the same order as the IWC anomaly in MIMAS. The IWC anomalies

Fig. 3



205 of CIPS and SOFIE do not show a clear response to the solar cycle. We notice that the year-to-  
206 year IWC variation in CIPS and SOFIE is larger than the IWC modulation during a solar cycle.  
207 IWC anomalies of SBUV and HALOE correlate well with MIMAS IWC anomalies before 2005  
208 and progressively weaken afterwards. Lübken et al. (2009) found a good agreement between  
209 NLC parameters calculated by MIMAS and satellite observations. The general agreement  
210 between the main characteristics and trends of the ice layers in MIMAS and the observations  
211 suggests that the microphysical and photochemical processes in MIMAS cover the main  
212 processes relevant to NLC formation (Lübken et al., 2009).

213

### 214 **3.2. Effect of NLC on water vapour (H<sub>2</sub>O)**

215

216 We calculated the zonal mean monthly averaged vertical profiles of H<sub>2</sub>O and temperature to  
217 investigate the impact of NLC formation on the H<sub>2</sub>O profile. Figure 4 shows the vertical H<sub>2</sub>O  
218 profile averaged for July at 68°N latitude and given on pressure altitudes  $z_p = H_p \ln(p_0/p)$ .  
219 Where  $p$  is the pressure of the model level,  $p_0$  is the pressure at the surface and  $H_p = 7$  km is  
220 the pressure scale height. This figure illustrates the effect of NLC formation on the background  
221 profile of water vapour since the H<sub>2</sub>O profile with NLC differs from that without NLC. In the  
222 presence of NLC there is a reduction in water vapour mixing ratio (dehydration) between 83-  
223 90 km, i.e. in the region where the saturation ratio of water vapour is larger than one. An  
224 enhancement in water vapour (hydration) is observed at altitudes between 79-83 km, where the  
225 saturation ratio of water vapour is smaller than one. An environment with a water vapour  
226 saturation ratio larger than one is supersaturated, meaning ice particles can grow under these  
227 conditions whereas a saturation ratio lower than one leads to ice sublimation. The degree of  
228 saturation depends on the background atmosphere's H<sub>2</sub>O concentration, and temperature. Ice  
229 particles formation starts at higher altitudes, where the temperature is the lowest, and then they  
230 sediment downward. During sedimentation, the ice particles grow by consuming H<sub>2</sub>O from the

Fig. 4

231 surrounding background, which decreases background H<sub>2</sub>O concentration. Then they approach  
232 a region with a saturation ratio smaller than one, where they sublimate, releasing the water  
233 vapour. This is the so-called freeze-drying effect well discussed in a number of papers (Hervig  
234 et al., 2003; Lübken et al., 2009; Bardeen et al., 2010). The results in Figure 4 illustrate the  
235 freeze-drying effect described above and also indicate that the effects of NLC on H<sub>2</sub>O are not  
236 present below ~79 km and above ~97 km. This is the novelty of the results in Figure 4. This is  
237 because the photochemical lifetime of water vapour below ~79 km becomes larger than  
238 dynamical characteristic times, and distributions of water vapour become dynamically  
239 determined. Above 97 km, the saturation ratio of water vapour is smaller than one;  
240 consequently, there is no NLC formation and consequently no effect on water vapour.

241

### 242 **3.3. Effect of solar cycle-induced temperature and photolysis changes on water vapour** 243 **(H<sub>2</sub>O)**

244

245 We investigate the temperature change between the solar minimum (1997) and maximum  
246 (2002) due to solar irradiance variation and how these changes affect the H<sub>2</sub>O profile. Different  
247 model runs performed for this study are summarized in Table 1. The differences (solar  
248 maximum - solar minimum) for H<sub>2</sub>O and temperature profiles are shown in Figure 5 for three  
249 model runs, namely E, F, and G. In run E, the solar cycle-induced temperature change and  
250 photolysis influences H<sub>2</sub>O concentration. In run F, only the temperature change caused by the  
251 solar cycle affects the H<sub>2</sub>O concentration, while in run G, only the photolysis caused by the  
252 solar cycle affects the H<sub>2</sub>O concentration (see Table 1). All of these runs are performed with  
253 constant CO<sub>2</sub> and CH<sub>4</sub> concentrations to avoid the effects of increasing GHG concentrations on  
254 temperature and H<sub>2</sub>O profiles.

255 In model run F, Ly $\alpha$  is held constant in MIMAS, so the photolysis of H<sub>2</sub>O is constant during  
256 the solar cycle. However, Ly $\alpha$  (and other bands) varies in the LIMA model, so the background  
257 temperature varies with the solar cycle. Therefore, the change in the H<sub>2</sub>O profile during the  
258 solar cycle is only due to the influence of the solar cycle on temperature and sequentially on  
259 microphysical processes. Figure 5a shows that the temperature increases during solar maximum  
260 compared to solar minimum through the entire altitude range (79-97 km). The difference of  
261 temperature amounts to ~0.5-1.7 K with maximum values at ~95 km. During solar maximum,  
262 increased solar irradiance leads to greater absorption of solar radiation in the MLT region by  
263 molecular oxygen and water vapour, which heats the background atmosphere. Temperature  
264 differences decrease as altitude decreases because the intensity of solar radiation decreases due  
265 to atmospheric absorption by molecular oxygen and water vapour. The solar cycle effect in the  
266 H<sub>2</sub>O profile with NLC (blue line) differs significantly from that without NLC (yellow line).  
267 Without NLC, the H<sub>2</sub>O profile difference is nearly zero at all altitudes, indicating that the  
268 temperature changes do not significantly affect the background H<sub>2</sub>O profile in the absence of  
269 NLC. With NLC, the H<sub>2</sub>O profile difference is positive in the altitude range of 82-87 km and  
270 slightly negative in the range from 79 to 82 km. The atmosphere is warmer during solar  
271 maximum; therefore, the ice formation rate is lower during solar maximum. When the ice  
272 formation rate decreases, the amount of water vapour consumed from the background  
273 decreases; hence, more H<sub>2</sub>O is left in the background during solar maximum compared to solar  
274 minimum, resulting in a slightly positive response at NLC forming altitudes above 83 km.  
275 Below that altitude, the slightly negative response is due to reduced ice formation in the  
276 nucleation region during solar maximum, which decreases H<sub>2</sub>O released at ice sublimation  
277 altitudes. The positive difference peak at ~83 km is located near the bottom of the H<sub>2</sub>O-saturated  
278 zone. Ice formation/sublimation is more sensitive to an increase in background temperature at  
279 this zone (where the degree of saturation is close to one) because at these altitudes the  
280 background temperature is almost equal to the frost point temperature, so an increase in

Fig. 5a

281 background temperature critically changes the degree of saturation. The change of the  
282 background temperature in a region where it is significantly lower than the frost point  
283 temperature is not critical for the degree of saturation. Overall, the temperature variation due to  
284 the solar cycle causes a positive H<sub>2</sub>O response on the solar cycle at ice-formation altitudes and  
285 a slightly negative response at ice-sublimation altitudes.

