

1 **Relation between the asymmetric ring current effect and the anti-sunward auroral** 2 **currents, as deduced from CHAMP observations**

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13 **Abstract.** During magnetically active periods the storm-time disturbance signal on ground
14 develops commonly an azimuthal asymmetry. Negative deflections of the magnetic horizontal
15 (H) component are enhanced in the 18:00 local time sector and smallest in the morning sector.
16 This is commonly attributed to the asymmetric ring current effect. In this study we are
17 investigating the average characteristics of anti-sunward net currents that are not closing in the
18 ionosphere. Their intensity is growing proportionally with the amount of solar wind input to the
19 magnetosphere. There is almost twice as much current flowing **across the polar region** in the
20 winter hemisphere as on the summer side. This seasonal dependence is more pronounced in the
21 dusk than in the dawn sector. Event studies reveal that anti-sunward currents are closely related
22 to the main phase of a magnetic storm. Since also the asymmetry of storm-time disturbances
23 build up during the main phase, we suggest a relation between these two phenomena. From a
24 statistical study of ground-based disturbance levels during magnetically active periods we obtain
25 support for our suggestion. We propose a new 3D current system responsible for the zonally
26 asymmetric storm-time disturbance signal that does not involve the ring current. The high-
27 latitude anti-sunward currents are connected at their noon and midnight ends to field-aligned
28 currents that lead the currents to the outer magnetosphere. The auroral net current branch on the
29 morning side is closed along the dawn flank **near the magnetopause**, and the evening side
30 currents flow along the dusk flank magnetosphere. Regardless through which loop the current is
31 flowing, near-Earth storm-time disturbance levels will in both cases be reduced in the morning
32 sector and enhanced in the evening.

33 34 35 **1. Introduction**

36 At auroral latitudes intense electric currents are flowing. Due to the anisotropic conductivity
37 distribution in the ionosphere different current types exist. Quite prominent are the field-aligned
38 currents (FACs), which can transfer energy and momentum over large distances from the

39 magnetosphere and deposit them in the high-latitude upper atmosphere. Horizontal Pedersen
40 currents are typically closing these FACs in the ionosphere. Furthermore, there are Hall currents,
41 flowing perpendicular to the electric and magnetic fields. These are generally regarded as source-
42 free, and they close in the ionosphere.

43 The intensity of currents that close FACs in the ionosphere can be estimated from magnetic field
44 measurements of low-Earth orbit (LEO) satellites on near-polar orbits. By integrating the along-
45 track magnetic field component over the full orbit the net current flowing transverse to the
46 orbital plane can be determined reliably. Corresponding results for anti-sunward net currents
47 have been obtained from Magsat (e.g. Suzuki and Fukushima, 1984). Stauning and Primdahl
48 (2000) used Ørsted magnetic field measurements for estimating the dawn to dusk net currents.
49 Equally, from CHAMP data Zhou and Lühr (2017) could determine the ionospheric net currents
50 for all local times. Net currents increase up to several Mega Ampère (MA) during magnetically
51 active periods. They can be divided into two principle types. Most prominent are the cross-polar
52 cap Pedersen currents closing excessive Region 1 (R1) FACs, which are not balanced by R2
53 FACs. About half as strong are the anti-sunward net currents connecting excessive downward
54 FACs on the dayside with upward FACs on the nightside. These anti-sunward currents, carried
55 predominantly by Hall currents, have first been confirmed observationally from Magsat data
56 (Suzuki and Fukushima, 1982, 1984). **Their intensity, derived from the ring integral of the
57 along-track field component, is clearly controlled by magnetic activity. Later Yamashita et
58 al. (2002) used a somewhat different approach. They interpreted the azimuthal, By
59 component of the Ørsted magnetic field data at middle and low latitudes for estimating
60 FACs flowing into and out of the ionosphere. These authors also deduced anti-sunward net
61 current intensities, dependent on magnetic activity, from their data. The advantage of this
62 latter approach is that it can also be applied to ground-based observations (see Nakano
63 and Iyemori (2005) and references therein). But the disadvantage is that important
64 assumptions have to be made for the interpretation in terms of net currents. Strangely, the
65 more intense dawn to dusk net currents are obviously not sensed by this approach.**

66 More recently Zhou and Lühr (2017) provided a detailed study on auroral zone net currents.
67 Making use of 5 years of high-resolution CHAMP magnetic field data, they could, for all local
68 times, derive the dependence of these currents on season, solar wind input and solar flux. In
69 particular, by estimating currents separately for the two hemispheres these dependences emerged
70 very clearly. The cross-polar cap duskward net current peaks at local summer when the
71 ionospheric conductivity is high. Conversely, the anti-sunward net current attains largest values
72 during local winter when conductivity gradients between the auroral region and the polar cap

73 maximise. At these gradients Hall currents can be diverted into FACs. The out-of-phase variation
74 of these two current types causes quite different responses of net current intensities in the two
75 hemispheres to magnetic activity.

76 There are still a number of open questions about the relationship between auroral zone net
77 currents and the asymmetric storm-time disturbances during the main phase. Suzuki and
78 Fukushima (1984) proposed a closure of the net anti-sunward current through the duskside
79 partial ring current. **Conversely, Crooker and Siscoe (1981) argued that the magnetic signals
80 from the excessive FACs around noon and midnight are sufficient to explain the
81 asymmetry signal, but they did not tell anything about current closure in the
82 magnetosphere. Ground-based measurements of the magnetic field eastward component at
83 mid and low latitudes have been used for estimating anti-sunward net currents (for a
84 review see Iyemori, 2000). This author offers several options for magnetospheric return
85 current on the duskside located somewhere between the magnetopause and the ring
86 current.** Furthermore, it has never been investigated how the anti-sunward net current flow is
87 split between the dawn and dusk side auroral regions. What is the effect of hemispheric
88 differences in current strength due to seasonal variation? Can a detailed consideration of all these
89 facts provide hints on the actual 3D geometry of the net anti-sunward current closure in the
90 magnetosphere?

91 The C/NOFS satellite on its low-inclination orbit can be used to investigate the ring current
92 asymmetry. On every revolution it samples ring current signals from all local times. Magnetic
93 field readings of C/NOFS during the years 2008 through 2010 have been considered by Le et al.
94 (2011) to study the ring current evolution during storms. The authors show that the disturbance
95 signal is azimuthally symmetric before and after the storm. But during the main phase a clear
96 asymmetry is building up, with enhanced amplitudes around the 18 LT sector and reduced values
97 around 06 LT. During the storm recovery phase, the disturbance signal returns to symmetric
98 distribution. The degree of asymmetry grows as the magnetic activity gets larger, but the local
99 time sector in which the largest amplitudes are observed stays around 18 LT. Similar results
100 concerning the asymmetry of the ring current effect have been derived from ground-based
101 observations (e.g. Love and Gannon, 2009). These authors claim that the dawn-dusk asymmetry
102 in the disturbance field is on average proportional to D_{ST} . Newell and Gjerloev (2012) made use
103 of a large number of magnetometers from the SuperMAG data repository. Their SMR index is
104 similar to D_{ST} but provides local time resolution with four sectors (SMR-00, SMR-06, SMR-12,
105 SMR-18). By means of a superposed epoch analysis Newell and Gjerloev (2012) determined the
106 response of their index to a magnetic storm. They found a clear dominance of the disturbance

107 signal at 18 LT and smallest deflections at 06 LT. All this is consistent with the notion of a
108 partial ring current on the duskside. For checking that inference Lühr et al. (2017) had a look at
109 *in situ* ring current density measurements by Cluster and other spacecraft. They could not
110 confirm the enhancement of ring current intensity in the dusk sector. The strongest ring current
111 parts during a magnetic storm are rather observed by these missions in the post-midnight sector.
112 The difference in ring current interpretation from near-Earth observations and *in situ*
113 measurements has been described in more details by Lühr et al. (2017), but it is still an open
114 issue.

