



Thermal electron anisotropy driven by kinetic Alfvén waves in the Earth's magnetotail

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Abstract. Thermal and subthermal electron populations in the Earth's magnetotail are usually characterized by pronounced field-aligned anisotropy that contributes to generation of strong electric currents within the magnetotail current sheet. Formation of this anisotropy requires electron field-aligned acceleration, and thus likely involves field-aligned electric fields. Such fields can be carried by various electromagnetic waves generated by fast plasma flows interacting with ambient magnetotail plasma.

- 5 In this paper we consider one of the most intense observed wave emissions, kinetic Alfvén waves, that often accompany fast plasma flows in the magnetotail. Using two tail seasons (2017, 2018) of MMS observations we have collected statistics of 80 fast plasma flows (or bursty bulk flows) events with distinctive enhancement of intensity of broadband electromagnetic waves (kinetic Alfvén waves). We show correlation the intensity of electric fields of kinetic Alfvén waves and characteristics of electron anisotropy distributions: the parallel electron anisotropy increases with magnitude of the wave parallel electric field.
- 10 Also the energy range of this electron anisotropic population is well within the expected acceleration range for assumed kinetic Alfvén waves. Our results indicate an important role of KAWs in production of thermal field-aligned electron population typically observed in the Earth's magnetotail.

1 Introduction

- 15 Electron population of the Earth's magnetotail plays an important role in magnetotail current sheet structure and dynamics. Being mostly magnetized (in contrast to well scattered and isotropic ions, e.g. Walsh et al. (2011); Wang et al. (2013)) electron



population can be quite anisotropic (Walsh et al., 2011; Artemyev et al., 2014), and such anisotropy in the high- β plasma sheet may be significant source of plasma currents (Hau, 1993; Artemyev et al., 2020; Kamaletdinov et al., 2020). These currents contribute to both the cross-field plasma sheet current flowing in the dawn-dusk direction (e.g., Zelenyi et al., 2011; Artemyev et al., 2016; Mingalev et al., 2018) and to field-aligned currents (Walsh et al., 2013; Artemyev et al., 2017) contributing to magnetosphere-ionosphere coupling.

There are two distinct types of electron anisotropic populations in the Earth's magnetotail. In the dipolarized magnetotail (e.g., at dipolarization fronts and behind them) the thermal (< 10 keV) electron population is predominantly transversely anisotropic (with $T_{e\perp} > T_{e\parallel}$; see Khotyaintsev et al. (2011); Fu et al. (2012); Runov et al. (2013); Zhang et al. (2018)), what results in whistler-mode wave generation (Khotyaintsev et al., 2011; Panov et al., 2013; Le Contel et al., 2009; Breuillard et al., 2016; Grigorenko et al., 2020). This anisotropy is due to electron betatron heating at the dipolarization front (Fu et al., 2013; Birn et al., 2012, 2013), and thus observations of such transversely anisotropic electrons are quite transient. In the quiet plasma sheet with stretched magnetic field lines the thermal electron population is predominantly field-aligned anisotropic (with $T_{e\parallel} > T_{e\perp}$; see Hada et al. (1981); Stiles et al. (1978); Artemyev et al. (2014)). The strongest patterns of such field-aligned anisotropic electrons, electron beams, contribute to electrostatic waves and nonlinear electron holes widely observed at plasma sheet boundary layers (Matsumoto et al., 1994; Tong et al., 2018). Origins of such anisotropy are not yet revealed. Adiabatic acceleration during magnetotail dipolarization may, in principle, contribute to field-aligned anisotropy (Sergeev et al., 2001; Birn et al., 2014), but adiabatic heating during magnetotail convection should mostly contribute to transverse anisotropy (Lyons, 1984; Zelenyi et al., 2013). The perspective mechanisms for field-aligned electron anisotropy generation are ionospheric outflow (Walsh et al., 2013; Khazanov et al., 2014) and electron acceleration by field-aligned electric fields. Such fields are carried by broadband electrostatic noise dominated by electron holes and solitary waves (Mozer et al., 2015; Vasko et al., 2017b; Malaspina et al., 2018) and by kinetic Alfvén waves (Chaston et al., 2012, 2014; Ergun et al., 2015; Stawarz et al., 2017). Although, electron holes and solitary waves can shape field-aligned electron populations (Vasko et al., 2017a), their generation results from the relaxation of electron beams (Omura et al., 1996; Goldman et al., 1999; Lotekar et al., 2020), i.e. electron field-aligned anisotropy should be origin of such nonlinear structures. Kinetic Alfvén waves are predominantly generated around energy release regions (e.g., magnetic reconnection, see Liang et al., 2016) and effectively accelerate electrons along magnetic field lines (Damiano et al., 2015; Damiano et al., 2016; Artemyev et al., 2015). Thus, these waves transfer ion kinetic energy to electron heating and can be responsible for generation of field-aligned electron anisotropy.

Kinetic Alfvén waves (KAWs) are widely observed (Chaston et al., 2014; Ergun et al., 2015) and modeled (Lin et al., 2017; Sharma Pyakurel et al., 2018; Cheng et al., 2020) in the plasma sheet, and their dispersion guarantees a finite field-aligned electric field (Stasiewicz et al., 2000), that potentially can accelerate plasma sheet electrons along magnetic field lines. Although the typical KAW phase speed (about Alfvén speed, that is $v_A \sim 100 - 500$ km/s in the plasma sheet) is insufficiently high for Landau resonance with plasma sheet electrons ($m_e v_A^2 / 2 < 1$ eV), the large effective wave potential $\sim E_{\parallel}/k_{\parallel} \sim 1$ keV (Damiano et al., 2015; Artemyev et al., 2015; Wang et al., 2019a) results in resonance widening (Palmadesso, 1972; Karimabadi et al., 1990). Electrons fast bouncing along magnetic field line can be trapped into such potential (Watt and Rankin, 2009) or reflect from it (Kletzing, 1994), and both effects would result in electron heating along magnetic field lines. However, so far



there were no statistical evidence of KAWs contribution to field-aligned electron anisotropy in the plasma sheet. In this study we use observations of Magnetospheric Multiscale (MMS) Mission (Burch et al., 2016) in the Earth's magnetotail to confirm the key KAWs' role in shaping plasma sheet electron anisotropy.

55 Following Chaston et al. (2012), we consider bursty bulk flows (Baumjohann et al., 1990; Angelopoulos et al., 1992) as a main source of KAWs. For statistics of such flows collected in 2017-2018 MMS tail seasons, we identify KAWs from their dispersion properties (e.g. Chaston et al., 2014) and investigate correlations of KAWs characteristics and thermal electron anisotropy.

2 Spacecraft instruments and dataset

60 We focus on the KAW signatures associated with the fast plasma flows in the Earth's magnetotail at radial distances of $\sim 15 - 30 R_E$. We pick up all intervals lasting more than 10 minutes and associated with increased plasma bulk velocities. Then we compare observed magnetic and electric field spectra with the theoretical predictions for KAW (see details below) and keep in dataset only intervals that are characterized by:

- plasma flow component along Earth-Sun line GSE $V_x > 100 \text{ km/s}$ during interval
- 65 – spacecraft are near the neutral plane (mean value of the GSE $|B_x| \leq 10nT$)
- Observed wave field spectra are close to theoretical predictions

For two tail seasons (2017, 2018) we have selected 81 events (see spatial distribution on Fig. 1 and two typical examples on Figs. 3 and 4).