286 In model run G (Fig. 5b), we consider only the effect of solar cycle-induced Ly $\alpha$  variation on  
287 water vapour photolysis. The background temperature is held constant. Photolysis of H<sub>2</sub>O by  
288 Ly $\alpha$  radiation molecules mainly ~~leads to produce~~ atomic hydrogen (H) and hydroxyl (OH) in  
289 the upper atmosphere (~90 %) and with less extent to O(<sup>1</sup>D) with molecular hydrogen (~10 %).  
290 The photolysis rate is higher during solar maximum due to the increased Ly $\alpha$  flux caused by  
291 the increased solar activity. Without NLC, the difference in the H<sub>2</sub>O profile is negative at all  
292 altitudes (yellow line), indicating that the background H<sub>2</sub>O is reduced during solar maximum  
293 due to increased photolysis. Figure 5b shows that the negative response peaks at an altitude of  
294 ~87.5 km. The solar cycle effect on the photolysis of H<sub>2</sub>O decreases above 87.5 km because the  
295 water vapour mixing ratio decreases with increasing altitude. The solar cycle variation of the  
296 photolysis effect decreases below 87.5 km because the solar Ly $\alpha$  radiation intensity decreases.  
297 With NLC (blue line), the H<sub>2</sub>O difference between the solar maximum and the solar minimum  
298 is essentially negative at ice sublimation altitudes (below ~83 km) and negligible at higher  
299 altitudes (above ~85 km). This is due to the redistribution of the H<sub>2</sub>O profile during NLC  
300 formation (“freeze-drying”). During solar maximum, the background H<sub>2</sub>O concentration  
301 available for ice formation is reduced due to enhanced photolysis. The lower H<sub>2</sub>O availability  
302 during solar maximum results in lower ice formation and, thus, lower H<sub>2</sub>O release during  
303 sublimation, leading to lower hydration in the sublimation zone. For this reason, the solar cycle  
304 variation of the photolysis effect is more pronounced at sublimation altitudes. Above 85 km,  
305 the effect of photolysis, in the case with NLC, is minimal because of the lower availability of  
306 H<sub>2</sub>O due to dehydration by NLC.

Fig. 5b

307 Figure 5c shows a combination of both effects, namely the solar cycle-induced temperature  
308 change and photolysis effects on H<sub>2</sub>O. Without NLC (yellow line), the H<sub>2</sub>O profile shows a  
309 negative response at all altitudes, peaking at ~87.5 km similar to run G (Fig. 5b, yellow line).  
310 We found that the variation of temperature has an almost negligible effect on the H<sub>2</sub>O in the  
311 absence of NLC (see Fig. 5a, yellow line), so the negative response of water vapour without  
312 consideration of microphysical processes (yellow line on Fig. 5c) is mainly caused by the  
313 photolysis effect. With NLC (Fig. 5c, blue line), the combined effect of temperature and  
314 photolysis has a slightly positive response on water vapour in the ice formation zone (83-89km)  
315 and a negative response in the ice sublimation zone (80-83km). The slightly positive response  
316 is caused by the temperature modulation, and the negative response is primarily due to the  
317 photolysis modulation throughout the solar cycle.

318 The study proves that the water vapour response to the solar cycle is affected by the re-  
319 distribution of water in the presence of NLC. There may exist regions with positive correlation  
320 of water vapour with Ly $\alpha$  when NLC formation occurs. Without NLC, the water vapour always  
321 shows a negative correlation to the solar cycle. When comparing the effects of solar cycle  
322 modulations of temperatures and photolysis on H<sub>2</sub>O, the photolysis has a stronger effect on  
323 water vapour, however, the variation of temperature induces a positive correlation of solar  
324 irradiance and H<sub>2</sub>O.

325

#### 326 **4. Increasing greenhouse gases and reducing solar cycle**

327

328 This section examines how the increase in GHGs affects the H<sub>2</sub>O response to the solar cycle.  
329 To distinguish the GHG effects, we compared the model results with increasing CO<sub>2</sub> and CH<sub>4</sub>  
330 (Run A) to the model run with constant CO<sub>2</sub> and CH<sub>4</sub> (Run E). It is noted already that an  
331 increasing CO<sub>2</sub> concentration leads to a cooling of the middle atmosphere, and an increase in

332 CH<sub>4</sub> concentration leads to an increase in H<sub>2</sub>O concentration (see Sec.2 for details). In Figure  
333 2, the concentration of CO<sub>2</sub> and CH<sub>4</sub> increase during the late period, and at the same time, the  
334 peak of the Ly $\alpha$  flux decreases. In order to filter out the effect of reduced Ly $\alpha$  intensity, we  
335 calculated the H<sub>2</sub>O response profile per unit of Ly $\alpha$  ( $\Delta$ H<sub>2</sub>O / $\Delta$ Ly $\alpha$ ). Figure 6 shows the result  
336 for the first (1997-2002, blue line) and the second period (2008-2014, orange line) for model  
337 runs E (Fig. 6a) and A (Fig. 6b) respectively. These profiles show positive and negative  
338 responses depending on altitude. Under the conditions of constant GHGs (run E) the sensitivity  
339 of water vapour to Ly $\alpha$  does not change from the early to the late period (Fig. 6a). As expected,  
340 for the case of growing methane and carbon dioxide (run A), the sensitivity of water vapour to  
341 Ly $\alpha$  increases during the late period (orange line, Fig. 6b) compared to the early period (blue  
342 line, Fig. 6b). This is because an increase in CO<sub>2</sub> (and consequently temperature decrease) leads  
343 to an intensification of microphysical processes, hence, to the increased freeze-drying. In  
344 addition, increasing methane leads to more water vapour in the upper mesosphere, which also  
345 leads to an increased water vapour variation with solar cycle.

