



## Signature of gravity wave propagations from the Troposphere to Ionosphere

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**Abstract.** We observed a gravity wave (GW) signature in the OH emission layer in the upper mesosphere, and 4 hours later, a medium-scale traveling ionospheric disturbance (MSTID) in the OI 630 nm emission layer. Spectral analysis of the two waves did show that both have the same wave characteristics: wavelength, period, phase speed and propagation direction, respectively, 200 km, 60 min, 50 m/s, propagating southeastward. During the MSTID occurrence, concentric wavefronts were also observed in the ionosphere by detrended total electron content (dTEC) maps. From the gravity wave ray-tracing simulation for the mesospheric gravity wave, we found that the wave came from a tropospheric deep convection spot and propagated up to the 140 km altitude. Regarding the same wave characteristics between mesospheric GW and MSTID, two possible cases are investigated: a direct influence of the GW oscillation in the OI 630 nm emission height and the generation of a secondary wave during the GW breaking process. The concentric wave structure suggests the generation of a secondary wave after the primary wave was dissipated in the lower thermosphere. This is the first time to report an observational event of gravity wave propagation from the troposphere, mesosphere to thermosphere-ionosphere in the south American region.

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### 1. Introduction

A deep cloud convection in the troposphere generates vertical (up and down) air-mass movement launching a variety of gravity waves into the stratosphere. Atmospheric gravity waves (GWs) have important roles in transporting the energy and momentum from the lower to upper atmosphere and ionosphere. A part of energy and momentum is deposited in the mesosphere lower-thermosphere (MLT) region through wave breaking and altering the background wind field. Some of the GWs produce secondary waves and propagate further upwards into the thermosphere where it modulates ionospheric plasma (Hocke and Schlegel, 1996, Nicolls et al., 2014). A part of Medium-scale Travelling Ionospheric Disturbances (MSTIDs) has its origin

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30 in the passage of gravity waves in the ionosphere (Otsuka, 2018). Observations of GW propagation in the thermosphere have been carried out by many researchers since Hines (1960) has presented theoretical background for the GW propagation in the ionosphere. Rottger (1973) suggested the role of GWs in the ionospheric irregularities.

GW observations in the mesosphere have been carried out by measuring short period temporal variation of the mesospheric airglow (Hydroxyl and atomic oxygen OI 557.7 nm emissions) by airglow photometers in 1970-1990 (e. g., Takahashi et al, 35 1999). After 1990 airglow digital imagers were used to monitor GWs in two-dimensional forms (e.g., Taylor et al. 2009, Dare-Idowu et al. 2020, Nyassor et al., 2021). In case of GWs in the stratosphere, satellite-onboard GPS radio occultation measurements have made it possible to observe GWs by vertical profile of the temperature variability on a global scale (Tsuda, 2014, Xu et al., 2017).

There are many previous works on the GW propagations in the stratosphere, mesosphere and ionosphere individually. 40 However, it has been difficult to monitor an event of GW propagating through the troposphere up to the ionosphere. Smith et al. [2013] observed GW waves in the OH and OI 557.7 nm emission layers in the mesosphere to lower thermosphere (MLT) region (85-100 km) and OI 630.0 nm emission layer (around 240 km altitude) in the thermosphere and discussed on the mountain waves from the mesosphere to ionosphere. They attributed the wave structure in the ionosphere as due to secondary waves. Azeem et al, (2015), for the first time, reported the occurrence of circular GW structures in the stratosphere, mesosphere 45 and ionosphere during a tropospheric convective storm. They observed concentric wave structures in the stratosphere by the Atmospheric Infrared Sounder(AIRS) onboard Aqua satellite (<https://airs.jpl.nasa.gov/>), and by an optical imaging radiometer (VIIRS) onboard Suomi satellite ([https://www.nasa.gov/mission\\_pages/NPP/main/index.html](https://www.nasa.gov/mission_pages/NPP/main/index.html)), and in the ionosphere by ground-based GPS receivers. Prior to this work, Nishioka et al. (2013) has reported concentric gravity waves in the ionosphere which were induced by a severe convective system (supercell) in the troposphere. The concentric waves lasted for more than 50 7 hours. Nyassor et al., (2021) reported the first mesospheric concentric gravity waves excited by thunderstorm. Takahashi et al. (2020) presented the generation and propagation of MLT-GWs and concentric MSTIDs in the ionosphere during a deep convection activity in the troposphere over the south American continent.

Regarding propagation of GWs from the lower to upper atmosphere, Vadas (2007), as the first time, studied propagation property of GWs from the troposphere to the thermosphere for the horizontal wavelength of 10 to 1000 km and the period of 55 10 to 100 min. The author presented the GW dissipation altitudes depending on their horizontal wavelength and period. In case of the horizontal wavelength of 200 km and its period of 60 min, for example, the model predicts dissipation above 120 km altitude. It means that the dissipation produces a body force and generating secondary waves.