We use burst-mode fields and plasma data measured on board MMS1 spacecraft (Burch et al., 2016): magnetic field data with 70 8ms time resolution from the fluxgate magnetometer (Russell et al., 2016); electric field data with 1ms time resolution from Spin-plane Double Probe instrument/Axial Double Probe instrument(SDP/ADP) (Ergun et al., 2016; Lindqvist et al., 2016; Torbert et al., 2016); plasma moments, energy and pitch-angle distributions from the Fast Plasma Investigation (FPI) (Pollock et al., 2016) with 150ms and 30ms for ions and electrons, respectively. Hereinafter in our calculations we use the field-aligned coordinate system: the parallel direction (\parallel) is along background magnetic field vector (we use low-pass filtering with the 75 cutoff frequency of 0.01 Hz), the first perpendicular component (\perp_1) is along ion bulk velocity component perpendicular to the background magnetic field, and the second perpendicular component (\perp_2) completes the right-hand system. We subtract background magnetic field from the measured magnetic field vector to estimate $> 0.01\text{Hz}$ magnetic field perturbations (see example on Fig. 2). The same procedure is applied to the electric field vector. We also construct continuous wavelet spectra (using Morlet wavelets) of the parallel and perpendicular components of magnetic and electric field perturbations.

80 In the rest reference frame KAWs have frequencies less than the ion cyclotron frequency ω_i , whereas plasma flow speed in our study usually exceeds 100 km/s (with mean value of $\sim 400 \text{ km/s}$). Thus, we expect that in the spacecraft frame all the spectrum properties will be dominated by the Doppler shift, i.e. $\omega_{sc} = \omega + \mathbf{k} \cdot \mathbf{v} \gg \omega$ and the upper frequency limit of the interest will be well higher than ω_i . So in our analysis we consider wave perturbations with frequencies up to the lower hybrid

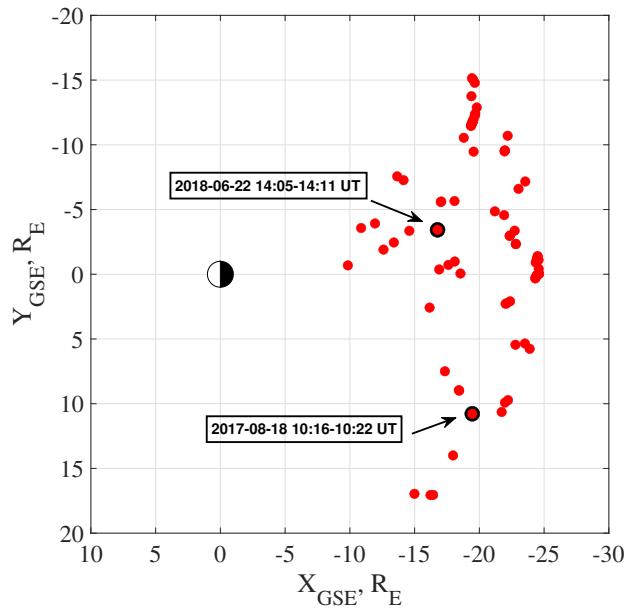


Figure 1. Coordinates of all events.

frequency (see details below). It has been shown (see Stasiewicz et al., 2000; Chaston et al., 2012) that for warm plasma in the low frequency limit ($\omega \ll \omega_i$) KAWs show following polarization properties:

$$\frac{E_\perp}{B_\perp} = \frac{v_A (1 + k_\perp^2 \rho_i^2)}{\sqrt{1 + k_\perp^2 (\rho_i^2 + \rho_s^2)}} \quad (1)$$

and

$$\frac{B_\parallel}{\sqrt{B_{\perp 1}^2 + B_{\perp 2}^2}} = i\beta \frac{k_\perp v_A}{\omega_i} \frac{1}{\sqrt{1 + k_\perp^2 (\rho_i^2 + \rho_s^2)}} \quad (2)$$

where v_A is local Alfvén speed, ρ_i is ion gyroradius, ρ_s is ion acoustic gyroradius, k_\perp is the transverse component of the wave vector \mathbf{k} , $\beta = 4\pi n(T_e + T_i)/B_0^2$, B_0 is the magnitude of background magnetic field, $B_{\perp 1}$ and $B_{\perp 2}$ are two transverse wave magnetic field components.

Note that we deal with quite inhomogeneous plasma and magnetic fields of fast plasma flows, but we use a local dispersion relation of KAWs (Chaston et al., 2012) and omit the effects of plasma gradients (Bashir et al., 2019) and nonlocal wave properties (Huang et al., 2018; Wang et al., 2019b). These effects can result in deviation of theoretical dispersion predictions from the observed wave characteristics. To estimate a possible role of plasma inhomogeneity we evaluate theoretical equations both for local plasma parameters and for parameters at some distance from the equatorial plane. We use the vertical stress balance condition $p = B^2/8\pi + n(T_i + T_e) = const$ and definition of the magnetotail lobe pressure $p = B_{lobe}^2/8\pi$ to determine magnetic fields at the distance where plasma pressure is 30% of the equatorial value: $B_{middle}^2/8\pi = 0.7p$, $n(T_i + T_e)_{middle} =$