346 To study the effect of a decreasing Ly $\alpha$  amplitude during the late period (2008-2014), we  
347 calculated the ratio of water vapour absolute deviations between solar minimum and solar  
348 maximum for the early and late period. The amplitude of Ly $\alpha$  variation is weaker during the  
349 late period ( $\sim 1.14 \cdot 10^{11}$  [phot.cm<sup>-2</sup>s<sup>-1</sup>]/solar cycle) compared to the early period ( $\sim 1.85 \cdot 10^{11}$   
350 [phot.cm<sup>-2</sup>s<sup>-1</sup>]/solar cycle). The intensity of Ly $\alpha$  during the late period solar maximum is  
351 reduced by  $\sim 40\%$  compared to the early period. As can be seen from Figure 7a the magnitudes  
352 of positive and negative H<sub>2</sub>O responses decreased during the late period for model runs with  
353 constant GHGs (Run E). In Figure 6a, we found that the H<sub>2</sub>O sensitivity to Ly $\alpha$  flux is the same  
354 in the early and late periods for the model run with constant GHGs (Run E). Therefore, the  
355 reduced response of H<sub>2</sub>O during the late period in model run E (Fig. 7a) is only due to the  
356 reduced solar Ly $\alpha$  variation. Comparing the late period H<sub>2</sub>O response to the solar cycle from

Fig. 7

357 model runs with constant GHG (Fig. 7a, orange line) to model runs with increasing GHG (Fig.  
358 7b, orange line) suggests that both the positive and negative peak responses are enhanced by  
359 increasing GHG concentration. Due to the increased solar Ly $\alpha$  flux and greenhouse gases, the  
360 NLC and water vapour response is expected to increase during the current solar cycle 25, as the  
361 Ly $\alpha$  radiance has already exceeded the peak value of the previous solar cycle 24.

362

### 363 **5. Missing H<sub>2</sub>O-solar cycle response**

364

365 A recent study by Hervig et al. (2019) reported a missing response in H<sub>2</sub>O concentration on  
366 solar cycle after 2005. In Figure 8, we compare our model results of H<sub>2</sub>O anomaly with the  
367 satellite observations. The H<sub>2</sub>O response is averaged over the geometric altitudes of 80-85 km  
368 at 68°N. For this comparison, we used MIMAS run A, where the increasing concentration of  
369 GHG is considered. The satellite observations are shown in Figure 8 from HALOE, SOFIE,  
370 and MLS according to Hervig et al. (2019). HALOE shows a strong negative response to Ly $\alpha$   
371 (-1.7 ppmv/solar cycle) during period 1, but in SOFIE and MLS the response is almost absent  
372 (+0.2 ppmv/solar cycle) during period 2 (Hervig et al., 2019). For MIMAS, no clear H<sub>2</sub>O-solar  
373 cycle anti-correlation is noticed in the early period, but it was slightly positive in the late period  
374 in agreement with SOFIE and MLS satellite observations. To investigate the H<sub>2</sub>O response on  
375 Ly $\alpha$  variation in more detail, we analysed the vertical H<sub>2</sub>O response profile at geometric  
376 altitudes similar to the satellite observations.

Fig. 8

377 Figure 9 shows the vertical profile of H<sub>2</sub>O response in geometric altitudes for the model run  
378 with constant GHGs (run E, Fig. 9a) and growing GHGs (run A, Fig. 9b). The magnitude of the  
379 H<sub>2</sub>O response at geometric altitudes (Fig. 9) differs from that at pressure altitudes (Fig.7). This  
380 is because the geometric altitude of constant pressure levels is not constant and varies

Fig. 9

381 throughout the solar cycle but also with time due to increasing GHG. Therefore, the magnitude  
382 of the H<sub>2</sub>O response differs when converted from pressure altitudes to geometric altitudes.  
383 We focus on the 80-85 km geometric altitude range (Fig. 9. shaded region). There are positive  
384 and negative H<sub>2</sub>O response zones within this altitude range similar to Figure 7. We calculated  
385 the average H<sub>2</sub>O response over 80-85 km altitude range for MIMAS runs A and E is given in  
386 Table 2. For the model run with growing GHGs (run A), the H<sub>2</sub>O response averaged over an  
387 altitude range of 80-85 km changed from -0.01 ppm/solar cycle in the early period to 0.10  
388 ppm/solar cycle in the late period (see Table. 2). The H<sub>2</sub>O response in the late period becomes  
389 slightly positive for run A, consistent with the satellite observations of SOFIE and MLS (see  
390 Fig. 8). The vertical profile of H<sub>2</sub>O-solar cycle response clearly show that H<sub>2</sub>O response to the  
391 solar cycle is not completely missing in the late period. The missing response in the MIMAS  
392 H<sub>2</sub>O as shown in Figure 8 occurred when averaging over the 80-85 km altitude range. Figure 9  
393 demonstrates that the H<sub>2</sub>O response shows nearly equal positive and negative responses within  
394 the 80-85 km altitude range (shaded region). Therefore, averaging the response in this altitude  
395 range becomes nearly zero, as the positive and negative responses cancel out each other. When  
396 averaging over the altitude range of 80 to 82 km in the early period we receive H<sub>2</sub>O response  
397 of -0.71ppm/solar cycle and an anti-correlation between H<sub>2</sub>O and Ly $\alpha$  . The results clearly  
398 shows that the small solar cycle response in MIMAS is a consequence of averaging over an  
399 altitude range of 80-85 km. It suggests that averaging H<sub>2</sub>O response over an altitude range  
400 containing positive and negative responses may not provide a detailed understanding of the  
401 H<sub>2</sub>O-solar cycle response.