There is a difficulty to observe a gravity wave from its origin (source) in the troposphere following up to the thermosphere. During the upward propagation, it could change its wave characteristics under the background atmosphere condition, 60 dissipating and producing secondary waves changing the horizontal wavelength, phase speed and propagation direction (Vadas and Crowley, 2010). It would take several hours to reach the mesosphere to the lower thermosphere (Vadas and Liu, 2013),



which makes it difficult to follow the wave step by step from the stratosphere to the thermosphere. Recent observation of concentric wavefronts in the stratosphere, mesosphere and thermosphere (Azeem et al, 2015) would be rather a rare case. Further observational evidence would be necessary to clarify the propagation processes. The purpose of the present work is to report a case of gravity wave propagation directly from Tropospheric convection to the mesosphere and thermosphere/ionosphere. For investigating the propagation of gravity waves, data from airglow OH imager in the mesosphere and OI 630 nm imager from the thermosphere/ionosphere, ionosondes and a GNSS receiver network are used.

## 2. Observations

Airglow observation has been carried out at Bom Jesus da Lapa (hereafter BJL), 13.3°S, 43.5°W, geomag. 14.1°S, since 2019. The observation site is located under the equatorial ionospheric anomaly (EIA) belt. Equatorial plasma bubbles can also be frequently observed. An all-sky airglow imager equipped with 3-inch optical interference filters (for 630.0 nm, 557.7 nm and OH-NIR(710-930 nm) takes 180° wide images with a time sequence of ~ 5 min. Exposure time for each filter is 15 s for the OH-NIR and 90 s for the OI 630.0 nm and 557.7 nm images. The imager characteristics have been presented by Wrasse et al. (2021).

Ionospheric Total Electron Content (TEC) maps are produced by using ground-based GNSS receiver network over the South America (Takahashi et al., 2016). The TEC maps over South America with 10 min intervals are available at EMBRACE's website (<http://www2.inpe.br/climaespacial/portal/en/>). To retrieve small amplitude of TEC variability, less than 1.0 TEC unit of spatial variation, we plot the vertical TEC time series from the individual receiver, and to pick up a fraction from running means as follows,

$$dTEC(t) = TEC(t) - \langle TEC(t \pm 30 \text{ min}) \rangle$$

where the symbol  $\langle \rangle$  means the time average between the  $\pm 30$  min interval. The temporal resolution of the dTEC map is one min. The spatial resolution is a square of  $0.25^\circ \times 0.25^\circ$ , then taking a running average with 5 points. The dTEC(t), therefore, can be plotted on the map with  $1.25^\circ$  square ( $\sim 130 \times 130$  km) by one min time interval. Any wave propagation mode, of which oscillation amplitude being less than one TEC and the temporal oscillation between 5 and 30 min can be plotted (Figueiredo et al. 2018, Essien et al. 2021).

In the present work, the data from the ionosondes were used to observe the vertical drift of the F-layer and to calculate the electron density profile. Three digital ionosondes (DPS-4) (<http://www.digisonde.com/instrument-description.html>) have been operated one at São Luís (2.6°S, 44.2°W, geomag. 3.9°S), Fortaleza (3.9°S, 38.4°W, geomag. 6°S), and Cachoeira Paulista (22.7 S, 45.0 W, geomag. 18.1° S). The DPS-4 sounder has a 500 W peak power, covering a frequency range from 0.5 to 30 MHz. Ionograms are taken at a time interval of 10 minutes.

## 3 Results

### 3.1 OH images



On the night of 18-19 January 2020, around 22:30 to 23:30 UT (19:30-20:30 Local Time), the airglow OH images showed two wave crests passed over BJL. Figure 1 shows the OH images at the moment of the wavefront passing over the zenith. The images are projected on the geographic coordinates. The horizontal extension is approximately 500 km, and the blue dot indicates the location of the BJL site. Two sequential images with a time interval of 10 minutes indicate that the wavefront is propagating toward the Southeast (indicated by an red arrow). The propagation mode can be clearly seen if keograms are made by using the meridional and longitudinal cuts as a function of time. Figure 2 shows the keograms (zonal and meridional cuts) between 22:30 and 28:30 (04:30) UT. Two bright bands propagating from North to South and West to East are the wavefronts in question. The broad bright band passing over the zenith from NE to SW through the night is the Galactic Milky way.