115 In this study we make use of CHAMP data and follow up on the results presented by Zhou and
116 Lühr (2017) for addressing the open questions listed above. Of special interest is the relation
117 between the net anti-sunward current and the asymmetric storm-time effect at low latitudes.
118 Prime basis for the investigations is the CHAMP magnetic field dataset from the 5 years, 2001-
119 2005. But also recordings from geomagnetic observatories are taken into account for
120 characterizing the near-Earth magnetic effects.

121 In the sections to follow we will first shortly introduce the data and basic processing algorithms
122 for determining net currents. Section 3 presents a statistical survey of net currents at all local
123 times. The dependence of anti-sunward net currents on solar wind input and season is analysed
124 in Section 4. Section 5 presents for one magnetic storm a direct comparison between anti-
125 sunward currents and ground-based disturbance levels. The mean characteristics of the ring
126 current signal during magnetically active periods ($K_p > 6$), as observed on ground, are outlined
127 in Section 6. In Section 7 the various observations are discussed, focusing on the comparison
128 between anti-sunward currents and storm-time disturbance signals. Finally, in Section 8 results
129 are summarised and a new 3D current system is proposed for closing the anti-sunward net
130 currents.

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132 **2. Dataset and calculation of net auroral currents**

133 The CHAMP satellite was launched into a near-circular polar orbit (inclination: 87.3°) with an
134 initial altitude of 456 km on 15 July 2000 (Reigber et al., 2002). By the end of the mission, 19
135 September 2010, the orbit had decayed to 250 km. The orbital plane covers all local times within
136 130 days when considering upleg and downleg arcs. The Fluxgate Magnetometer (FGM) on
137 board CHAMP recorded the vector magnetic field every 0.02 s with a resolution of 0.1 nT. The
138 FGM magnetic field readings are calibrated routinely by using the observations of the onboard
139 absolute scalar Overhauser Magnetometer. In this study the fully calibrated Level-3 magnetic

140 field products (product identifier: CH-ME-3-MAG) are used (Rother and Michaelis, 2019),
 141 which are provided in the North-East-Center (NEC) frame with a time resolution of 1 Hz. The
 142 time period used in this study comprises the five years from 2001 to 2005, experiencing solar
 143 and magnetic activities from high to moderate levels. Five years of CHAMP magnetic field
 144 observations are just needed to sample all local times 14 times, evenly distributed over all
 145 seasons.

146 The approach for deriving net currents in the auroral region from CHAMP magnetic field data
 147 has been described in detail by Zhou and Lühr (2017). Here we use the same dataset and adopt
 148 their processing algorithm. Calculations are based on Ampère's law in integral form

$$149 \quad I = \frac{1}{\mu_0} \oint_L B_{AT} dl \quad (1)$$

150 where I is the net current flowing through the closed integration contour, μ_0 is the permeability
 151 of free space, B_{AT} is the along-track magnetic field component caused by the current I , dl is a
 152 differential path element along the CHAMP orbit. Equation (1) can be written in discrete form
 153 as

$$154 \quad I = \frac{1}{\mu_0} \sum_{m=1}^n B_{ATm} \cdot \Delta l \quad (2)$$

155 where m is the summation index, and Δl is the path length per increment (here 7.56 km for 1s).
 156 For deriving the along-track magnetic field component, B_{AT} , we have subtracted from the
 157 CHAMP data the main field, crustal field and large-scale magnetospheric field, as represented
 158 by the high-resolution model POMME-6 (Maus et al., 2010). From the set of magnetic
 159 residuals, the component aligned with the velocity vector is calculated.

160 Zhou and Lühr (2017) derived net currents from integration along full CHAMP orbits. In
 161 addition, they applied integration loops confined to one hemisphere and could study
 162 hemispheric differences. Here we go one step further by estimating net currents flowing through
 163 a loop from subauroral latitudes up to the geomagnetic pole. In this way we get current estimates
 164 for all local times and can compare directly net current intensities on the dawnside with those
 165 on the duskside and noon with midnight results. The penalty for the further detailing of the
 166 results is that we have to make certain assumptions on the magnetic fields along parts of the
 167 integration path where no direct observations are available. The considered integration paths
 168 for the two local time sectors along the orbit are sketched in Figure 1. CHAMP magnetic field
 169 readings are taken from 50° magnetic latitude (MLat) (point A) up to the highest MLat reached

170 along the orbit (point B). From there the virtual return path goes vertically down to point C,
171 follows the Earth's surface until point D and goes vertically up to the start point A. The second
172 loop follows the same scheme, taking CHAMP readings along the track from E to F and closing
173 the loop along the virtual path (F-G-H-E).

174 Since there are no measurements along the return path, we have to make assumptions about the
175 magnetic field along that track. Here we follow the same reasoning and approach as
176 successfully applied in the work of Zhou and Lühr (2017). Auroral net currents are connected
177 to FACs on both ends. According to Fukushima's theorem (Fukushima, 1976) magnetic
178 signatures from a pair of antiparallel FACs closed by ionospheric currents vanish at the Earth's
179 surface. The current configuration in our case, however, differs somewhat from the ideal case
180 presented by Fukushima (1976), therefore the theorem might not be fully applicable here.

181 For estimating the contributions from the unsampled parts the following assumptions are made:
182 (1) The contribution from C→D is similar in shape to that from A→B. (2) The contribution
183 from D→A are proportional to the vertical field component B_z at point A since the radial
184 magnetic field varies only smoothly through the current sheet. **We have modelled E-region**
185 **currents flowing along the auroral oval connected to FACs on the noon and midnight**
186 **ends. Resulting magnetic signals along a closed path as outlined in Figure 1 were**
187 **calculated.** An outcome of this exercise is that the integral over A→B has to be multiplied by
188 1.2 for including the contributions from path C→D and that the vertical magnetic field
189 component, B_z , has to be multiplied by 11 times the orbital altitude and divided by the
190 permeability of free space to represent the contributions from path D→A. For further validation
191 of these corrections see Zhou and Lühr (2017), Section 4.2. The same approach described here
192 is also applied to the contour E-F-G-H-E. The remaining paths in the integration loops are B-C
193 and E-H. Here again, the observed B_z component at the point B (E) has been taken as a measure
194 for scaling the missing contribution. We have tested a series of different factors multiplied to
195 the B_z value at the top-side corners. There is a statistical way to validate the suitability of the
196 applied factors. Each local time sector is sampled in two ways, on upleg and 130 days later on
197 downleg passes. In these two groups the ring integral is calculated in opposite directions. Only
198 in the case of a proper scaling of this vertical contribution, both results are on average identical.
199 From this test we found that the best agreement is obtained when the contributions from the
200 vertical path elements in the middle are neglected. Figure 2 shows the final comparison for both
201 hemispheres and all local times. For the northern hemisphere (left frame) we obtain, when
202 ignoring the vertical paths, an almost perfect match between upleg and downleg results. The

203 agreement is not as good for the southern hemisphere, but any additional contribution from this
204 vertical path element makes the agreement between the curves worse. **Our resulting**
205 **assumption of insignificant contributions from the vertical path elements at the poles does**
206 **not affect the total net current flowing over a polar region. It may just affect the**
207 **partitioning of anti-sunward currents between the dawn and dusk sides.**

208

209 **3. Statistical survey of net current distribution**

210 For obtaining the average distribution of net currents at all local times we consider CHAMP
211 magnetic field data from the 5 years, 2001-2005. Overall 24,440 orbits with clean data are
212 available. From each orbit we obtain two net current results for both hemispheres. This results
213 in a large number ($\sim 10^5$) of samples for this study. Figure 3 shows the average local time
214 variations of net currents in the northern and southern hemispheres (upleg and downleg results
215 are combined). Positive values represent eastward currents. On average we find somewhat
216 larger values in the northern hemisphere than in the southern. This is consistent with the
217 observations of Zhou and Lühr (2017). Positive (eastward) net currents prevail within the local
218 time sector 07-19 MLT, representing a dawn to dusk flow. The opposite sign is found around
219 the 24 MLT sector, reflecting also dominant dawn to dusk currents.