402

## 403 6. Conclusions

404

405 In this study, we used our ice particle model MIMAS along with atmospheric dynamicstransport  
406 model LIMA to investigate the response of H<sub>2</sub>O to the solar cycle from 1992-2018. We



407 investigated how NLC formation affects vertical H<sub>2</sub>O profiles by running model simulations  
408 with and without microphysics. NLC formations are shown to redistribute H<sub>2</sub>O profiles by  
409 consuming H<sub>2</sub>O from the background at ice-forming altitudes (dehydration) and releasing it at  
410 ice-sublimating altitudes (hydration) is known as the “freeze-drying” effect. To investigate the  
411 missing solar cycle response in satellite observations reported by Hervig et al. (2019), we  
412 divided the entire study period into an early (1992-2005) and late (2005-2018) period. We first  
413 investigated how the Ly $\alpha$  variation affects the H<sub>2</sub>O profile between solar minimum and  
414 maximum in the early period. The solar Ly $\alpha$  variation affects the H<sub>2</sub>O concentration at NLC  
415 altitudes mainly in two ways: through the effect of temperature change and through the effect  
416 of photolysis. To distinguish these two effects, we performed additional model simulations with  
417 different background conditions (see Table 1). We found that the modulation of water vapour,  
418 which comes through the temperature changes with solar cycle, causes a slight positive H<sub>2</sub>O  
419 response at ice-forming altitudes and a negative response at ice-sublimating altitudes. The solar  
420 cycle photolysis effect has only negative responses on the H<sub>2</sub>O profile, and this response  
421 dominates at ice sublimation altitudes with NLC conditions. Our results for the case of  
422 photolysis effect only are supported by previous simulations, which also suggest that freeze-  
423 drying significantly reduces the potential effect of Ly $\alpha$  photolysis on H<sub>2</sub>O above 82 km, while  
424 the effect is enhanced at 80-82 km, where ice particles sublimate (von Zahn et al., 2004, Lübken  
425 et al., 2009).

426 To the best of our knowledge, we have for the first time identified a positive response of water  
427 vapour to Ly $\alpha$  variation in the MLT region which is due to microphysical processes. It was  
428 assumed for a long time that water vapour only anti-correlates with the solar cycle at mesopause  
429 altitudes (e.g. Sonnemann and Grygalashvyly, 2005; and references therein). We should note  
430 that in the Martian atmosphere, where microphysical processes have a crucial role in water  
431 vapour distributions through the entire atmosphere in all seasons (e.g. Shaposhnikov et al.,  
432 2018), this effect may be important.

433 We have made a comparison between the model and satellite observations of the H<sub>2</sub>O response  
434 to the solar cycle, averaged over an altitude range of 80-85 km. The satellite observations from  
435 HALOE show a strong anti-correlation to the solar cycle in the early period, but the model  
436 shows a very small response in both the early and late periods. The vertical H<sub>2</sub>O response  
437 profiles from MIMAS show that within the 80-85 km altitude range, the positive and negative  
438 responses are almost equal in magnitude and symmetric. Therefore, averaging the response over  
439 this altitude range reduces the overall response in model, as positive and negative responses  
440 cancel each other out.

441 We also investigated the role of increasing GHG on the H<sub>2</sub>O-solar cycle response. From the  
442 early to the late period, there are mainly two factors that affect the long-term H<sub>2</sub>O solar cycle  
443 response: increasing CO<sub>2</sub> and CH<sub>4</sub> concentrations and the lower intensity of the solar cycle (see  
444 Figure 2). We found that increasing GHG concentration increased the H<sub>2</sub>O response to Ly $\alpha$ .  
445 The Ly $\alpha$  intensity during the late solar maximum decreased by 40% compared to the early solar  
446 maximum. Therefore, the overall response of H<sub>2</sub>O to the solar cycle is also decreased in the late  
447 period. It should be noted that our results have limitations as they use constant dynamics for all  
448 years. We are looking forward to a new gravity wave resolving model to investigate the effects  
449 on changing dynamics due to changing GHGs and solar activity.

450

## 451 **Appendix**

452

453 **Data availability.** The satellite data shown in this paper are reproduced from the paper by  
454 Hervig et al., 2019. Lyman- $\alpha$  data are available at <http://lasp.colorado.edu/lisird/lya/> from  
455 LASP. The data utilized in this manuscript can be downloaded from [https://www.radar-](https://www.radar-service.eu/radar/en/dataset/ArvFyujQbPGYfRqv?token=UEOfafmhOFFfWBRKONmZ)  
456 [service.eu/radar/en/dataset/ArvFyujQbPGYfRqv?token=UEOfafmhOFFfWBRKONmZ](https://www.radar-service.eu/radar/en/dataset/ArvFyujQbPGYfRqv?token=UEOfafmhOFFfWBRKONmZ)

457

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459 **Competing interests.**

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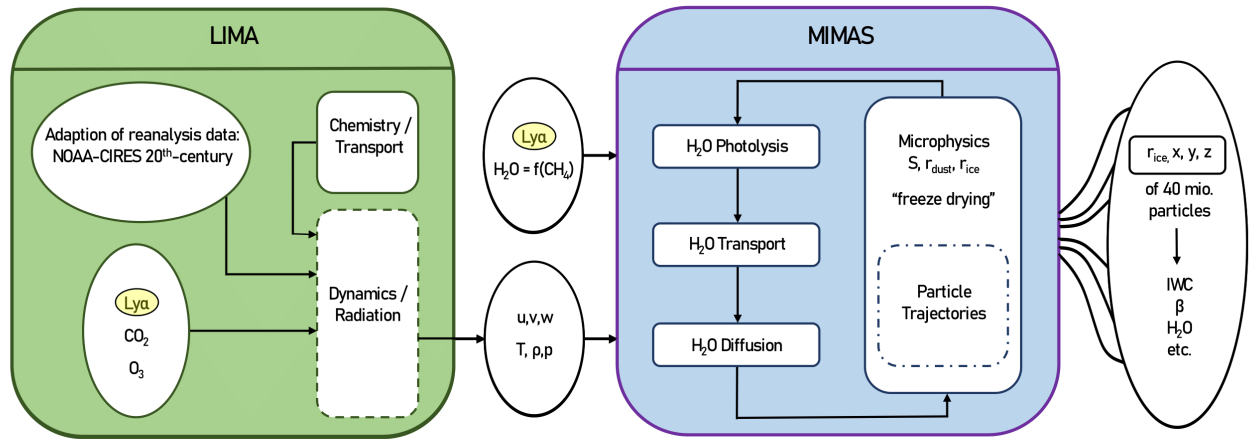
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630 **Figures**