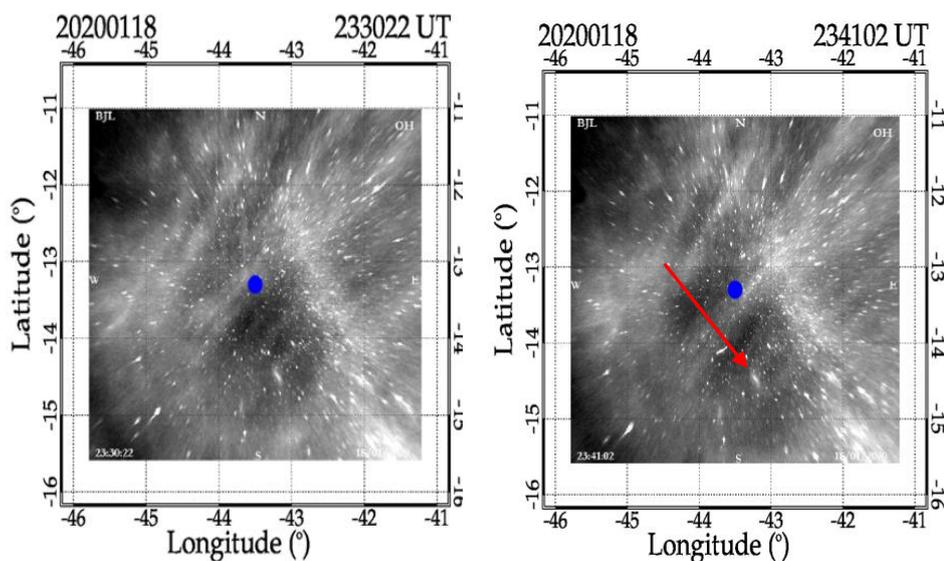


Fig. 1. Geographically coordinated OH images observed at Bom Jesus da Lapa (BJL), 13.3°S, 43.5°W, geomag.14.1°S, at 23:30 UT (left) and 23:41 UT (right) on the night of 18-19 January 2020. Blue dots indicate the zenith of BJL. The red-arrow indicates the propagation direction of the waves.

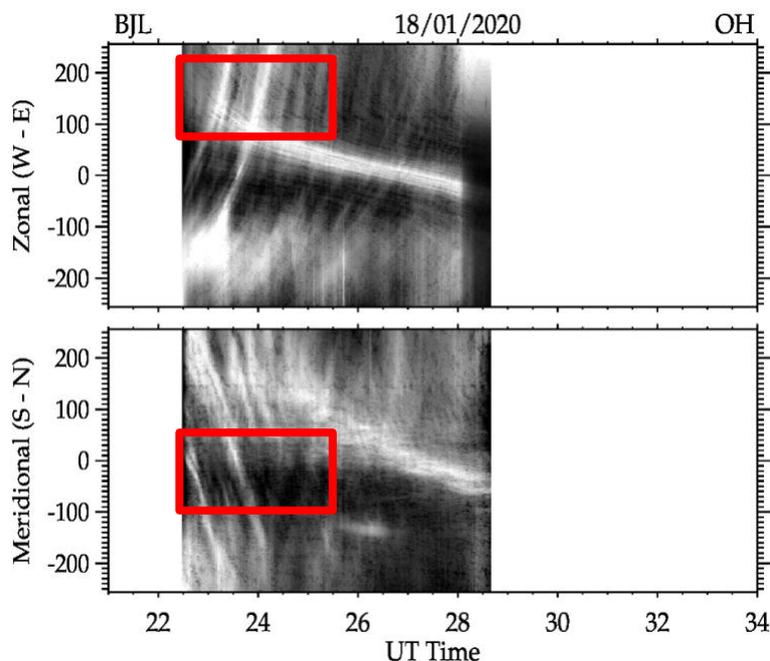


Fig. 2. Keograms of OH images observed at BJJ in the night of 18-19 January 2020. Meridian cuts along the W-E direction (upper panel) and the S-N direction (lower panel) are shown as a function of time from 22:00 to 30:00 (06:00) UT. The red color rectangular boxes are where the FFT spectral analysis were taken.

For calculating the wave characteristics (horizontal wavelength, period, phase velocity), we used Fast Fourier Transform (FFT) spectral analysis (Wrasse, 2007, Figueiredo et al., 2018, Essien et al., 2018). Image samples used in the calculation are indicated by red boxes in Figure 2. The wave characteristics obtained are, the horizontal wavelength of  $217.9 \pm 12$  km, the period of  $60.6 \pm 03$  min, the phase speed of  $59.9 \pm 5$  m/s, and the propagation direction of  $148.5 \pm 10^\circ$ .

### 3.2 OI 630 nm image

Figure 3 presents two sequential images of the OI 630 nm emission at 03:20 UT and 03:41 UT.