220 There is not only a local time variation of the net currents but also a dependence on season.
221 Figure 4 shows the distribution of current strength in a magnetic local time (MLT) versus Month
222 of Year frame. We clearly find strongest currents during local summer months in particular
223 around noon at both hemispheres. This is primarily due, as explained by Zhou and Lühr (2017),
224 to the enhanced ionospheric conductivity during that season.

225 As outlined by Zhou and Lühr (2017), the large net currents derived from noon/midnight
226 orbits can be related to the cross-polar cap Pedersen currents closing the excessive Region 1
227 (R1) FACs. The positive values around noon and the negative around midnight are both
228 consistent with that notion. In this study we are more interested in the net currents on the
229 dawn and dusk sides. Therefore, we consider the average values from orbits within the local
230 time sectors 03-09 MLT and 15-21 MLT as dawnside and duskside net currents, respectively.
231 From Figure 3 it is evident that a negative (westward) average current results from the 03-09
232 MLT sector and a positive (eastward) from the 15-21 MLT sector. This means, both sides
233 contribute to an anti-sunward net current. The characteristics of these anti-sunward currents
234 are of prime interest for this study.

235

236 4. Dependence of net current on solar wind input and on season

237 Similar to Zhou and Lühr (2017) we also investigate the dependence of anti-sunward net
238 currents on magnetic activity. Different from them we look at the fractions flowing on the
239 dawn and dusk sides separately. As measure for the solar wind input, we use the coupling
240 function as defined by Newell et al. (2007). By somewhat rescaling this function we obtain
241 the so-called merging electric field, E_m , which represents approximately the solar wind
242 electric field in units of mV/m

$$243 \quad E_m = \frac{1}{3000} V_{SW}^{\frac{4}{3}} (\sqrt{B_y^2 + B_z^2})^{\frac{2}{3}} \sin^{\frac{8}{3}}\left(\frac{\theta}{2}\right) \quad (3)$$

244 where V_{SW} is the solar wind velocity in km/s, B_y and B_z both in nT are the IMF components
245 in GSM coordinates, θ is the clock angle of the IMF. E_m values have been smoothed over 15
246 min, and the propagation time from the bow shock to the ionosphere has been considered by a
247 delay of 20 min (for more details see Zhou and Lühr, 2017).

248 Figure 5 shows the mean dependence of the eastward net currents on the dawn and dusk sides
249 on the merging electric field, E_m , separately for the northern and southern hemispheres. The
250 current values had been grouped into five activity classes ($0 < E_m \leq 1$, $1 < E_m \leq 2$, $2 < E_m \leq 3.5$,
251 $3.5 < E_m \leq 5$, $5 < E_m \leq 7$ mV/m). Blue dots represent the mean values within these classes and the
252 blue bars reflect the standard deviations. The mean values infer a good linear relationship
253 between current intensity and merging electric field in all cases, as confirmed by the fitted red
254 lines. On the dawnside westward currents get stronger with growing E_m and correspondingly
255 eastward currents intensify on the duskside. This confirms in all four cases an increase of anti-
256 sunward currents with growing activity. Slopes are somewhat steeper on the dawnside than on
257 the duskside. Interestingly, the net currents on the dawnside show a small positive bias (~ 52
258 kA) for vanishing solar wind input. We relate that to the effect of net anti-sunward plasma
259 flows driven by intense day-to-night winds in the early morning sector (e.g. Lühr et al., 2007)
260 during very quiet periods.

261 As expected, the net currents on the flanks depend also on season. Figure 6 shows the mean
262 annual variation of eastward net currents on the dawn and dusk sides separately for the two
263 hemispheres. **Vertical bars represent the formal uncertainty of the mean value for each**
264 **month.** This analysis is based on data from more active periods with $E_m > 3$ mV/m
265 (approximately $Kp > 4^+$) since anti-sunward net currents are phenomena increasing with
266 magnetic activity. We find in both hemispheres weaker anti-sunward currents in the summer
267 hemisphere than at local winter. This holds for the dawn and dusk sides and is consistent with

268 the results of Zhou and Lühr (2017). Compared to the mean values, the relative annual
269 variations are not too large (15% - 20%) and have comparable sizes in both hemispheres. In
270 the northern hemisphere a semi-annual signature is quite prominent, commonly referred to as
271 the Russel-McPherron effect (Russel and McPherron, 1973). It reflects the typical annual
272 variation of magnetic activity with maxima at equinoxes and a minimum around June solstice.
273 The semi-annual variation is not so obvious in the southern hemisphere, but the annual
274 amplitude is larger.

275 For completeness we have also calculated the dependence of the dawn and dusk side net
276 eastward currents on solar wind input separately for June and December solstice months and
277 for the two hemispheres. Obtained results are listed in Table 1. The negative signs of the
278 slopes on the dawnside and the positive on the duskside represent both increasing anti-
279 sunward current intensity with enhanced solar wind input. When comparing the slopes of the
280 dawn and dusk sides between the two solstices, one finds a smaller seasonal difference on the
281 dawnside than on the duskside. At dusk the factor is partly reduced to less than a half during
282 local summer with respect to local winter. Net currents in the dawn sector are obviously less
283 dependent on sunlight in the ionosphere. This is consistent with Guo et al. (2014), who report
284 that the eastward auroral electrojet intensity shows a larger seasonal variation (stronger in
285 local summer) than the westward jet. Finally, it is interesting to note that in Table 1 the
286 intercepts on the dawnside show systematically large sunward net currents (82 kA) in the
287 summer hemispheres. This is consistent with the stronger day-to-night wind in the sunlit polar
288 region (e.g. Lühr et al., 2007) which seem to control the anti-sunward plasma flow over the
289 dawnside polar cap during quiet times.

290

291 **5. Variation of net currents during a magnetic storm**

292 It is suggested since quite some time that the anti-sunward currents are connected via FACs to
293 the ring current (e.g. Suzuki et al., 1985). In particular, it is believed that net currents feed the
294 partial ring current on the duskside. Here we want to check, to which degree the CHAMP data
295 support this inference. The partial ring current generally forms during the main phase of a
296 magnetic storm.

297 For investigating these connections in more details, we have selected the geomagnetic storm
298 on 17 August 2003. This event is well suited because CHAMP is crossing the auroral oval on
299 orbits close to dawn/dusk. The storm is initiated by a sudden storm commencement (SSC) at
300 14:20 UT on 17 August. From the solar wind and interplanetary magnetic field (IMF)

301 variations, shown in Figure 7 (bottom), we can deduce that a sudden increase of solar wind
302 speed from about 420 km/s to more than 500 km/s is responsible for the SSC. About an hour
303 later, when IMF B_z turns negative, the main phase of the storm starts and extends into the
304 next day. On that day the storm time disturbance index reached a minimum of $D_{ST} = -148$ nT
305 (see Fig. 7, top frame). It follows a typical recovery phase lasting several days. During part of
306 that time IMF B_z is still negative, but the solar wind speed has returned to pre-event levels.

307 For comparison we present in the top frame of Figure 7 the storm-time evolutions of the total
308 anti-sunward net currents (blue curves), including contributions from both hemispheres,
309 together with the SYM-H index (red curves). The SYM-H values are averages over the 10-
310 min intervals when CHAMP crossed the polar regions. Right after the southward turning of
311 IMF B_z intense anti-sunward currents (negative values) commence. About 4 hours later
312 currents recover to a moderate value, but intensify again early next morning. This intermittent
313 occurrence of net current continues into the recovery phase of the storm but with decreasing
314 amplitudes.