631

632 Figure 1. Sketch of the LIMA (green) and MIMAS (blue) models (from Lübken et al., 2021)



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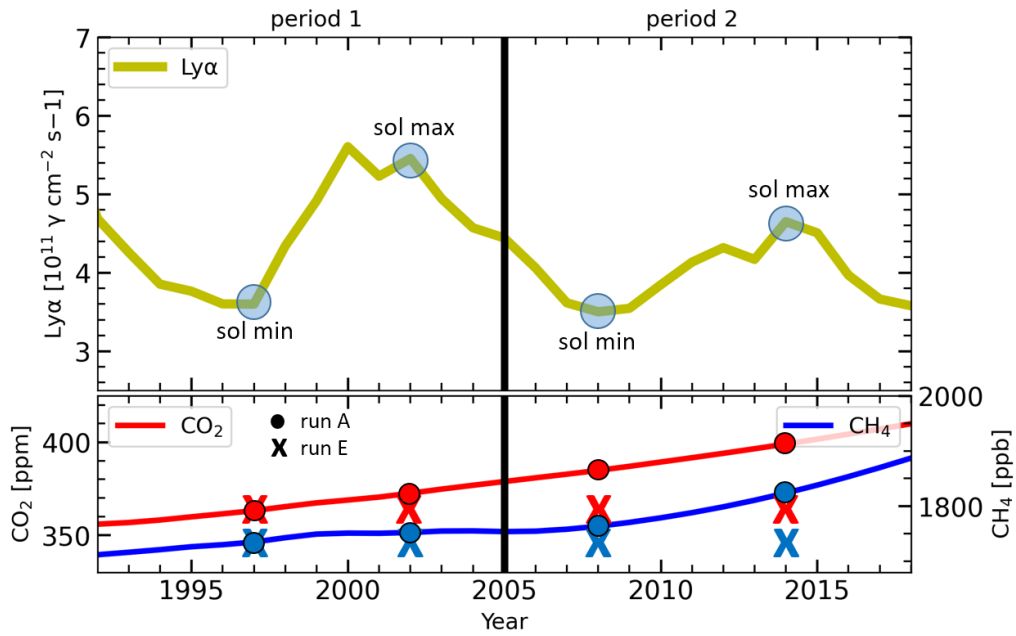
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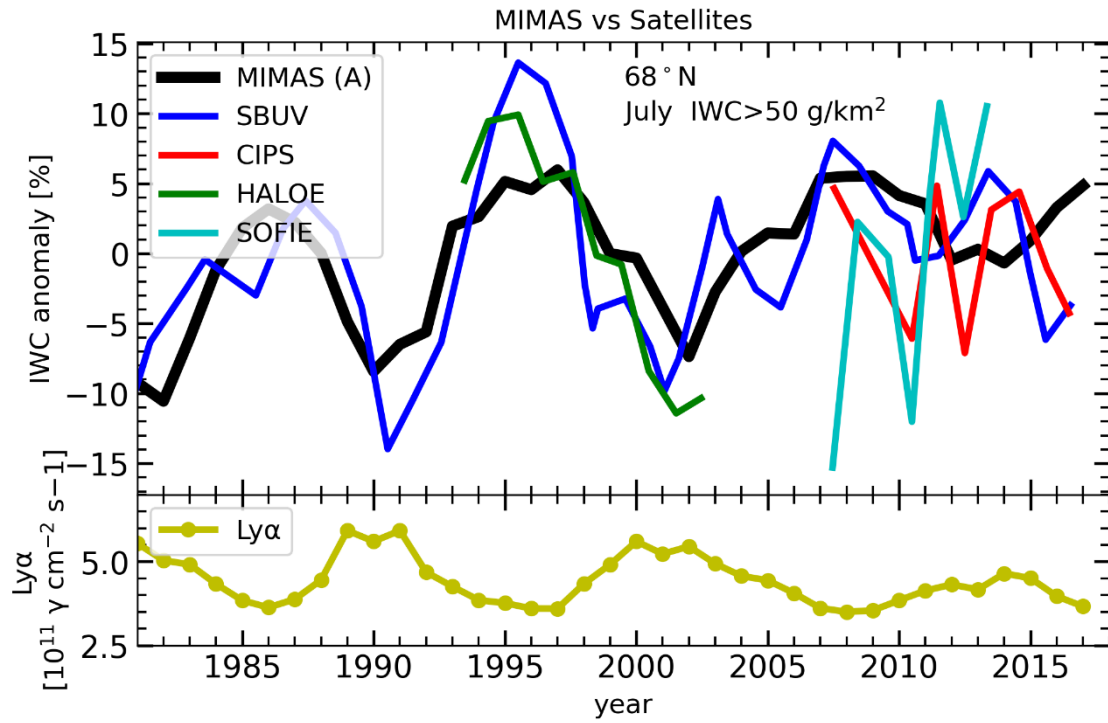
641 Figure 2. Time series of solar  $\text{Ly}\alpha$ ,  $\text{CO}_2$ , and  $\text{CH}_4$  for 1992-2018. The corresponding  $\text{Ly}\alpha$ ,  
 642  $\text{CO}_2$ , and  $\text{CH}_4$  values for the solar cycle maximum and minimum years used for this study are  
 643 marked. The  $\text{CO}_2$  and  $\text{CH}_4$  values for run A are represented with dots, and for run E with  
 644 crosses. The study period is divided into period 1 as early (1992-2005) and period 2 as late  
 645 (2005-2018).