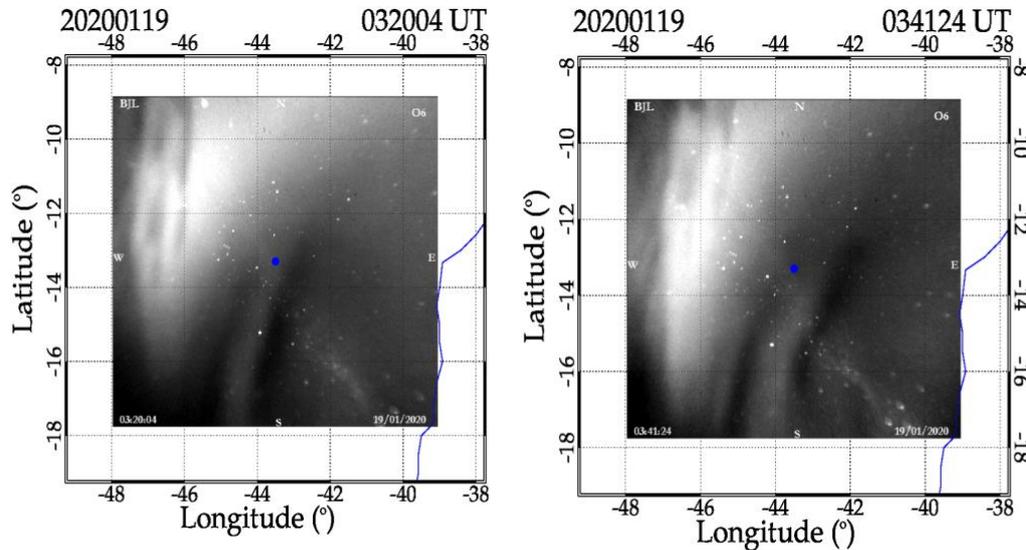


Fig. 3. Geographically coordinated OI 630 nm images observed at Bom Jesus da Lapa (BJL) at 03:20 UT (left) and 03:41 UT (right) on the night of 18-19 January 2020. Blue dots indicate the location of BJL. MSTID (lower/middle), EPBs (upper left corner) and MTM (upper left) can be seen (see Text).

125 The two dark band structures in the southwest of BJL are propagating toward the East. We checked any contamination of the OH emission in the OI 630 nm image. No such wavelike structure could not observe in the OH images in the same period. The bright OI 630 emission intensity over the northwest part of the sky should be the midnight downward drift of the F-layer accompanied by Midnight Temperature Maximum (MTM) in the thermosphere (Colerico et al., 1996, Figueiredo et al., 2017). One can also notice the presence of the equatorial plasma bubbles (EPBs) (at least two) in the northwest of BJL, which is also drifting toward the East. From the FFT spectral analysis of the OI630 keogram (not shown here), the wave characteristics were obtained: the horizontal wavelength of  $201.7 \pm 13$  km, the period of  $64.2 \pm 33$  min, the phase speed of  $52.4 \pm 27$  m/s with the propagation direction of  $113.2 \pm 10^\circ$ . The characteristics of the wave propagation in the OH emission layer and OI 630 nm emission layer are summarized in Table 1.

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**Table 1.** Wave Characteristics obtained by the OH images in the MLT region, OI 630 nm images in the thermosphere and dTEC maps in the ionosphere. (The values in parentheses indicate the error range.)

	OH (23-00 UT)	6300 (03-04 UT)	dTEC (02-04 UT)
$\lambda_H$ (horiz. wave length) (km)	217.9 (12.3)	201.7 (13.4)	360 (50)
$\tau$ (period) (min)	60.6 (3.0)	64.2 (33.2)	26.5 (2)
$V_p$ (horiz. phase speed) (m/s)	59.9 (4.5)	52.4 (26.9)	226 (20)
Az (azimuth) (deg)	148.5 (10.0)	113.2 (10.0)	Circular

### 3.3 dTEC map

140 A small amplitude of perturbation of the ionospheric total electron content (TEC) was retrieved from the dTEC analysis mentioned in the previous section. Figure 4 depicts a snapshot of the dTEC map at 03:44 UT (00:44 LT) on 18-19 January 2020. The red squares indicate where the TEC disturbance was larger than 0.2 TECu from the running average level. The spatial resolution of the map, in this case, is around 130 km, and the temporal resolution is 1 min. From previous works (Figueiredo et al., 2018, Essien et al., 2021), it is known that the amplitude of oscillation of TEC during the passage of MSTID  
145 is between 0.1 to 0.2 TECu. Our present results, the dTEC variation of -0.4 to +0.2 TEC are similar to the previous works.

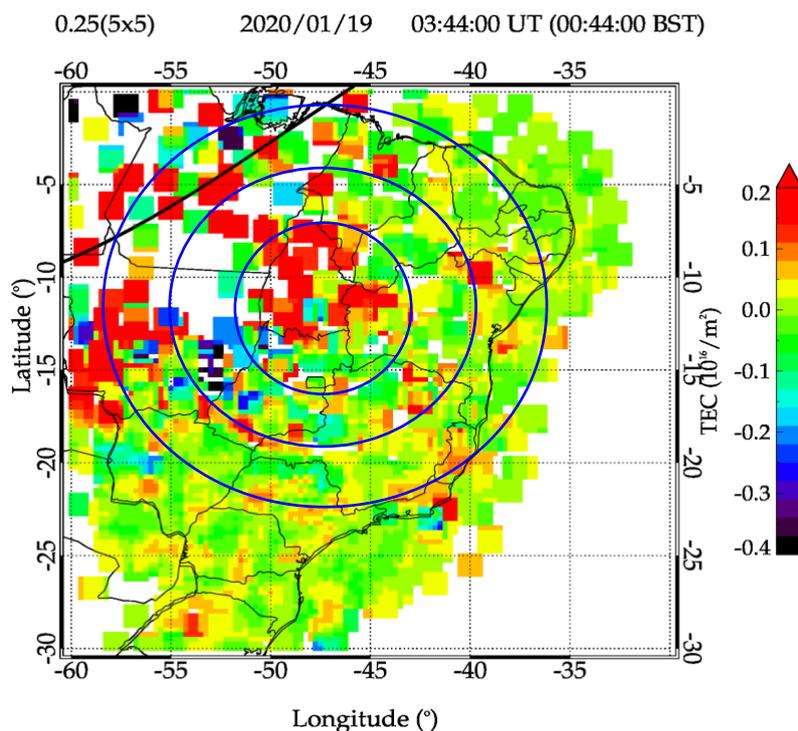


Fig. 4 dTEC Map in the equatorial and low latitude of Brazil, taken at 03:44 UT (00:44 LT at 45W), 18-19 January 2020. The black line indicates the geomagnetic equator and the blue circular lines indicate wave crests of MSTIDs. The top/left corner “0.25(5x5)” indicates a spatial resolution of the map as being  $1.25^\circ \times 1.25^\circ$  (~130 x 130 km).