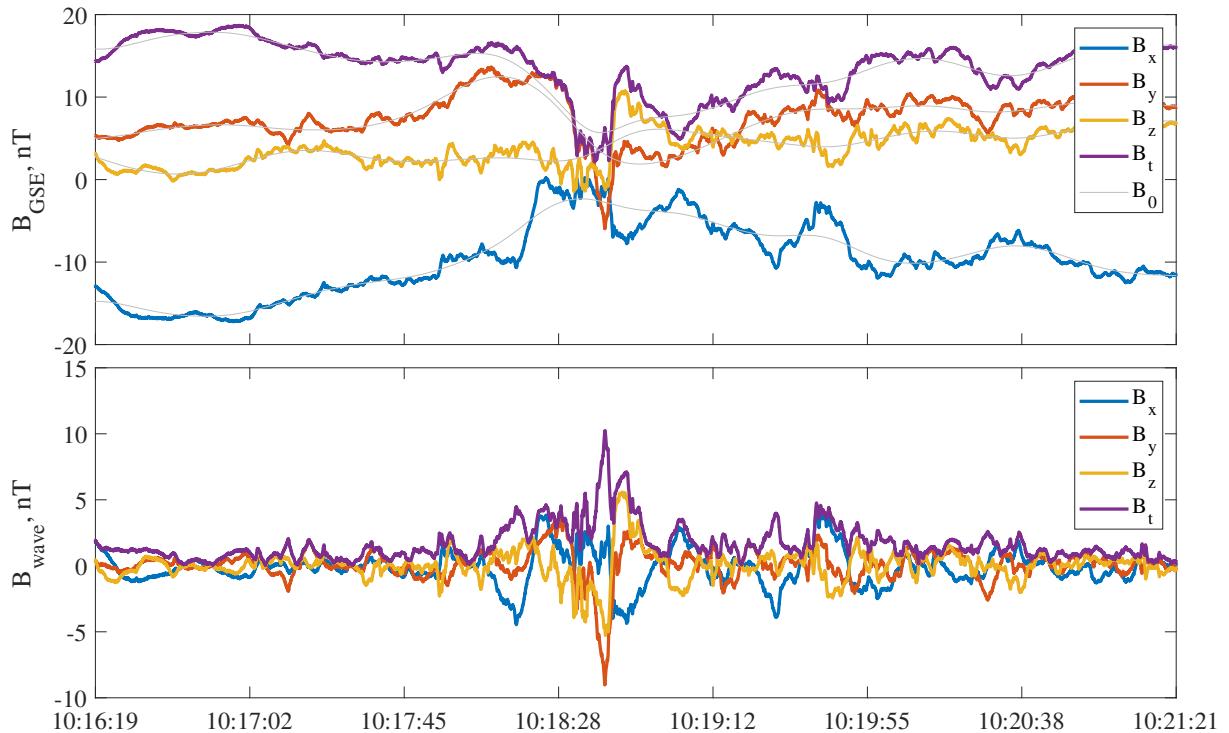


Figure 2. Wave (perturbation) magnetic field determination: gray lines show low-pass filtered background field.

0.3 p . We additionally assume $n = const$ (based on statistics of current sheets in the middle tail, see Runov et al. (2006)) and $T_e/T_i = const$ (based on comparison of $T_e(B_x)$ and $T_i(B_x)$ profiles in the middle tail, see Artemyev et al. (2011b, a)). Thus, we provide theoretical estimates for $B_{\parallel}/\sqrt{B_{\perp 1}^2 + B_{\perp 2}^2}$ at some distance from the equatorial plane together with estimates for local plasma parameters (see yellow curve in the Fig. 5, 6). Estimates at some distance from equatorial plane are used only to show the range of theoretical expectations for wave dispersion across the plasma sheet.

We aim at statistically analyzing the relation of electron flux (thermal) anisotropy and KAW electric field intensity. Such analysis should reveal a possible KAW role in shaping of the anisotropic electron populations. For such statistical comparison, we consider normalized electron energy spectra and electric field spectra. We use mean absolute value of convectional electric field $|V_i \times B_0|$ for electric field magnitude normalization, mean value of electron temperature T_e and mean absolute value of KAW electrostatic potential $|\phi|$ for electron energy normalization (averaging is performed for the MMS1 probe within the time interval of enhanced ion bulk velocity; $|V_x| > 200$ km/s, typically lasting for ~ 5 minutes). To estimate ϕ we consider the



110 parallel component of wave vector k_{\parallel} :

$$\phi(t) = -Im \left(\int_{-\infty}^{\infty} \frac{E_{\parallel}(\omega)}{k_{\parallel}(\omega)} e^{i\omega t} \frac{d\omega}{2\pi} \right) \quad (3)$$

Thus, knowing the transverse component k_{\perp} and wave frequency in the spacecraft frame ω , we estimate the parallel k_{\parallel} from equation (Chaston et al., 2008):

$$k_{\parallel}^2 v_A^2 k_{\perp}^2 \rho_s^2 = \left(\omega^2 - \frac{k_{\parallel}^2 v_A^2}{\xi_i^2} \right) \left(1 - \frac{k_{\parallel}^2 v_s^2}{\omega^2} I_0(\alpha_i) e^{-\alpha_i} \right) \quad (4)$$

115 where $\alpha_i = k_{\perp}^2 \rho_i^2$, $\xi_i = (1 - I_0(\alpha_i) e^{-\alpha_i}) / \alpha_i$, v_s is the ion acoustic speed and I_0 is the zeros order modified Bessel function.

3 Case-study

We start with two example events from our dataset shown in Figs. 3, 4. Both events are characterized by enhanced electric and magnetic field fluctuations and strong plasma flow with $V_i \sim 500$ km/s. There is magnetic field dipolization front with B_z increase (Runov et al., 2009; Nakamura et al., 2001) at $\sim 18 : 18 : 20$ for the first event and at $\sim 14 : 07 : 00$ for the second 120 one. The dipolization front separates two electron populations with quite different pitch-angle distributions and energy spectra (the sharper difference is seen for the second event); e.g. for the [0.2, 2] keV energy range (panel (f) of Fig. 4) the pitch-angle distribution becomes dominated by field-aligned flows behind the front. Panel (h) in Fig. 4 shows energy distribution of flux anisotropy calculated as

$$A_{\perp} = \frac{F_{\perp} - F_{\parallel}}{F_{\perp} + F_{\parallel}} \quad (5)$$

125 where $F_{\parallel} = \max(F_{parallel}, F_{anti})$,

$$\begin{cases} F_{parallel} : \alpha \in [0^\circ, 18^\circ]; \\ F_{anti} : \alpha \in [162^\circ, 180^\circ]; \\ F_{\perp} : \alpha \in [72^\circ, 108^\circ]; \end{cases} \quad (6)$$

and α is pitch-angle. We use SPEDAS software (Angelopoulos et al., 2019) to calculate F_{\perp} , F_{\parallel} fluxes (note FPI pitch-angle resolution is 6°). The energy range of [0, ~ 200] eV is contaminated by photoelectrons. To exclude this population from the consideration, we use electron data only for > 100 eV energy range.

130 Figure 4 shows parameters of magnetic and electric field perturbations in the field-aligned coordinate system after subtraction of the low-pass filtered background field. Black vertical lines here bound time intervals which we use for spectral analysis. Previous investigations of KAWs with Cluster (Escoubet et al., 2001) and THEMIS (Angelopoulos, 2008) measurements (Chaston et al., 2005, 2008, 2012, 2013) rely on theoretical relations that provide an estimate of the parallel electric field component from the measured perpendicular one. In our work we use direct measurements of the parallel electric component 135 available on MMS (Torbert et al., 2016; Ergun et al., 2016).



There is a clear enhancement of the wave activity in the wide frequency range for the electric field components (see panels (d,e)) and mainly up to ion cyclotron frequency for the magnetic field components (see panels (f,g)). Time-averaged wave spectra are shown in Figs. 5, 6: up to the $\sim 0.1\text{Hz}$ magnetic and electric spectral densities (transverse components) have almost the same profiles (left panels of Figs. 5, 6), whereas $E_{\perp}/B_{\perp} \sim \text{const} \sim v_A$, i.e. we deal with Alfvénic perturbations.

140 The ratio E_{\perp}/B_{\perp} grows with frequency, as it is expected for KAWs (Chaston et al., 2012). Panel (b) compares observed ratio E_{\perp}/B_{\perp} and predicted from the KAW's dispersion relation: up to the lower-hybrid frequency observations follow the theoretical model quite well, and for frequencies larger than the lower-hybrid frequency we likely deal with another wave mode (e.g., lower-hybrid drift waves or whistler-mode waves often observed around the dipolarization fronts, see Le Contel et al. (2009); Deng et al. (2010); Khotyaintsev et al. (2011); Panov et al. (2013); Grigorenko et al. (2020); Wendel et al. (2019);
145 Huang et al. (2016)). Thus, we consider electric and magnetic wave characteristics only below the lower-hybrid frequency.