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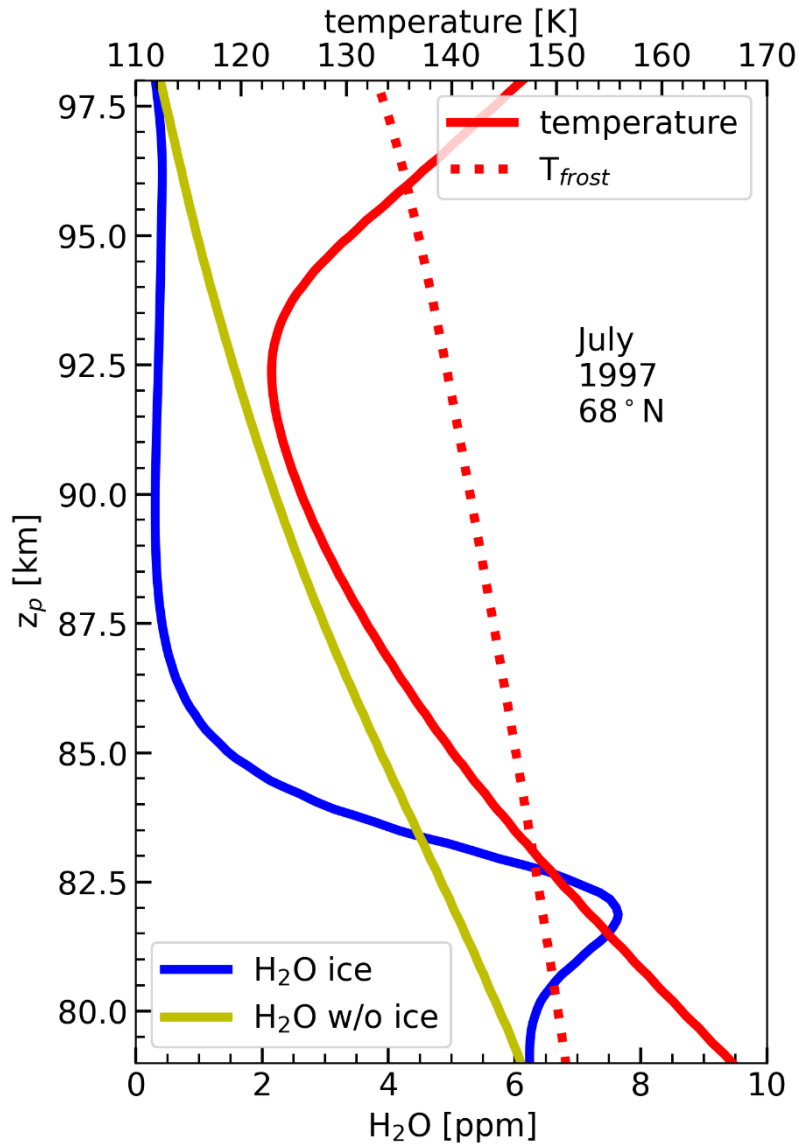
648 Figure 3. Time series of July mean IWC anomalies at 68°N from model and satellites based on  
 649 Hervig et al., (2019). Anomalies for each data set are calculated as the difference from their  
 650 long-term mean. To reduce year to year variability, the time series of SBUV and HALOE are  
 651 smoothed using the sliding average method of window size 3. Ly $\alpha$  -solar cycle modulation is  
 652 shown in the bottom panel.



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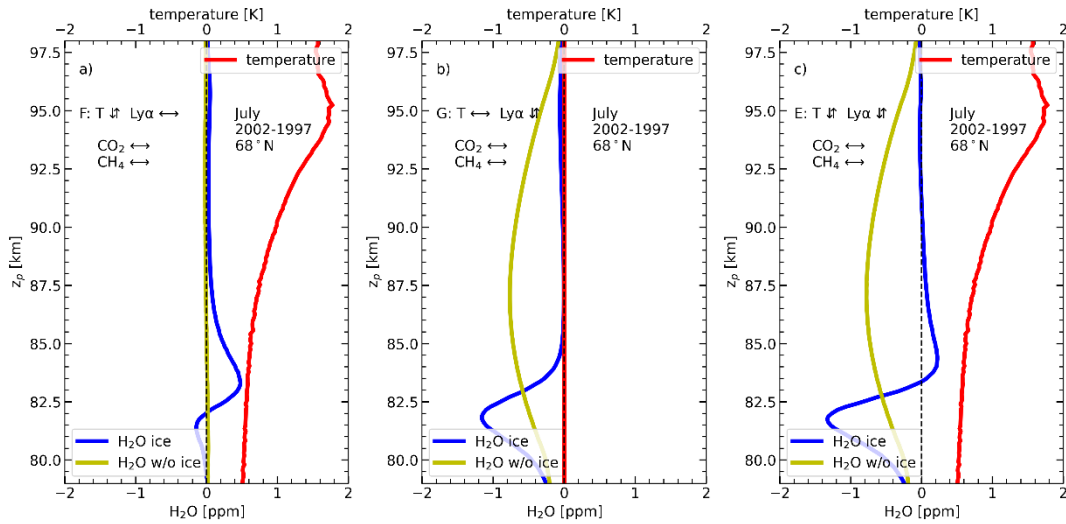
655 Figure 4. Zonally and monthly averaged H<sub>2</sub>O and temperature profile for July at 68°N from  
656 MIMAS with and without NLCs. The red dotted line represents frost point temperature. The  
657 blue lines show the background H<sub>2</sub>O concentration with NLC, and the yellow lines show the  
658 H<sub>2</sub>O concentration without NLC.