The accumulated red squares at around (7- 13°S, 45- 50°W) in Figure 4 indicate a region where the ionosphere is disturbed. One can further notice that there is some increase of disturbance (the color of yellow to red squares) on the circles indicated by light blue lines, mainly from the East to South. The center of the circles falls on (12°S, 47°W) where the dTECs are disturbed. Although the circles are not whole around but partial, these suggest a concentric wave structure. Such partial  
155 concentric structures have been reported by Nishioka et al. (2013), who attributed it to the wind filtering effect. From the distance between the circles, the horizontal wavelength was estimated to be  $360 \pm 20$  km. To obtain the wave characteristics, keogram analysis was applied. Longitudinally (latitudinally) sliced dTECmap at 7.5°S (47.5°W) as a function of time is shown in Figure 5. During the time interval of 02:30 to 03:30 UT in the 8° to 14°S region (highlighted by rectangular boxes) one can notice a periodic structure of dTEC perturbation in both longitudinal and latitudinal cuts, indicating a passage of MSTID. The  
160 disturbance appeared around 01:30 to 02:30 UT in the 3° to 7°S region was due to the equatorial plasma bubbles. Although the spatial resolution of the keogram image is not enough to estimate the horizontal wavelength, the periodicity could be obtained, which is  $26.5 \pm 2.0$  min. The phase velocity is, therefore calculated to be  $226 \pm 15$  m/s.

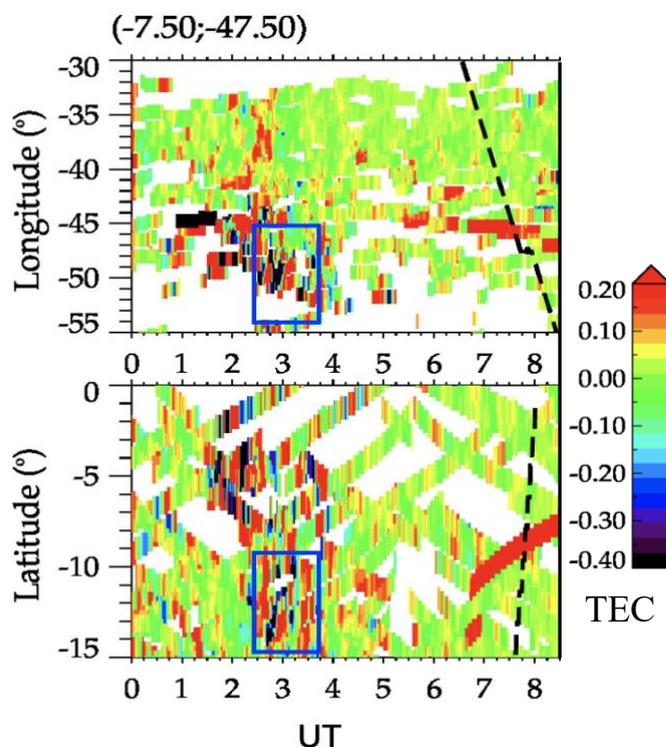


Fig. 5. Keograms of dTEC cut longitudinally at 7.5°S (top) and latitudinally at 47.5°W (bottom) as a function of Time (UT).

165 The rectangular box highlights where the FFT spectral analysis was applied.



## 4 Discussion

Through the evening to midnight over the Bom Jesus da Lapa (BJL) airglow observation site (17°S, 38°W) on 18-19 January 2020, we observed a relatively slow speed GW in the OH emission layer (~87 km altitude) at around 23:00 UT. In 4 hours  
 170 later (03:00 UT), we observed a wave structure in the OI630 nm emission layer in the lower ionosphere (~240 km altitude). The two different waves, one from MLT and the other from the thermosphere, had almost the same wave characteristics, i. e., a same horizontal wavelength ( $210 \pm 10$  km), same period ( $62 \pm 5$  min) and the same phase speed ( $55 \pm 5$  m/s) with the propagation direction of  $130 \pm 20^\circ$  (SE). The OH wavefronts are extended longer than 500 km. On the other hand, the OI 630  
 175 structure called our attention to further investigate whether these waves have the same origin from the lower atmosphere, i. e., both are primary waves, or one of the waves in the thermosphere was due to the secondary wave generated in the lower atmosphere.