315 So far, we have seen the evolution of total net current intensity during the magnetic storm on
316 17 August. More details can be derived from Figure 8, where the contributions from the two
317 hemispheres are shown separately. The current signatures are quite different in the four
318 sectors. Before the SSC net currents in all frames are close to zero. Particularly intense anti-
319 sunward currents, up to 2 MA, appear in the southern hemisphere (SH) on the dawnside
320 during the main phase. Some hours before this strong signal, less intense anti-sunward
321 currents are observed on the dawnside in the northern hemisphere (NH) and the duskside SH.
322 It is interesting to note that there is in general a synchronous variation of net currents in these
323 two antipodal sectors with somewhat smaller amplitudes in the south. For example, the
324 prominent negative peaks around 42h Event Time (ET) in both hemispheres, which occur at
325 the start of the recovery phase. Even later in the recovery phase (~55h ET) a sizable anti-
326 sunward current appears in the SH dawn sector. Different to the other sectors there is only
327 little net current activity on the NH duskside. Quite common for all four sectors, there is
328 hardly any net current activity during times of northward IMF.

329 For the interpretation of the observations we have to remind that the event takes place towards
330 the end of northern summer. More intense anti-sunward currents are therefore expected in the
331 SH. Also, the quietness on the NH duskside is consistent with our previously shown statistical
332 results for that season. The quasi-synchronous variation of net currents at NH dawn and SH
333 dusk could convincingly be explained with a control by IMF B_y on related FACs in the polar

334 cap. Stronger anti-sunward currents are expected in the NH dawnside for negative IMF B_y
335 and in the SH dawnside for positive IMF B_y . A direct comparison with the IMF observations,
336 shown in Figure 7, reveals a qualitative agreement. For example, the intense SH dawn current
337 matches well the positive excursion of IMF B_y around 30h ET, but the details of phasing do
338 not fit so well in other cases. At least for this event we can state that in both hemispheres
339 more intense anti-sunward net currents are observed on the dawnside than on the duskside.

340 It would have been desirable to study more individual storms in this detail. But an event has
341 to satisfy a number of conditions for providing instructive results on the temporal evolution of
342 anti-sunward currents during a storm. The storm should occur close to one of the solstice
343 seasons, and the local time of the CHAMP orbit has to be close to dawn/dusk. We have
344 considered all storms during the CHAMP era (2000-2010) with D_{ST} exceeding -100 nT. Just
345 the presented event satisfied all these requirements reasonably well.

346

347 **6. Ground-based signature related to anti-sunward net current**

348 The observed anti-sunward currents are connected on both ends to FACs. These field-aligned
349 currents have to close somewhere in the magnetosphere. Depending on the route these
350 currents take corresponding magnetic signatures are expected at Earth surface. Traditionally
351 the D_{ST} index (or SYM-H, as shown in Figs. 7 and 8) is used for describing the evolution of a
352 storm. But this index reflects only the azimuthally symmetrical part of the magnetospheric
353 fields. Therefore, it is not well suited for quantifying the asymmetric effects, possibly caused
354 by the auroral net currents. More appropriate for this purpose seems to be the SuperMAG ring
355 current index, SMR. It is a quantity comparable to D_{ST} or SYM-H but provides local time
356 resolution from four sectors (SMR-00, SMR-06, SMR-12, SMR-18). More details about the
357 SMR index can be found in Newell and Gjerloev (2012). By comparing the evolution of
358 magnetic signatures on the evening and morning sides (SMR-18 and SMR-06) we may see
359 the effect of a partial ring current. Figure 9 shows in the top frame the field deflections in
360 these two time sectors during our storm. As expected, there are larger amplitudes observed on
361 the evening side, in particular towards the end of the main phase. In the lower frame the
362 differences between the two traces, SMR-18 minus SMR-06, are plotted. In this way we try to
363 eliminate the contribution of the symmetrical ring current. Before and after the active phase of
364 the storm the difference stays close to zero. Shortly after the SSC we find first positive
365 deflections, i.e. a dominance of the dawn sector, and around the end of the main phase the
366 prominent minimum, i.e. larger effects on the duskside. Thereafter the difference signal is

367 more variable. A closer comparison between SMR difference signal and the net currents in
368 Figure 8 reveals that the best (but not perfect) match is found with the CHAMP SH dawnside
369 currents. However, these should, according to the traditional picture, weaken the ring current
370 on the morning sector. At least for this storm the asymmetric D_{ST} effect cannot be explained
371 by an intensification of the duskside ring current. We will revisit this issue in the Discussion,
372 Section 7.

373 From our study of the anti-sunward net currents we know that the effects can be significantly
374 different in the two hemispheres mainly depending on the season. Although SMR provides
375 information on local time differences, it does not distinguish between hemispheric sources.

376 In order to obtain more information on the net current seasonal effects in ground observations
377 we analysed magnetic field data from a meridional chain of observatories. Stations involved
378 are Wingst (WNG, 54.15° DLat), L'Aquila (AQU, 42.45° DLat), Tamanrasset (TAM, 24.80°
379 DLat), Bangui (BNG, 4.36° DLat), and Hermanus (HER, -33.86° DLat), where DLat is the
380 latitude in dipole coordinates. Our study has shown that net currents are particularly strong
381 during magnetic storms. We are therefore interested in magnetic field deflections at the
382 observatories during disturbed times. The disturbance signal is determined from times with a
383 magnetic activity index $K_p \geq 6$. Here the values around 06 and 18 MLT are considered since
384 they are expected to show the largest difference. For studying them we had a look at the
385 hourly averages of the horizontal component, H, from 03-06 UT and 15-18 UT, respectively.
386 A quiet-time background field is subtracted, determined from hourly averages of the same UT
387 times as above, but only data within the $K_p = 0-1$ range are selected. In order to make the
388 result well comparable with our net currents we considered the same 5 years (2001-2005) as
389 for CHAMP.

390 The obtained mean horizontal disturbance fields **are shown in Figure 10** separately for the
391 three Lloyd seasons: June solstice (May-Aug), December solstice (Nov-Feb) and combined
392 equinoxes (Mar+Apr, Sep+Oct). As expected for such active periods, we get negative mean
393 values (southward fields) in all the cases. The values in the evening sector are more negative
394 than those from the morning sector. An exception makes the station WNG. Here the fields on
395 the duskside are more positive, opposed to the other observatories, than those from dawn.
396 This observatory is located obviously too far north. Therefore, its readings are affected also
397 by the auroral electrojet during severe storms, not only by the ring current. For that reason, we
398 have not considered it any further in the analysis.

399 The larger amplitudes at dusk than on the dawnside are traditionally attributed to the effect of
400 the partial ring current. For a more quantitative evaluation of the asymmetry we
401 calculated the mean values of the hours 03+04 UT, representing the dawn levels and the
402 means of 16+17 UT for the dusk values. These UT periods take into account the
403 longitudes of the stations. The obtained values of H component deflections in the two
404 time sectors are shown for the considered observatories in Figure 11, separately for the
405 three seasons. As can be seen, the expected dawn/dusk difference is consistently
406 observed at all stations up to AQU. In addition, numerical values are listed in Table 2.
407 The overall largest negative disturbance fields are obtained for the months around
408 December and smallest around June. This reflects the seasonal distribution of strong
409 storms during the 5 years considered.