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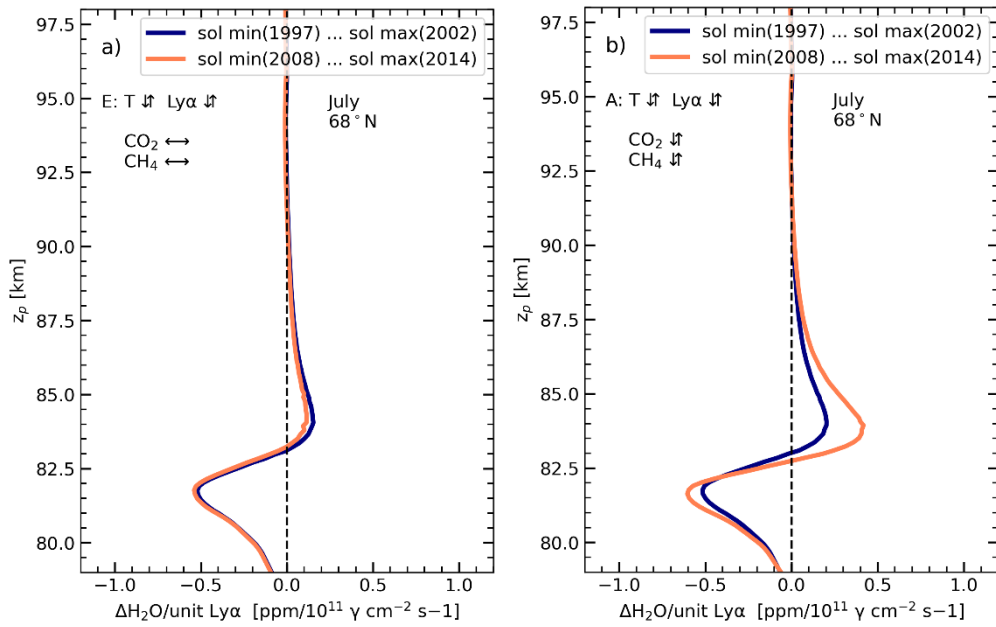
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661 Figure 5. The difference in profiles between solar maximum (2002) and minimum (1997) for  
 662 July mean H<sub>2</sub>O and temperatures. The blue and yellow lines represent NLC and non-NLC  
 663 conditions. In all cases, CO<sub>2</sub> and CH<sub>4</sub> values are constant corresponding to 1997. (a) Run F:  
 664 only temperature change effects on H<sub>2</sub>O, (b) Run G: only Photolysis change effect on H<sub>2</sub>O, (c)  
 665 Run E: both temperature change and photolysis change effects on H<sub>2</sub>O



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667 Figure 6. H<sub>2</sub>O-response per unit Ly $\alpha$  variations in July at 68°N during the years between solar  
 668 minimum and maximum in the early (1997 to 2002) and late (2008 to 2014) periods. (a)  
 669 MIMAS model run E with constant CO<sub>2</sub> and CH<sub>4</sub>, (b) MIMAS model run A with varying CO<sub>2</sub>  
 670 and CH<sub>4</sub>.



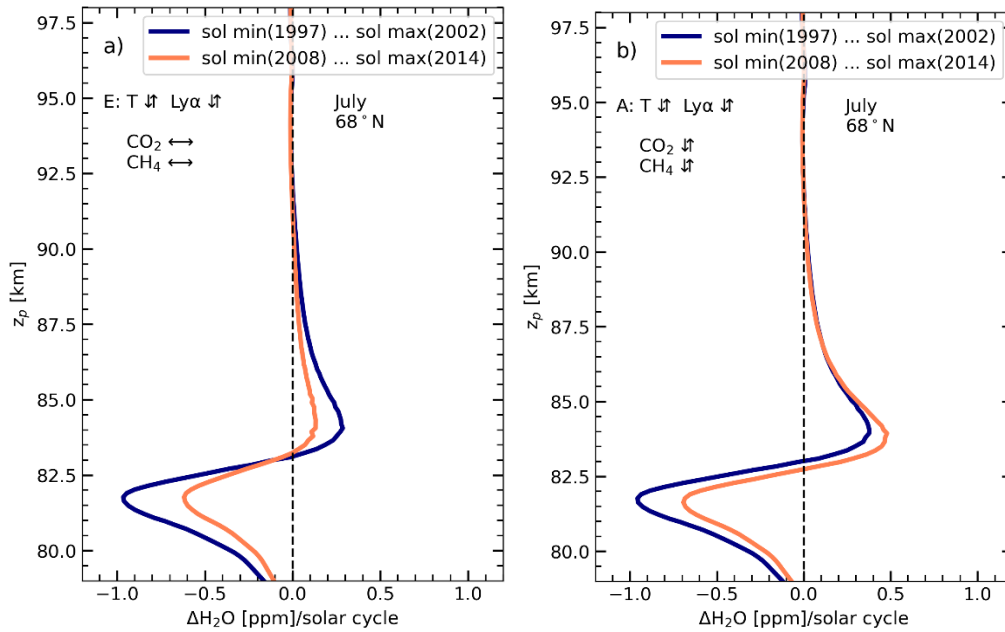
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674 Figure 7. H<sub>2</sub>O-response to absolute solar cycle Ly $\alpha$  variations in July at 68°N during the years  
 675 between solar minimum and maximum in the early (1997 to 2002) and late (2008 to 2014)  
 676 periods. (a) MIMAS model run E with constant CO<sub>2</sub> and CH<sub>4</sub>, (b) MIMAS model run A with  
 677 varying CO<sub>2</sub> and CH<sub>4</sub>.

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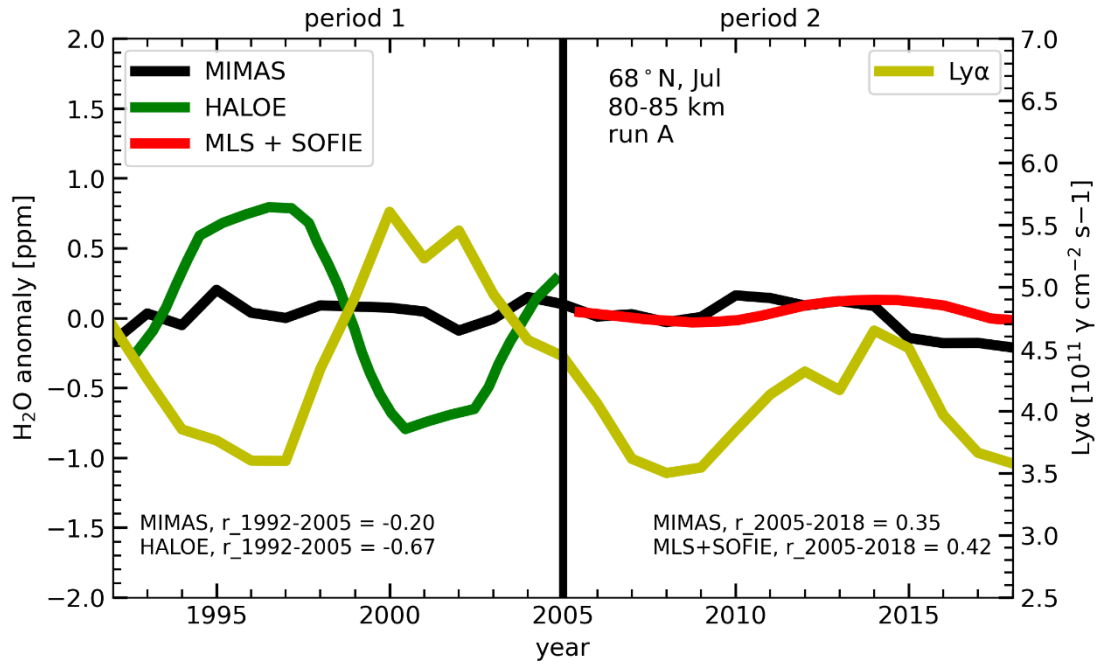