### 4.1 GW Ray-Tracing and Tropospheric Convection Origin

For studying the wave propagation, we used a wave ray-tracing method (Paulino et al. 2013, Vadas et al., 2019, Nyassor et al.,  
 180 2021) to find out the source of the waves in the lower atmosphere. The wind model used in this work was according to the NRLMSISE-00 (Picone et al., 2002) and Horizontal Wind Model (HWM14) (Drob et al., 2015). Figure 6 presents the Ray-Tracing trajectories of the wave for the case of no-wind and with-wind model for the GWs in the OH layer.

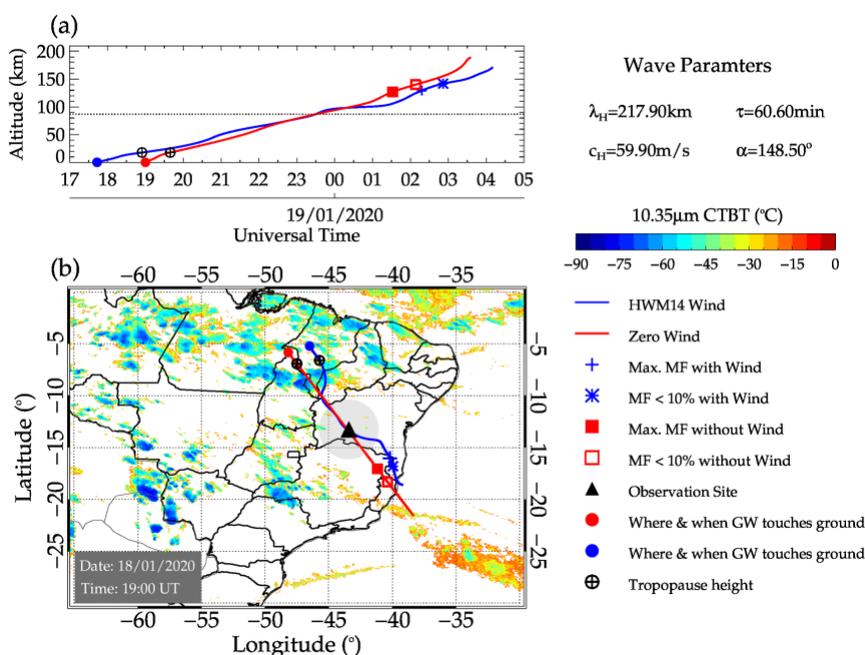




Fig. 6. Ray tracing (backward and forward) of the observed gravity wave ( $\lambda_h=217.9$  km) starting at 87 km altitude at BJL on the night of 18 January 2020 at 23:00 UT. The vertical trajectory versus time (top) and horizontal distance (bottom) are shown. The red line is of the case of no-wind and the blue line is of with-wind model. The blue triangle is the starting point of the ray tracing. The background map is Cloud top temperature from NOAA Goes 16 meteorological satellite data ( $10.35 \mu$ ) on 18 January 2020 at 19:00 UT.

The simulation started from 87 km altitude at 23:30 UT and went down (backward tracing) to the ground level at around 19:00 UT, crossing the tropopause ( $\sim 15$  km) at around ( $7.0^\circ\text{S}$ ,  $46.5^\circ\text{W}$ ) at 19:30 UT. There is a difference of around 200 km of the tropopause crossing positions between the no-wind and with-wind trajectories, which can be assumed to be an error range in the present study. Then, we search for any convective system in this region. Figure 6 also shows the cloud top temperature map produced by GOES-16 ( $13.3 \mu$  radiation map) at 19:00 UT (<https://www.cptec.inpe.br/>). It can be seen the convection system spread over the tropical zone  $0$  - $10^\circ\text{S}$ . It is the Intertropical Convergence Zone (ITCZ). One can notice that there is the lowest temperature spot ( $-80^\circ\text{C}$ ) at ( $8.5^\circ\text{S}$ ,  $46.5^\circ\text{W}$ ), where a deep convection was in progress. One can notice that this convection spot is located very close (in an error range of 200 km) to the GW trajectory at the tropopause height. According to the GOES-16 maps, this convection spot started at around 18:00 UT, developing into a much larger area from 19:00 to 23:00 UT and decreasing the intensity after 00:00 UT. During the 5 hours of activity, the convection spot should generate up and down streams inside of the convection cell producing a variety of GWs. The present Ray-Tracing suggests that the observed gravity wave in the OH emission layer started from this convection spot propagating up to the lower thermosphere.