At the same time, some compressional modes are presented in the ultra-low frequency range ($B_{\parallel}/B_{\perp} > 1$). Most likely KAWs are mixed with some ultra-low frequency compressional (diamagnetic perturbations) that are typical for high plasma pressure magnetotail (Volwerk et al., 2003; Runov et al., 2014). Observed frequencies of such perturbations are dominated by the Doppler shift and we cannot distinguish them from KAWs. However, contrast to KAWs, ultra-low frequency diamagnetic
150 perturbations are MHD modes that are not expected to carry any significant field-aligned fields. Thus, we can assume that electron field-aligned acceleration (if any) is associated with observed KAWs, that are well seen in E_{\perp}/B_{\perp} dispersion for the intermediate frequency range.

Figures 5, 6 demonstrate that observed wave activity can be associated with KAWs in the range of $\sim [10^{-2}, 10^1]\text{Hz}$, up to the lower-hybrid frequency. However, there are likely additional compressional (with larger B_{\parallel}) low frequency modes that
155 contribute to the observed spectra around $\sim 10^{-2}\text{Hz}$. Panels (c) show that KAWs dispersion underestimate B_{\parallel}/B_{\perp} for low frequencies and slightly overestimate it for $\sim [10^{-1}, 10^1]\text{Hz}$ frequency range. We also plot theoretical curve for plasma and magnetic field parameters corresponding to some distance from the equatorial plane (see details above), and the observed B_{\parallel}/B_{\perp} values are between these two theoretical curves. Thus, although difference between observed B_{\parallel}/B_{\perp} and theoretical predictions for KAWs in $\sim [10^{-1}, 10^1]\text{Hz}$ frequency range can be explained by strong variation of plasma characteristics
160 across the plasma sheet, KAWs theory underestimates compressional magnetic fields at lower frequencies, where observations are likely dominated by diamagnetic fluctuations typical for dipolarization fronts (Runov et al., 2014) and fast plasma flows (Bauer et al., 1995; Panov et al., 2012; Volwerk et al., 2004a).

To reveal the relation between KAWs' electric fields and electron anisotropy, we consider 2D distributions of flux anisotropy in energy and electric field space. For this we bin a field amplitude range and average all energy distributions over the time
165 interval bounded by black vertical lines on Figs. 3, 4 within the same bin of field amplitude shown in Figs. 7, 8. There are three distinguishing energy ranges with different anisotropy (note we normalize energy to instantaneous electron temperature measurements):

1. $E/T_e \lesssim 10^{-1}$
2. $10^{-1} \lesssim E/T_e \lesssim 2 - 4$



170 3. $E/T_e \gtrsim 2 - 4$

The first (lowest) range can be filled by transversely anisotropic electrons of unknown origin (this could be an effect of real transverse acceleration of cold electrons at dipolarization fronts, see Birn et al. (2014)).

175 The second (intermediate) energy range shows the parallel electron flux anisotropy. This energy range is filled by thermal ($E/T_e \in [0.1, 3]$) electrons that are expected to be most influenced by KAWs electric fields (Damiano et al., 2015; Artemyev et al., 2015). Moreover, this electron population is responsible for the thermal electron anisotropy in the magnetotail (Walsh et al., 2011; Artemyev et al., 2014).

The third (highest) energy range is characterised by perpendicular anisotropy that is presumably associated with betatron electron acceleration at the dipolarization front (Fu et al., 2013; Birn et al., 2012, 2013).

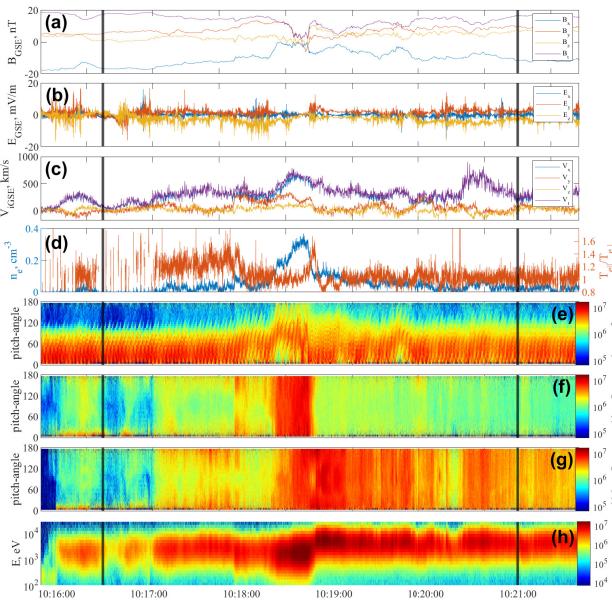
180 Note, that in the considered frequency range perpendicular electric field component usually have larger values than parallel one (both for our observations and for the theoretical predictions), so right panels on Figs. 7 and 8 cover a wider range of amplitudes than left panels do. Moreover, events with larger parallel field are rarer, i.e in the rightmost bin on the left panels we have very poor statistics.

185 In this study we focus on the second (thermal and subthermal) electron population. Figures 7, 8 show that the field-aligned anisotropy of this population increases with the growth of the electric field amplitude. Note we use both parallel and transverse fields: whereas the best correlation of anisotropy is expected for field-aligned fields, accuracy of this field measurements in the dipolarization configuration (when the main magnetic field component is almost along the spacecraft spin axes) is lower than accuracy of perpendicular component measurements. KAW dispersion assumes that larger perpendicular electric field component would correspond to larger field-aligned component.

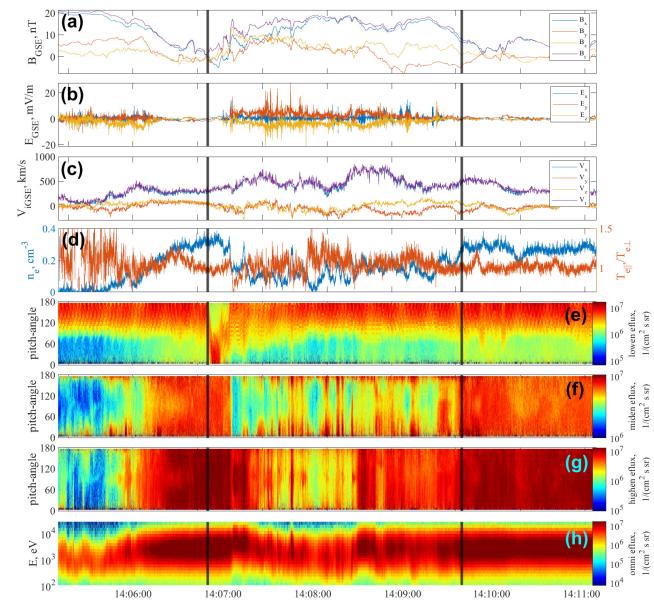
190 The electron acceleration by KAWs is controlled by wave field-aligned electric fields E_{\parallel} , whereas typical accelerated energy is scaled with the wave scalar potential ϕ that can be defined as $E_{\parallel} = -\nabla_{\parallel}\phi$ (Damiano et al., 2015; Artemyev et al., 2015). Note that effective potential describes a combined contributions of electrostatic and electromagnetic fields to E_{\parallel} (Stasiewicz et al., 2000). To compare energies of field-aligned electron population with model predictions of electron acceleration by KAWs, we determine ϕ from E_{\parallel} spectra and KAW dispersion relation (see Eq. (3)). Panel (c) in Figs. 4 shows time series of such reconstructed ϕ that is used to normalize electron energy spectra. Figures 7, 8 show that energies of anisotropic electron 195 population are well within $\sim 3|\phi|$, whereas the most anisotropic electrons have energies below $|\phi|$. This result confirms the scenario of the field-aligned electron acceleration by KAWs.