410 Here we are more interested in the asymmetry of the disturbance. The mean values of
411 $H_{\text{dawn}} - H_{\text{dusk}}$ considering all seasons are: BNG: 71.3 nT, TAM: 71.4 nT, HER: 63.3 nT,
412 AQU: 49.2 nT. The differences decrease with the distance from the geomagnetic
413 equator. However, these morning/evening differences derived vary from season to
414 season. In Table 3 the values are listed separately for seasons. Clearly largest
415 asymmetries result for June solstice months. This is surprising because the relatively
416 small negative H deflections around June compared to those of December months (see
417 Table 2) indicate stronger storms in the latter season. For an explanation of this
418 apparent inconsistency we may have a look at the magnetic activities prevailing during
419 the relevant periods. It has to be realized that the deflections on the dawn and dusk sides
420 are not measured by our single chain of observatories simultaneously. As a consequence,
421 our selection criterion ($K_p \geq 6$) is commonly fulfilled only in one time sector during a
422 day. By the statistical approach we hoped that variations in magnetic activity average
423 out. Table 4 lists the mean a_p values for the three seasons and two local times. For
424 equinox conditions the activity levels in the two sector match well, but this is not the case
425 for the solstice seasons. During months around June, duskside measurements are from
426 clearly larger activity periods than the corresponding dawnside samples. In that case a
427 stronger ring current effect at dusk will add to the asymmetry and therefore cause the
428 enhanced differences in Table 3. Just the opposite scenario is true for the events around
429 December solstice. Here the reduced ring current activity on the duskside compared to
430 dawn reduces the actual asymmetry effect. As a consequence, we have to state, our
431 ground-based observations are not sufficient to reveal a possible seasonal effect of the
432 storm-time disturbance asymmetry. The average result (2001-2005) of the ground-based

433 **observations reveals a mean disturbance field asymmetry of about 72 nT at the equator**
434 **results for a weighted mean magnetic activity level of $a_p = 134$ nT ($K_p \sim 7$).**

435 The ring current signal has also been measured by the C/NOFS satellite. On its low-latitude
436 orbit (inclination: 13°) it samples the H component deflections at all local times on every
437 orbit. In that way, any azimuthal asymmetries of the signal can well be detected. In a
438 dedicated study, Le et al. (2011) investigated the evolution of the ring current signals during
439 several geomagnetic storms. They clearly could confirm the appearance of an asymmetry
440 during the storm main phase. During the recovery phase the signal became symmetric again.
441 In a later study Lühr et al. (2017) performed a statistical survey on the type of asymmetry. For
442 different classes of magnetic activity, the mean difference between dawn and dusk deflections
443 were determined and the local time where the maximum appeared. For high activity, $K_p > 6$,
444 they obtained a center displacement of 38 nT, half the difference between dawn and dusk
445 signals. This can be compared with the difference of disturbance levels that we derived here
446 for periods of $K_p > 6$ from the observatories. We obtained a mean value of 72 nT near the
447 equator, which is slightly less than the corresponding result from C/NOFS (76 nT). **An**
448 **explanation for the difference between the two values could be our averaging over two**
449 **hours of dawn and dusk side measurements, while in case of C/NOFS the actual**
450 **minimum and maximum values are compared. Overall, the two independent types of**
451 **measurements confirm each other including the level of disturbance asymmetry.**

452

453 **7. Discussion**

454 In this study we investigated the statistical properties of anti-sunward net currents in the
455 auroral regions and their relation to ground-based signatures at middle and low latitudes. The
456 general properties of auroral net currents had been presented by Zhou and Lühr (2017). Here
457 we go one step further by determining the anti-sunward currents flowing on the dawn and
458 dusk sides separately.

459

460 **7.1 Dependence on season and solar wind input**

461 As expected, the net current intensity is directly proportional to the solar wind coupling
462 function, E_m . **This has earlier been reported (e.g. Nakano and Iyemori, 2005).** When
463 looking at annual averages the resulting net currents are about the same for enhanced activity
464 (e.g. $E_m > 3$ mV/m) in the dawn and dusk sectors and in both hemispheres (see Fig. 6).
465 However, obvious differences appear when taking the local seasons into account. From Table

466 1 we can deduce that the slopes of the current intensity curves with respect to E_m are similar
467 on the dawn sides for local summer and local winter. Conversely on the duskside, the obtained
468 E_m dependences are clearly steeper for winter than for summer conditions. This is valid for
469 both hemispheres. We interpret it as an indication that the conductivity gradient on the
470 duskside between the auroral region and the polar cap is much steeper in the winter
471 hemisphere than in the sunlit summer. Different from that dawn side conductivity gradients
472 seem to be less season dependent.

473 When evaluating the average hemispheric net current characteristics from Table 1 we obtain
474 for $E_m = 6$ mV/m ($Kp \approx 6^+$) intensities of about 640 kA and 810 kA in each hemisphere for
475 summer and winter conditions, respectively. It has been reported earlier (e.g. Guo et al., 2014)
476 that the intensity of the eastward electrojet on the duskside is depending more directly on the
477 sun-induced conductivity. Obviously, the stronger summer-time eastward electrojet
478 contributes less to the anti-sunward net currents. The closure of those electrojet currents
479 across the polar cap seems to be quite efficient during the sunlit season.

480 A detail, interesting to note, is that for vanishing solar wind input, $E_m = 0$, i.e. due northward
481 IMF, we obtain, in particular on the dawn side during summer season, sunward net currents of
482 about 80 kA in both hemispheres. Reason for that is probably the day-to-night wind over the
483 polar cap that is driving **anti-sunward plasma drift, overcoming the dawn to dusk electric**
484 **field effect and causes** net currents in opposite direction. More dedicated studies would be
485 needed for elucidating the details of a high-latitude wind dynamo under such special
486 conditions.

487

488 7.2 Comparison with ground-based observations

489 We have shown that the magnetic field effects of anti-sunward currents are also observable on
490 ground. Our satellite results imply that the asymmetry between dawn and dusk disturbance
491 signals during magnetically active periods should be larger in the winter hemisphere, and the
492 seasonal effect more prominent at mid-latitude ground stations than near the equator.

493 **However, our statistical study of recordings from a single European-African meridional**
494 **chain is not sufficient to confirm the seasonal difference between hemispheres. It would**
495 **require at least two meridional chains separated by about 180° in longitude for**
496 **monitoring storm-time disturbances on the morning and evening sectors simultaneously.**

497 **Nakano et al. (2002) had deduced anti-sunward net currents from eastward magnetic**
498 **field deflections at mid-latitude stations around noon and midnight. They report,**

499 **consistent with our satellite results that ground-based signals are larger in the winter**
500 **hemisphere than in the sunlit hemisphere. Since the asymmetric storm-time disturbance**
501 **signal is expected to result mainly from the connected field-aligned currents, mid-**
502 **latitude stations in the same hemisphere are predominantly affected by it.** This implies,
503 recordings in the summer hemisphere underestimate during active periods the D_{ST} value
504 because of the northern hemisphere station dominance. An over-proportional reduction of the
505 mean D_{ST} index during months around June solstice, compared to other activity indices, e.g.
506 K_p , has earlier been reported (e.g. Mursula and Karinen, 2005). In their Figure 1 they show
507 that the average H component deflections at the northern hemisphere index observatories
508 reaches almost 0 nT at the beginning of July. While at Hermanus the zero level is attained
509 around New Year. In our view this northern hemisphere D_{ST} minimum can be explained by
510 the combined effect of the well-known annual July magnetic activity minimum with the
511 weaker disturbance signal in the summer hemisphere. **At Hermanus this July minimum is**
512 **much less prominent, but therefore December, January values are reduced.** Since three
513 out of four D_{ST} observatories are located in the northern hemisphere, the excessive summer
514 minimum in D_{ST} is expected to result from a hemispheric bias. Just for completeness we may
515 note that Mursula and Karinen (2005) offered another explanation for the D_{ST} July minimum
516 which we do not regard as so convincing.