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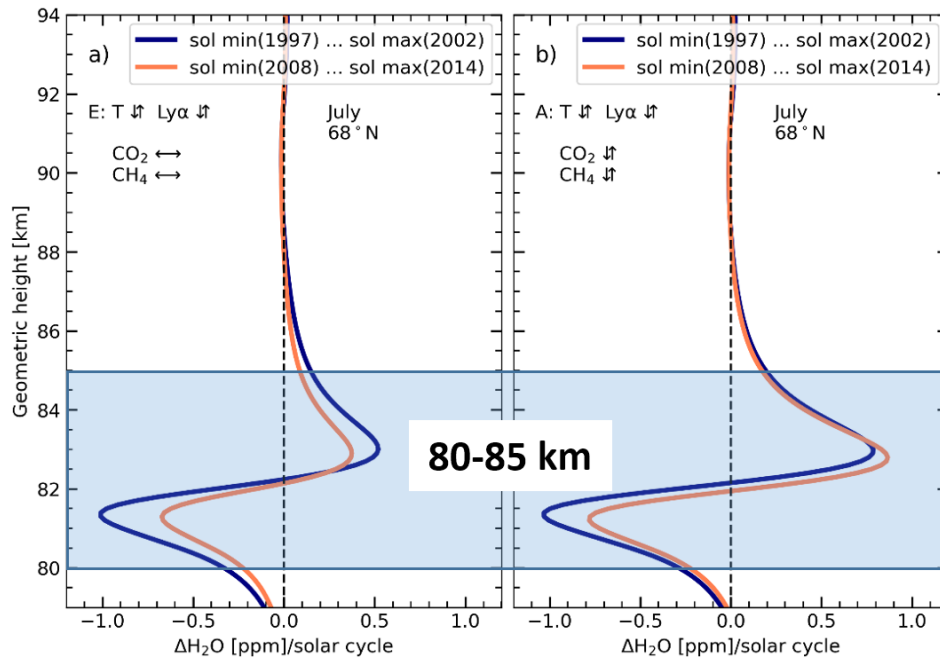
681 Figure 8. Time series of Ly $\alpha$  and H<sub>2</sub>O anomalies as monthly averages for July at 68°N for the  
 682 altitude range of 80 km to 85 km from MIMAS run A and satellites (HALOE and the composite  
 683 data (MLS and SOFIE)). Satellite observations are according to Hervig et al., 2019. The H<sub>2</sub>O-  
 684 Ly $\alpha$  correlation is calculated for the early and late periods (see inlet).



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687 Figure 9. H<sub>2</sub>O-response to absolute solar cycle Ly $\alpha$  variations in July at 68°N during the years  
 688 between solar minimum and maximum in the early (1997 to 2002) and late (2008 to 2014)  
 689 periods represented in geometric altitudes. The shaded region represents the altitudes range  
 690 used for calculating an average solar cycle response. (a) MIMAS model run E with constant  
 691 CO<sub>2</sub> and CH<sub>4</sub>, (b) MIMAS model run A varying CO<sub>2</sub> and CH<sub>4</sub>.



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706 **Tables**

707 Table 1. MIMAS simulations were carried out under different background conditions. The  
 708 horizontal arrow stands for constant values for the given year; the vertical arrow is for varying  
 709 parameters. How Ly $\alpha$  affects H<sub>2</sub>O is given for each run in the last column.

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	LIMA		MIMAS		
<b>Model run</b>	<b>CO<sub>2</sub></b>	<b>Ly<math>\alpha</math> T effect</b>	<b>CH<sub>4</sub></b>	<b>Ly<math>\alpha</math> photolysis effect</b>	<b>Water vapour solar cycle response affected by</b>
<b>A</b>	↓	↓	↓	↓	<ul style="list-style-type: none"> <li>• <b>Temperature change (Ly<math>\alpha</math> + CO<sub>2</sub>)</b></li> <li>• <b>Photo dissociation</b></li> <li>• <b>Varying CH<sub>4</sub> (H<sub>2</sub>O source)</b></li> </ul>
<b>E</b>	↔ <b>1997</b>	↓	↔ <b>1997</b>	↓	<ul style="list-style-type: none"> <li>• <b>Temperature change</b></li> <li>• <b>Photo dissociation</b></li> </ul>
<b>F</b>	↔ <b>1997</b>	↓	↔ <b>1997</b>	↔ <b>1997</b>	<ul style="list-style-type: none"> <li>• <b>Temperature change</b></li> </ul>
<b>G</b>	↔ <b>1997</b>	↔ <b>1997</b>	↔ <b>1997</b>	↓	<ul style="list-style-type: none"> <li>• <b>Photo dissociation</b></li> </ul>

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716 Table 2. The solar cycle H<sub>2</sub>O response averaged over 80-85 km geometric altitude at 68°N  
717 for model runs A and E.

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<b>Model run</b>	<b>ΔH<sub>2</sub>O (ppm)/solar cycle (80-85km)</b>	
	<b>Early period</b>	<b>Late period</b>
<b>MIMAS with constant CO<sub>2</sub> and CH<sub>4</sub> (run E)</b>	<b>-0.11</b>	<b>-0.06</b>
<b>MIMAS with increasing CO<sub>2</sub> and CH<sub>4</sub> (run A)</b>	<b>-0.01</b>	<b>0.10</b>

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