4 Statistical analysis

In this section we analyze of flux anisotropy and KAWs' electric fields for the entire statistics of 81 events (see Fig. 1). Fig. 200 9 shows averaged wave spectra and ratios of wave spectra for the entire statistics. The ratio $B_{\parallel}/B_{\perp} \sim 1$ for a wide frequency range, i.e. there is some contribution of compressional waves in the low frequency range ($\sim 10^{-2} - 10^{-1}$ Hz where measured compressional magnetic field fluctuations well exceed expectations from KAWs' dispersion relation, see Fig. 9 (c)). Such low frequency compressional fluctuations are quite typical for high-beta plasma sheet, especially during strong plasma flows (Bauer



(a) First example (2017-08-18)



(b) Second example (2018-06-22)

Figure 3. Two intervals with wave activity sharing properties of KAWs. (a): GSE magnetic field vector; (b): GSE electric field vector; (c): GSE ion bulk velocity; (d): electron density and temperature anisotropy, (e)-(g): electron pitch-angle distributions for three energy ranges: (e)[0, 200] eV, (f) [200, 2000] eV, (g) [2, 30] keV; (h): electron energy spectrum. Vertical black lines bound time interval used for wave spectral analysis.

et al., 1995; Zelenyi et al., 2015). Besides flapping waves (Fruit et al., 2002; Volwerk et al., 2004b; Runov et al., 2005), there is a wide range of *magnetic holes*, localized dips of magnetic field magnitude, contributing to compressional fluctuations on 205 the low frequency range (Ergun et al., 2020). Due to Doppler shift of all these low frequency fluctuations, these are well mixed with KAWs and we cannot distinguish KAWs contribution to B_{\parallel} in the low frequency range. However, already at frequency > 0.1 Hz KAWs' contribution dominates and observed ratio B_{\parallel}/B_{\perp} follows theoretical expectations.

The effect of non-KAWs contribution to magnetic and electric field fluctuations at low frequencies is also seen in E_{\perp}/B_{\perp} ratio that is $\sim 2 - 8v_A$ instead of theoretical prediction $E_{\perp}/B_{\perp} = v_A$ (Fig. 9(b)). Transverse electric field fluctuations are 210 stronger than one would expect for Alfvén waves, and these additional fluctuation intensity likely comes from convection of compressional magnetic field structures, i.e. strong plasma flow converts almost stationary magnetic field fluctuations (low frequency flapping waves, magnetic holes) to electric field fluctuations (Volwerk et al., 2003; Zelenyi et al., 2015). For larger frequency (> 0.1 Hz) the ratio E_{\perp}/B_{\perp} follows well the prediction of KAWs dispersion, i.e. transverse electric field fluctuations grows with frequency (dominated by the Doppler shift) faster than magnetic field fluctuations (Chaston et al., 2005, 2012). 215 Thus, Fig. 9 suggests that magnetic and electric fields (for frequency > 0.1 Hz) in selected intervals are dominated by KAWs.

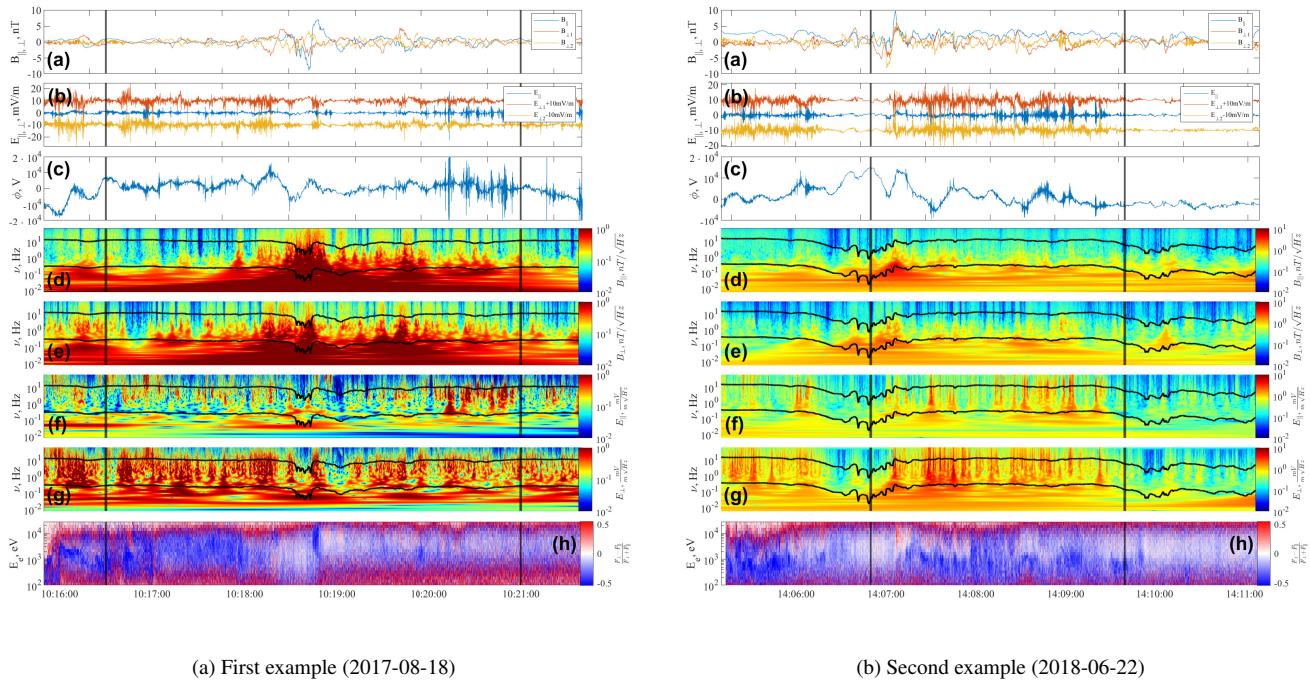


Figure 4. The same intervals as on Fig. 3. (a): wave magnetic field vector in field-align coordinate system; (b): wave electric field vector in field-align coordinate system; (c): electrostatic potential estimated from the equation 3; (d)-(g): magnitudes of wavelet spectra of parallel and perpendicular wave magnetic and electric field powers; (h): electron flux anisotropy. Vertical black lines bound time interval used for wave spectral analysis. Black curves show ion cyclotron and lower hybrid frequencies.