517 Rather interesting features are revealed from the event study of the magnetic storm on 17
518 August 2003. The evolution of sunward currents, as shown in Figure 8, is quite different on
519 the dawn and dusk sides in the two hemispheres. Several of the statistical features presented
520 in the previous sections can also be found in this event that occurred during northern summer
521 conditions. Largest currents are detected in the southern, winter hemisphere on the dawnside
522 during the storm main phase. In the northern summer hemisphere, the duskside currents
523 exhibit only small amplitudes. This is consistent with the mean seasonal dependences of this
524 local time sector (see Table 1). Sizable net currents appear on the dawnside in the northern
525 hemisphere at times when they are low in the southern hemisphere. This hemispheric
526 alternation in current flow can be related to the varying direction of the IMF B_y component.

527 For checking the magnetic effects of the net currents on ground we had a look at the SMR
528 index for this event (see Fig. 9). We expected a clear dominance of SMR-18 over SMR-06.
529 But only a moderate negative difference appears towards the end of the main phase in the
530 lower frame of that figure. Over large parts of the storm-time the signal is varying about the
531 zero-line. For the interpretation of this result we have to note that most of the observatories

532 contributing to the SMR index are located in the northern hemisphere. Because of the
533 prevailing summer season, the asymmetry is expected to be underestimated.

534 There is a certain anti-phase variation of the SMR difference in Figure 9 with the sunward
535 currents in Figure 8 on the NH dawnside and SH duskside. Prominent peaks appear around
536 19h and 41h ET in both figures but with opposite sign. This indicates that at the listed peak
537 times the negative deflections in the northern hemisphere are stronger on the dawnside than
538 on the duskside. The largest negative peak in the SMR difference signal, around 30h ET, is
539 well aligned with the strong anti-sunward current on the SH dawnside, but it is not as large as
540 expected from the strong SH net current deduced from CHAMP data. This observation
541 provides clear evidence that the effect of auroral net currents can be recognised by the
542 asymmetry of mid-latitude observatory readings, but the two hemispheres should be
543 interpreted separately. **It may be more instructive to have separate asymmetry values**
544 **from the SuperMAG stations for the northern and southern hemispheres.** With the
545 present distribution of stations, contributing to SMR, it is expected that this effect is
546 underestimated around June solstice and overestimated during December months.

547

548 7.3 Suggestion for a 3D current circuit

549 When comparing the CHAMP net currents at the four quadrants with the temporal evolution
550 of the SYM-H or SMR indices we find strongest net currents in the dawn sector and
551 particularly in the southern, winter hemisphere (see Fig. 8) during the storm main phase. The
552 traditional suggestion was that the auroral net currents, in particular those from the evening
553 sector, are connected to the ring current and intensify the part in the dusk sector (e.g. Suzuki
554 et al. 1985). But just on the duskside we find only weak anti-sunward currents during our
555 August 2003 storm. In previous works the term “partial ring current effect” is frequently used.
556 This was mainly meant as an acronym for an azimuthally asymmetric disturbance signal
557 during magnetic storms (e.g. Iyemori, 2000). The presented observations in this paper and
558 previous publications considering *in situ* ring current density distributions (see Lühr et al.,
559 2017 for a review) provide little evidence for a direct connection between auroral net currents
560 and the ring current. Here we want to introduce our idea of the 3D current circuit connected
561 with the anti-sunward currents.

562 From electrodynamic considerations it can be assumed that the FACs on the nightside are
563 connected to the net currents at steep conductivity gradients. This locates them at fairly high
564 latitudes near the border between auroral oval and polar cap. Field lines from this border do

565 not connect to the ring current but reach out close to the magnetopause. During the storm
566 main phase a lot of current flows along the electrojets from the day to the night side, which
567 cannot be returned to the dayside across the poorly conducting polar cap (in particular in the
568 dark hemisphere). The excessive current flows out along field lines to the outer magnetosphere
569 on the dawn and dusk side flanks.

570 Figure 12 presents a schematic drawing of the envisaged 3D current circuit. Shown is a view
571 onto the northern hemisphere. Equivalent current routes are assumed on the southern side.
572 There is no connection to the ring current foreseen.

573 **For the field-aligned currents flowing on the dayside into the ionosphere we assume that**
574 **they originate from a dynamo region in the low-latitude boundary layer (LLBL). In a**
575 **comprehensive review Lundin (1988) describes important properties of the**
576 **magnetospheric boundary layer. In his Sections 6 and 7 he outlines dynamo action and**
577 **the connection between the LLBL dynamo and the ionosphere. Following an injection of**
578 **magnetosheath plasma into the LLBL, due to reconnection, the initially existing plasma**
579 **at rest is accelerated tailward, which will set up polarization electric fields in the**
580 **dynamo region. As a consequence, FACs are flowing into the ionosphere. Dependent on**
581 **the orientation of IMF B_y , the injection takes place before or after local noon. Already**
582 **Bythrow et al. (1981) noticed from observations the excessive FACs flowing, besides the**
583 **Region 1 and Region 2 systems, into the ionosphere near noon. In our schematic picture,**
584 **Figure 12, the LLBL is depicted by grey shading and the dynamo regions are indicated**
585 **by the bulges around the Earthward FACs. Current closer is envisioned through the**
586 **LLBL from the tail region to the dayside.**

587 A current flowing through our dawnside circuit will generate a northward magnetic field on
588 ground, thus reducing the D_{ST} effect. Conversely, net currents through the dusk loop cause a
589 southward field, enhancing the ring current effect. Regardless on which side the net currents
590 close, the same kind of asymmetry results. The near-Earth disturbance signals at middle
591 latitudes from these current circuits are dominated by the magnetic effects of the connecting
592 FACs. For the resulting asymmetries, it does not make a big difference at which distance in
593 the magnetosphere the currents close, in the ring current or further out near the magnetopause.

594 With the 3D current circuit suggested here, it makes no problem to understand, why enhanced
595 disturbance levels always appear around 18 MLT (see Le et al., 2011) independent of the
596 magnetic activity level. Already Love and Gannon (2009) had noticed that storm-time
597 disturbances are commonly higher around the 18 MLT sector. They even suggested a linear

598 relation between the asymmetry amplitude and the D_{ST} value. The asymmetry should amount
599 on average to about 20% of the D_{ST} value. This claim was challenged by Siscoe et al. (2012).
600 These authors tried to identify a magnetospheric process that could systematically enhance the
601 ring current intensity in the dusk sector. In the end they were not able to offer a convincing
602 explanation.

603 We claim that our 3D current circuit, driven by **plasma injection through magnetic**
604 **reconnection on the dayside**, can better explain the observed features of the asymmetry
605 signal. It seems to be a quite stable circuit in space. Therefore, the localisation to 18 MLT,
606 independent of activity is achievable. We do not believe in a dependence of the asymmetry
607 amplitude on the D_{ST} value. But in a statistical sense, E_m and D_{ST} are related, therefore the
608 result of Love and Gannon (2009) can be explained. More correlated studies of magnetic
609 fields and currents in the outer magnetosphere and near-Earth observations are needed to
610 confirm our 3D current configuration.

611

612 **8. Summary and Conclusions**

613 In this study we have investigated the auroral net currents flowing anti-sunward. For the first
614 time, we present the partitioning of contributions from the dawn and dusk sides and from the
615 two hemispheres to the total net current. These magnetic storm-time phenomena show
616 significant dependences on solar wind input, season, and IMF B_y orientation. Of particular
617 interest here is the complete current circuit including the field-aligned currents attached to the
618 anti-sunward currents and the closure in the magnetosphere. Important results may be
619 summarised as follows:

- 620 1. Anti-sunward currents grow on average proportionally with the solar wind input (merging
621 electric field, E_m). This is valid for the dawn and dusk sides and for all seasons.
- 622 2. More intense currents are observed in the winter than in the summer hemisphere. We
623 relate that to the steeper conductivity gradients between auroral zone and polar cap during
624 dark seasons. In the winter hemisphere a larger part of the electrojet return current has to
625 be by-passed through the magnetosphere via FACs.
- 626 3. The seasonal dependence of net currents is significantly larger on the dusk than on the
627 dawn side. In the sunlit summer hemisphere the anti-sunward current intensity in that
628 sector is greatly reduced compared to its value during winter conditions (see Table 1). On
629 annual average, more anti-sunward current is flowing on the dawnside (10%-20%).