#### 4.2 GW breaking in the thermosphere and generation of a secondary wave

The forward Ray-Tracing shown in Figure 6, on the other hand, went up to 130 km with the momentum flux in maximum and then it lost the amplitude of oscillation at around 140 km, indicating dissipation of the wave energy. The wave dissipation occurred at the location of ( $18^\circ\text{S}$ ,  $40^\circ\text{W}$ ) where we observed the wave structure in the OI 630 nm image starting at around 03:00 UT. According to Vadas and Crowley (2010), GW dissipation produces a body force and generating secondary waves. The secondary waves have a variety of wave characteristics. In our present case, we understand that the primary wave observed at MLT region dissipates in the lower thermosphere then a secondary wave reaches at the OI 630 emission height, which is located at around 240 km altitude. If this is the case of what happened, the secondary wave had the same characteristic as the primary wave. Vadas and Becker (2018) have discussed in the small- and large-scale secondary waves. According to them, there will be two kinds of secondary waves, one is small scale waves (short horizontal wavelengths) that will be produced during the primary wave breaking process, and the other one is the much longer wavelength (thousands of km) which are produced by a body force generated after the primary GW dissipation. The latter is dependent on the spatial scale of the body force. Bossert et al. (2017), for example, observed secondary waves near the primary (mountain) wave breaking area and found that the horizontal wavelengths are shorter than the primary waves. Smith et al. (2013) reported nearly simultaneous observation of mesospheric GWs by OH airglow and thermospheric GWs by OI 630 nm images. According to their



observation, the horizontal wavelength of OH against OI 630 nm is 106 vs. 255 km, the phase speed of 49.5 vs. 104 m/s, and the period of 36.5 vs. 42.7 min. Our present case (relatively short wavelength and low phase speed) could be the first case, i.e., secondary wave generated through the breaking process.

#### 220 4.3 Possible direct influence of primary wave in the ionosphere

The other possibility of the presence of GWs is a direct influence of the primary wave in the OI 630 nm emission layer. The Ray-Tracing simulation for the MLT GWs did show its dissipation at around 140 km (Figure 5). According to Vadas (2007), the signature of GWs in the thermosphere could be observable even at one or two local density scale heights (15-20 km at around 150 km altitude) above the dissipation altitude. If this is the case, the influence of the primary wave could reach at  
225 least at the altitude of 170-180 km where the F-layer bottom height is located. It is worth to check, therefore, the OI 630 nm emission height over BJL during the GW occurrence (03:00-04:00 UT).

The airglow OI 630 nm emission is produced by the dissociative recombination process in the ionosphere:



Where O(<sup>1</sup>D) is in an excited state of atomic oxygen responsible to emit a photon of 630,0 nm. The emission rate, therefore,  
230 depends on the concentration of the electron density [e] and its height profile (Chiang et al., 2018). The electron density profile, especially its bottom side profile, could be estimated by ionosonde. Unfortunately, there is no ionosonde at BJL. During the passage of the waves at 03:30 UT (00:30 LT), two DPS ionosondes, one at Fortaleza (7°S, 38°W), North of BJL, and the other at Cachoeira Paulista (22.7°S, 45.0°W), South of BJL, were in a routine observation mode. The Fortaleza ionogram showed the F-layer peak height (hmF2) at 220 km. It is very low altitude, due to the Midnight Collapse of the Ionosphere  
235 (Gong et al., 2012). On the other hand, the ionosonde at Cachoeira Paulista observed the peak at 260 km. From the estimated electron density profiles, we calculated the OI 630 nm volume emission profiles based on the equation presented by Chiang et al. (2018). The peak emission height at Fortaleza was at 200 km. On the other hand, at Cachoeira Paulista it was at 240 km. The peak height at Fortaleza is 40 km lower than Cachoeira Paulista. The BJL site is located between the two ionosonde sites. Therefore, we assume that the OI 630 nm emission peak height at BJL might be at around 200 to 240 km. Therefore, the  
240 possibility of the effect of the primary wave breaking reaches at 180 km resulting perturbation of the OI 630 nm emission layer could not be ruled out.

#### 4.4 GWs observed by dTECmap

As shown in Figure 4, we could see a concentric wave structure in the dTEC map. The center of the circular waves is located at around (12.0°S, 47.0°W), which is about 4.0° South of the convection spot. If one compares the wave characteristics obtained  
245 by the the dTEC maps and the 630 nm images, the horizontal wavelength of the former is much longer (360 km) and the period is shorter (26.5 min). It means that the phase speed of the former is much faster (226 m/s). Although the characteristics of the two waves are different, they seem to have a same origin. During the period of 19:00 to 23:00 UT, the ITCZ convection spot developed quickly forming a cluster of convection at around (8.0°S, 46.0°W). It could be possible that these convection spots launched a variety of GW spectra. In order to simulate the wave propagation path of the circular fronts, the backward Ray-



250 Tracing was also applied for one of the concentric waves. The wavefront location and its altitude were assumed to be at (17.0°S, 43.0°W) and 300 km altitude, respectively, and started from 03:40 UT. The simulated backward trajectory is shown in Figure 7, where one can see the trajectory went down toward NW and reached the minimum height at around 100 km altitude at (12.5°S, 48.0°W). It means that this concentric wave did not come from the lower atmosphere but was generated above 100 km. Vadas and Azeem (2020) presented (see Figure 21 of Vadas and Azeem (2020)) the backward ray-tracing of the GWs  
 255 from 300 km altitude for the waves with the different phase speeds, and found that GWs with a phase speeds faster than 200 m/s could not propagate below ~100 km altitude. Our present case is similar to theirs. The horizontal distance between the wavefront and the location of the minimum height (at 100 km altitude) is approximately 805 km. This location almost coincided with the center of the concentric waves. These facts and the GW propagation simulation suggest that the concentric waves started above the mesopause and propagated upwards. It indicates that they are not primary waves from the troposphere,  
 260 but a secondary gravity wave generated in the lower thermosphere.