Fig. 10 shows statistical results for electron flux anisotropy, that are quite similar to results shown for individual events in Figs. 7, 8. Energy range (normalized on electron temperature or KAWs' electrostatic potential) of anisotropic electron population increases with the increase of field-aligned and transverse electric field amplitude. This anisotropic population occupies $E/T_e < 1$ energy range, that well coincides with the $E/|\phi| < 1$ range expected from models of electron field-aligned acceleration by KAWs (Damiano et al., 2015; Damiano et al., 2016; Artemyev et al., 2015).

5 Discussion

We show that electron interaction with KAWs is the perspective candidate for generation of the field-aligned anisotropic electron population. This population is often observed in the Earth's magnetotail (Walsh et al., 2011; Artemyev et al., 2014), but its origins remain unestablished. Energies of field-aligned electrons can reach several keV (Sergeev et al., 2001; Yushkov et al., 2017), whereas the ionosphere outflow is expected to be limited to sub-keV energies (Walsh et al., 2013). Thus, some acceleration mechanism is needed for shaping this population. The most popular acceleration mechanism is the electron acceleration by field-aligned electric fields within the reconnection region (Asano et al., 2010; Egedal et al., 2013; Angelopoulos et al.,

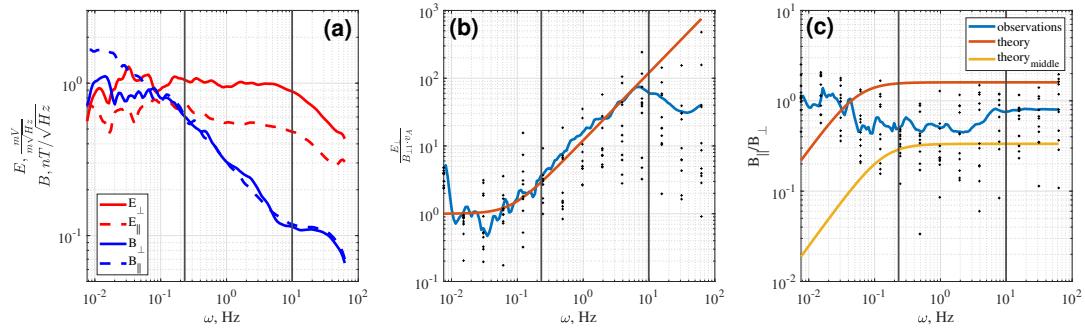


Figure 5. Theoretical and observed wave field spectra and spectra ratios for the first event (2017-08-18) averaged over 10:16:30-10:21:00. (a): averaged over selected time interval wavelet spectra of parallel/perpendicular electric and magnetic field components. (b): averaged over selected time interval ratio of perpendicular electric and magnetic field components wavelet spectra and theoretical prediction. (c): averaged over selected time interval wavelet spectra ratio of parallel and perpendicular magnetic field components and two theoretical predictions: red line shows results for observed plasma parameters and yellow line shows results for plasma parameters estimates at some distance from the equatorial plane (see text for details). Black markers show spread in time on several frequencies. Vertical lines show averaged values of ion cyclotron and lower hybrid frequencies.

2020) but observations suggest that field-aligned anisotropic electrons are not necessarily seen around spatially localized reconnection region (Artemyev et al., 2020). Fast plasma flows are known as primary regions of charged particle acceleration in the magnetotail (Birn et al., 2013, 2014; Ukhorskiy et al., 2017; Gabrielse et al., 2014, 2016), but these flows are mainly associated with compressional transverse electron heating (Khotyaintsev et al., 2011; Zhang et al., 2019). Our results suggest that KAWs associated with fast plasma flows (Chaston et al., 2012) can support the mechanism of the field-aligned electron acceleration. Thus, fast flows are indeed the primary acceleration regions, but field-aligned electron acceleration is not directly resulted from electron interaction with dipolarization fronts (fast plasma flow fronts) and require KAWs' generation and further dissipation due to Landau resonance with electrons (Sharma Pyakurel et al., 2018; Gurram et al., 2020).

KAWs as a natural mode of Alfvén waves in hot ion plasma (Stasiewicz et al., 2000) are suggested to be generated by several different mechanisms, e.g., mode conversion (Hasegawa and Chen, 1975; Hasegawa, 1976; Lin et al., 2012; Johnson and Cheng, 1997), phase mixing (Guo et al., 2015) and magnetic reconnection (Chaston et al., 2005, 2009; Liang et al., 2016; Wang et al., 2019b; Cheng et al., 2020). These waves are believed to be responsible for energy transfer from the magnetotail to

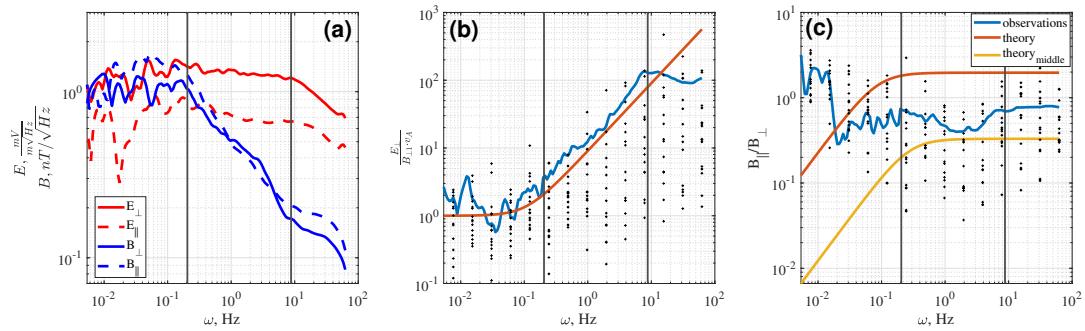


Figure 6. Same as in Fig. 5 for the second event (2018-06-22). Spectra averaged over 14:06:50-14:09:40.

240 the ionosphere (Lysak and Song, 2003; Keiling et al., 2003) and ionospheric electron outflow can be a response for such energy deposition (Lysak, 1990; Keiling, 2009). Thus, KAW electric field correlation with the field-aligned electron population can be, at least partially, due to the outflow mechanism. However, such outflow is expected to be north-south asymmetric (due to asymmetry of ionosphere properties Haaland et al. (e.g., 2017); Laundal et al. (e.g., 2017)) and low-energy (< 100 eV). Indeed, MMS FPI measurements show parallel/anti-parallel asymmetry for low energy electrons for a wider energy range up to one
 245 keV. Presence of accelerated field-aligned anisotropic electrons and good correlation of their energy range with the upper limit of KAWs acceleration, the scalar potential amplitude ϕ , suggests that electron acceleration by KAWs are not less important in production of the electron field-aligned anisotropy than more traditionally considered ionosphere outflow (Walsh et al., 2013).