- 630 4. Event studies of magnetic storms confirm the connection between anti-sunward auroral
631 currents and the asymmetric storm-time disturbance signal. From the event studied we see
632 that this claim holds for the total net current. But the partitioning of current through the
633 different loops can change during a storm several times between dawn and dusk sides and
634 the two hemispheres. Responsible for the preferred path is the prevailing season and the
635 IMF *B_y* orientation.
- 636 5. We propose a 3D current system **causing the asymmetric storm-time disturbances that**
637 **is driven by reconnection-related plasma injections on the dayside. Earthward**
638 **directed field-aligned currents around noon feed the anti-sunward high-latitude net**
639 **currents, and around midnight FACs carry the currents into the outer**
640 **magnetosphere on the tailside. A closure of the loops is anticipated by currents**
641 **through the low-latitude boundary layers** on the dawn and dusk flanks. We do not find
642 evidence for a connection of this circuit with the ring current.

643 For confirming our claims about the large-scale current system causing the asymmetric storm-
644 time magnetic disturbances more observations in the outer magnetosphere should be analysed.

645

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651 available at ftp://spdf.gsfc.nasa.gov/pub/data/omni/high_res_omni/. The SMR data are
652 downloaded from <http://supermag.jhuapl.edu/indices/>. The ground observations of magnetic
653 field are available at www.intermagnet.org. The work of Yun-Liang Zhou is supported by the
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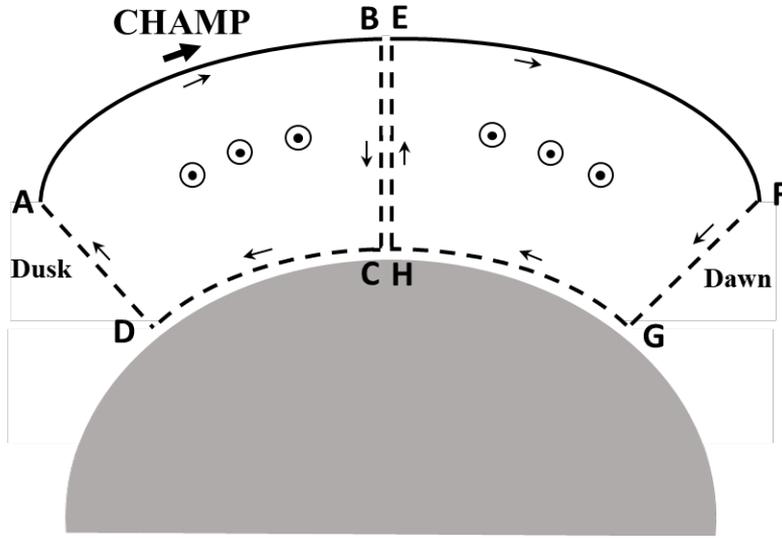
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Figures

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735

736 **Figure 1.** Schematic drawing of the anti-sunward net current (small circles with dots)

737 **determination by the ring-integral approach at auroral latitudes, separately for the**

738 **dawn and dusk local time sectors. Small arrows indicate the direction of integration.**

739 **The unsampled but estimated virtual return paths are shown as dashed lines.**

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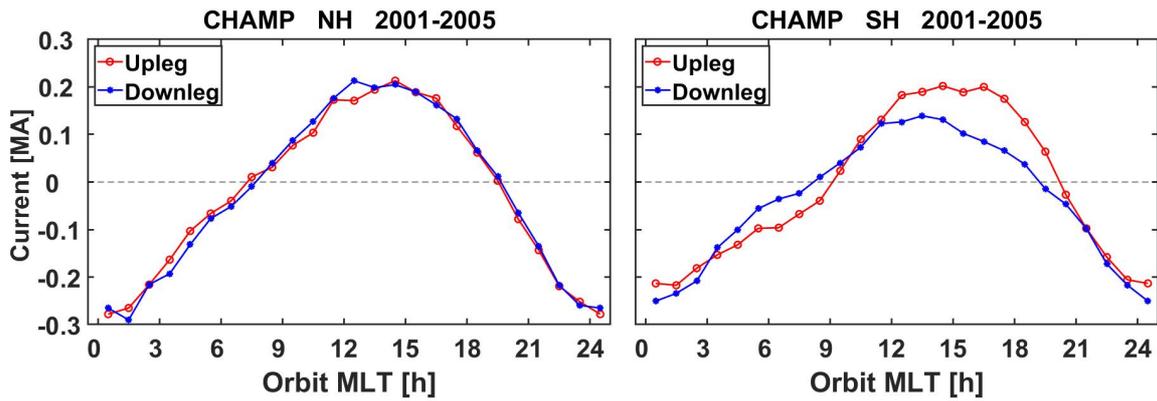
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749 **Figure 2.** Orbital local time dependence of auroral net currents separately for results

750 **from upleg and downleg passes. Best matches, shown here, are obtained when the**