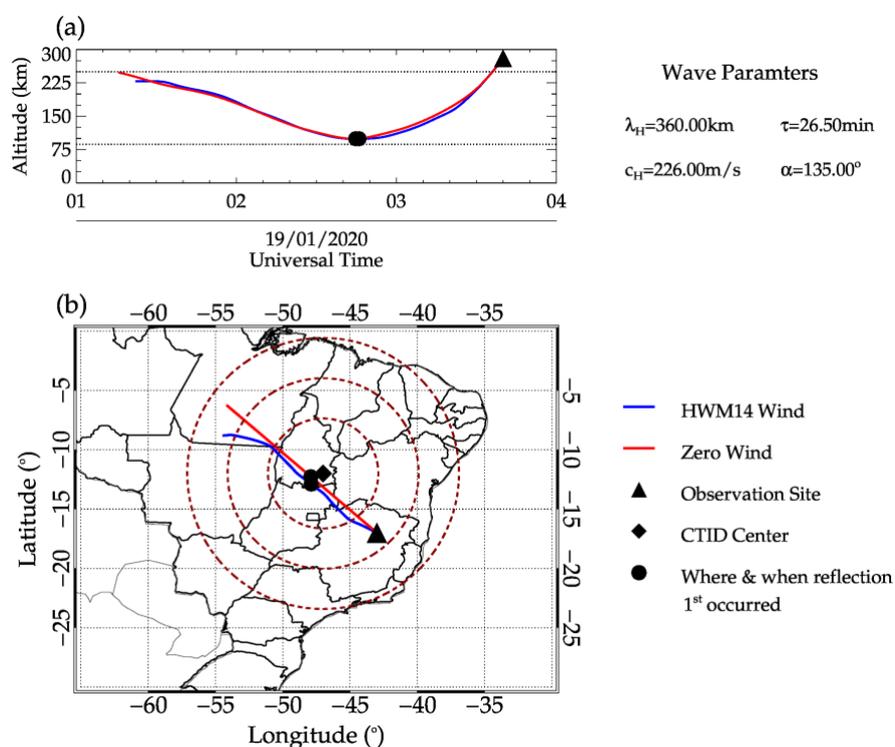


Fig. 7. Backward Ray-Tracing of the MSTID circular waves, starting from 300 km altitude at (17.0°S, 43.0°W) at 03:40 UT. The vertical trajectory versus time (top) and horizontal distance (bottom) are shown. The red line is in the case of no-wind and the blue line is of with-wind model. The black triangle is the starting point of the ray tracing. The black circle at (12.5°S, 48.0°W) indicates where the wave was reflected.  
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## Conclusion

We observed a gravity wave signature at the OH emission height (~ 87 km) and OI 630 nm emission height (~240 km), those showed the same wave characteristic. Although the two waves look to be similar, the wave observed in the ionosphere might be a secondary wave. However, a direct influence of the primary wave in the OI 630 nm emission layer at around 200 km altitude cannot be ruled out. We also observed fast concentric waves by the dTEC maps, which might be a secondary wave generated by a primary wave dissipated in the lower thermosphere. Both slow- and high-speed waves have their origin from a convective spot in the ITCZ region. This is the first time to report the direct evidence of GW propagation from the troposphere to the ionosphere by optical imaging and dTEC measurements in the south American region.

## 275 Data availability

GNSS ground-based receiver data, airglow image data, and ionosonde data used in the present study are available at the EMBRACE data center website (<http://www2.inpe.br/climaespacial/portal/en/#>). The satellite infrared thermal images (Figure 7) are obtained from the Geostationary Operational Environmental Satellite System 16 (GOES 16) data (<http://satellite.cptec.inpe.br/home/index.jsp>), provided by the Center for Weather Forecasting and Climate Studies (CPTEC) in Brazil. Two atmospheric models were used in computing the gravity wave ray tracing (Figure 6): one is the MSIS-E-00 Atmospheric Model: [https://ccmc.gsfc.nasa.gov/modelweb/models/msis\\_vitmo.php](https://ccmc.gsfc.nasa.gov/modelweb/models/msis_vitmo.php) and the other is the empirical Horizontal Wind Model (HWM14) (Drob et al., 2015).

## Acknowledgements

The present work was supported by CNPq (Conselho Nacional de Pesquisa e desenvolvimento) under the grants 310927/2020-0, 150569/2017-3, 161894/2015-1, 303511/2016, 300322/2022-4 and 306063/2020-4; Fundação de Amparo à Pesquisa do Estado de São Paulo (FAPESP) under the grant 2018/09066-8 and 2019/22548-4; and Coordenação de Aperfeiçoamento de Pessoal de Nível Superior (CAPES) under the process BEX4488/14-8. Paulino thanks to Fundação de Amparo à Pesquisa do Estado da Paraíba for the grants Demanda Universal Edital 09/20221 and Edital PRONEX termo de concessão 002/2019.

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