6 Conclusions

In this study we have collected statistics of 81 fast plasma flow events accompanied by observations of KAWs. Case and
 250 statistical studies show a correlation between the KAWs' electric field magnitudes and thermal/subthermal electron ($\sim 1 - 2$ keV) flux anisotropy: wider energy range of electron anisotropic population corresponds to higher KAWs' electric field intensity. Energies of this population are well within the upper limit of electrons accelerated by KAWs, i.e. KAWs' electrostatic potential. These results indicate an important role of KAWs in production of thermal field-aligned electron population typically observed in the Earth's magnetotail.

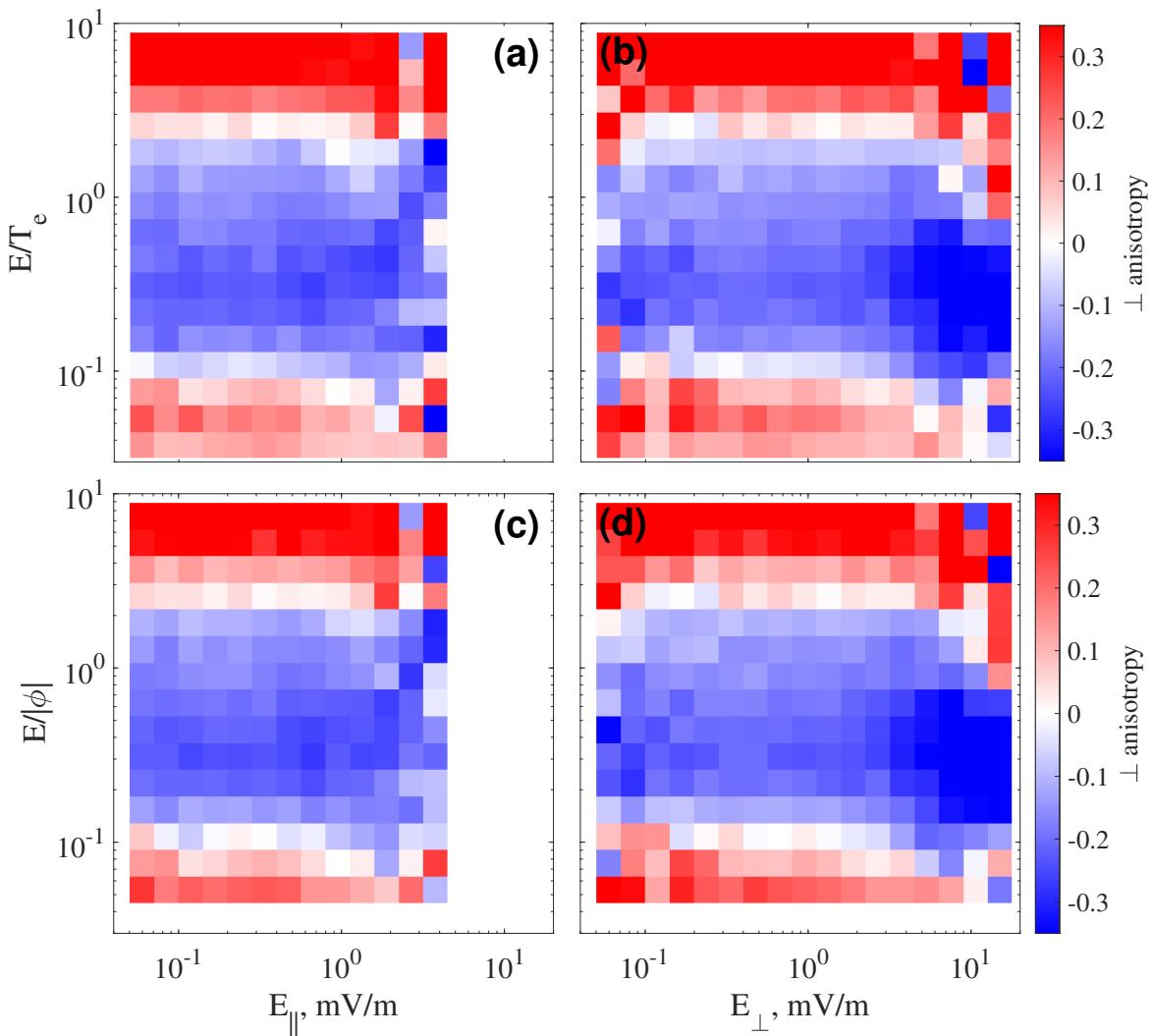


Figure 7. Electron flux anisotropy as function of parallel and perpendicular wave electric fields for the first event (2017-08-18) averaged over 10:16:30-10:21:00. (a), (b): electron energy normalized on the electron temperature; (c), (d): electron energy normalized on absolute value of electrostatic potential, estimated from the Eq. 3.

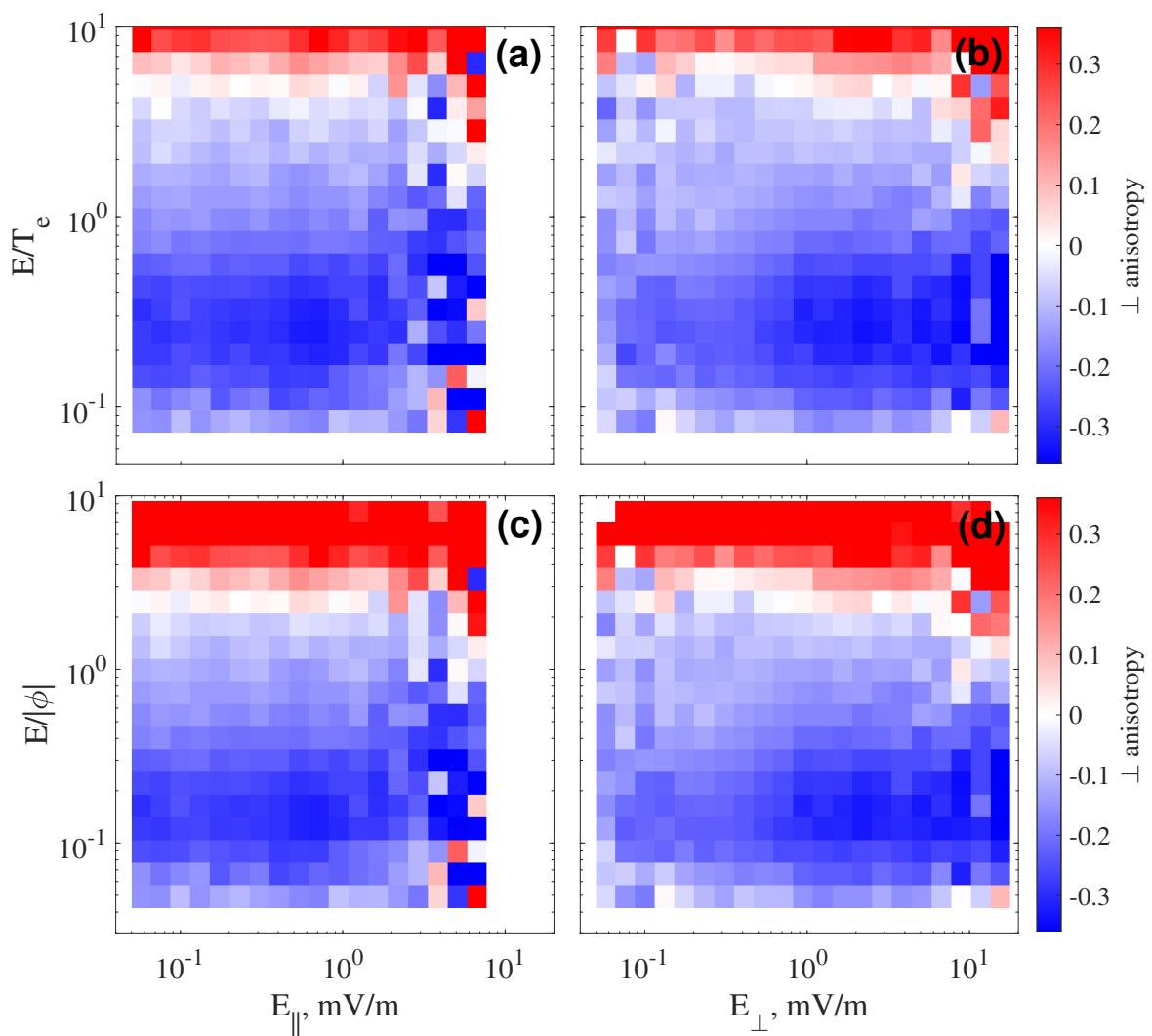


Figure 8. Same as in Fig. 7 for the second event (2018-06-22). Data averaged over 14:06:50-14:09:40.

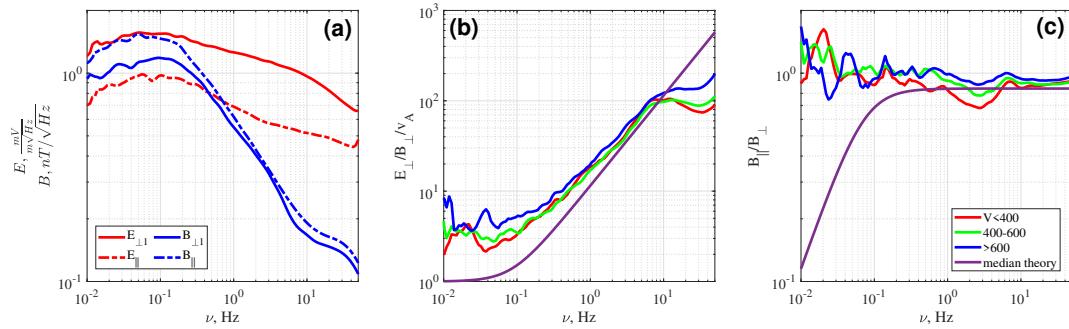


Figure 9. Averaged wave field spectra for the entire statistics. (a): averaged wavelet spectra of parallel and perpendicular electric and magnetic field components. (b): averaged ratios of spectra of perpendicular electric and magnetic field components and theoretical estimates. (c): averaged ratios of spectra of parallel and perpendicular magnetic field components and theoretical estimates.

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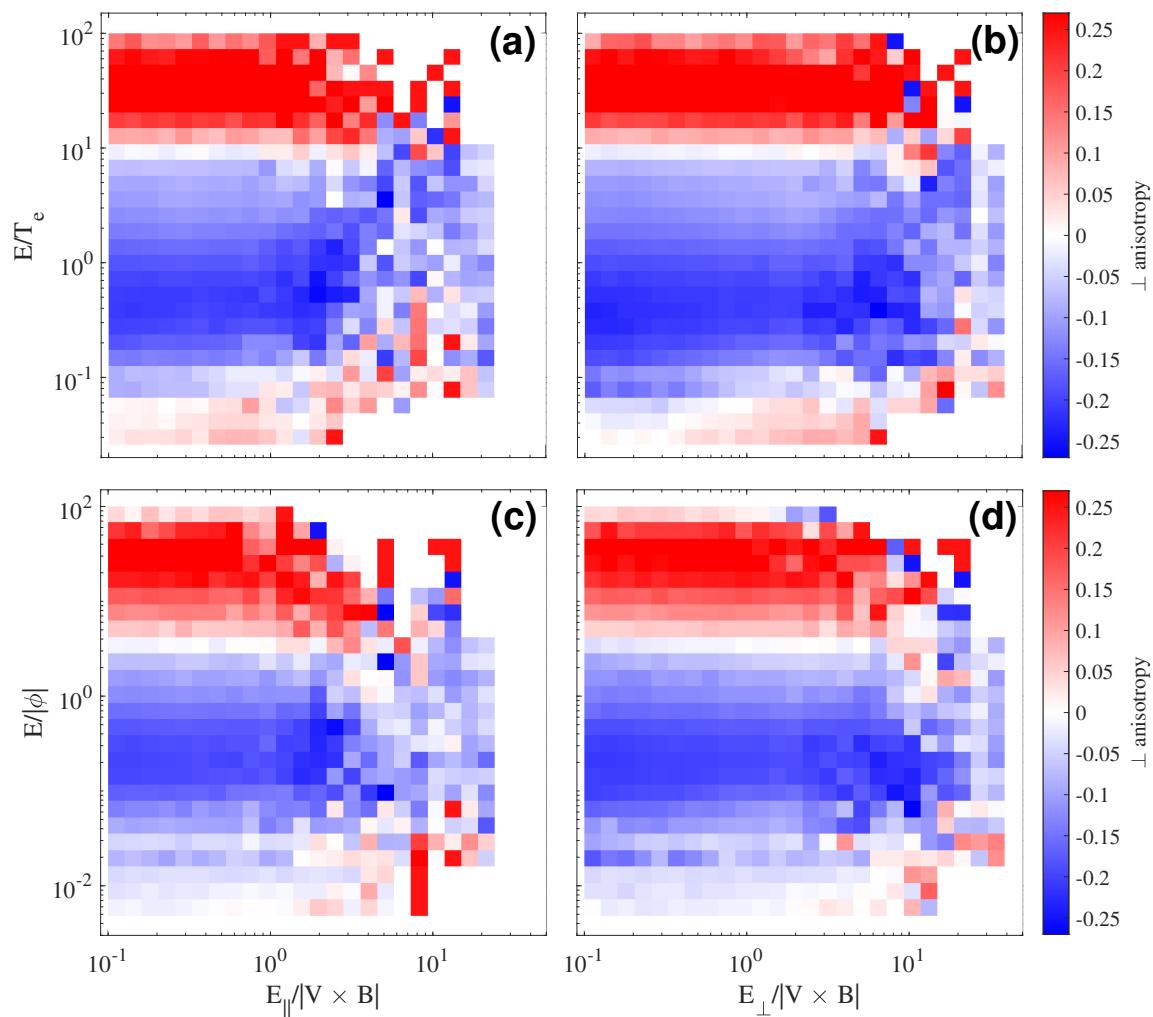


Figure 10. Statistical study of electron flux anisotropy. Figure format repeats the one shown in Fig. 7, but here we use dataset averaged over the entire statistics of 81 events.



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