751 **contributions from the vertical paths B-C and H-E are neglected.**

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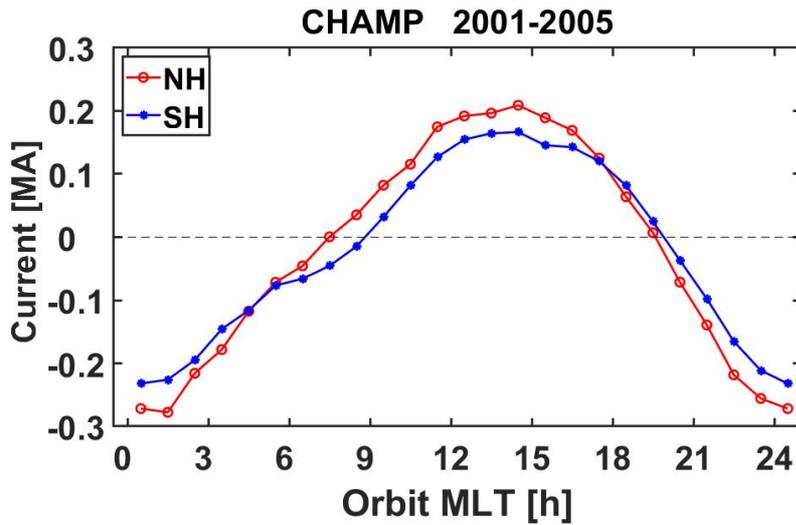
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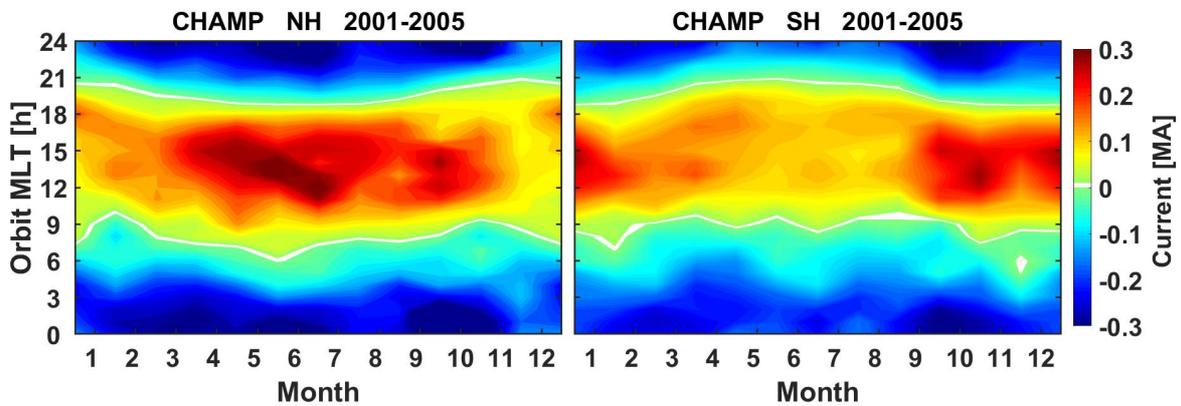
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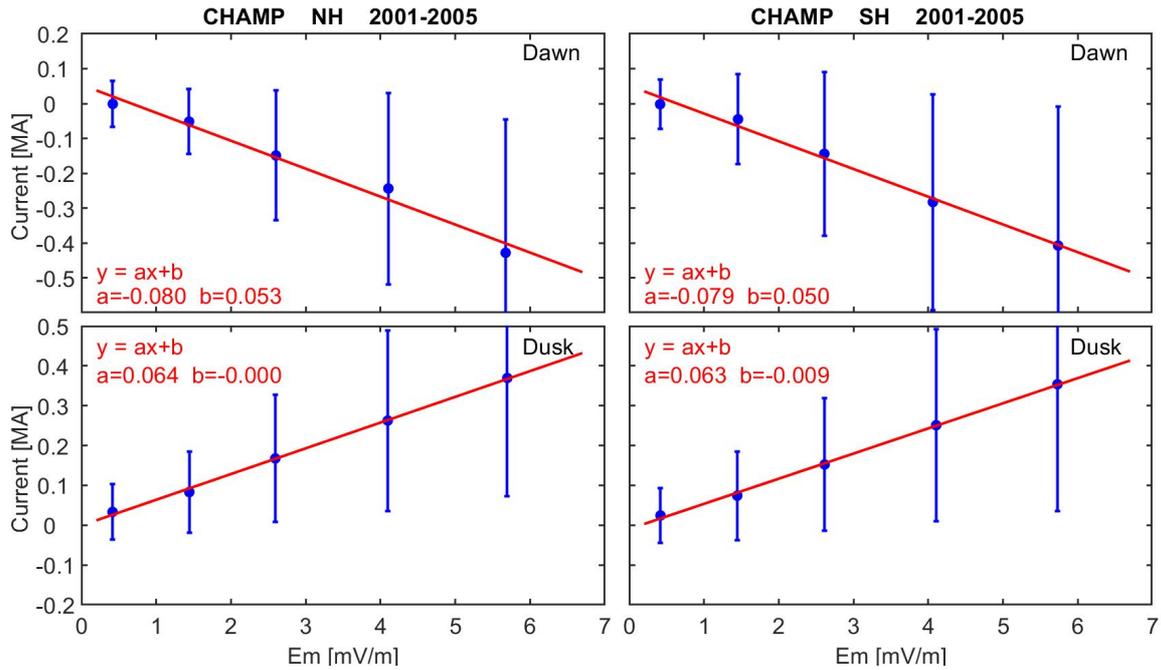
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Figure 3. Orbital local time dependence of mean auroral net currents; comparison between the two hemispheres.



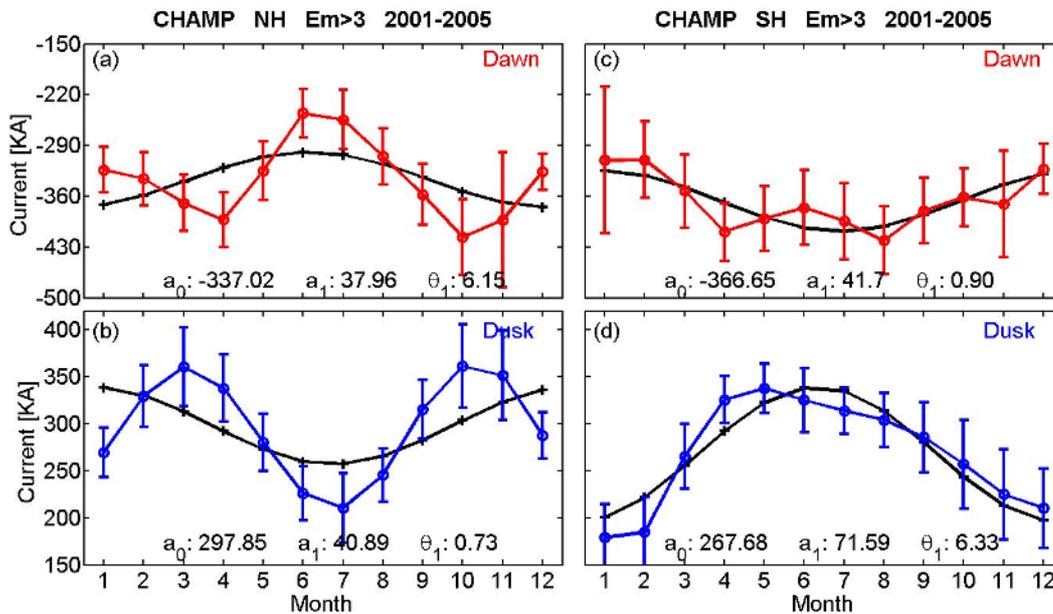
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Figure 4. Distribution of mean eastward net currents in local time versus Month of Year frames. Currents from noon-time orbits are strongest during local summer season in the two hemispheres.



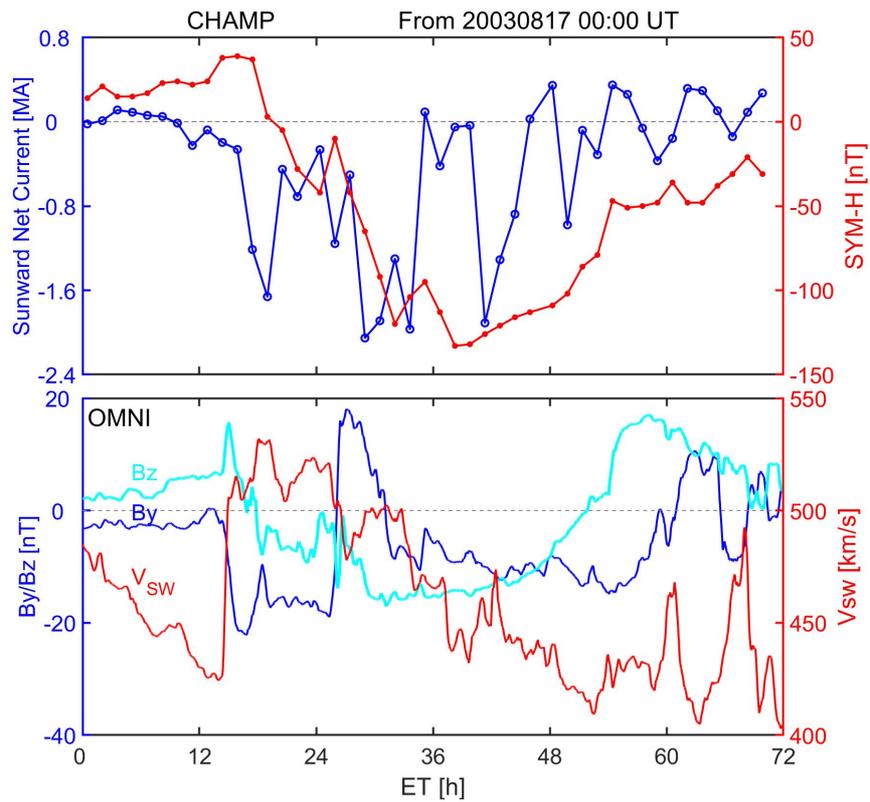
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Figure 5. The merging electric field, E_m dependence of net currents on the dawn and dusk sides, separately for the Northern (left) and Southern (right) hemispheres. The solid dots with vertical bars indicate the mean values and standard deviation of the net eastward current for five levels of E_m . Parameters of the linear fits (red lines) are listed in the top left corner of each frame.



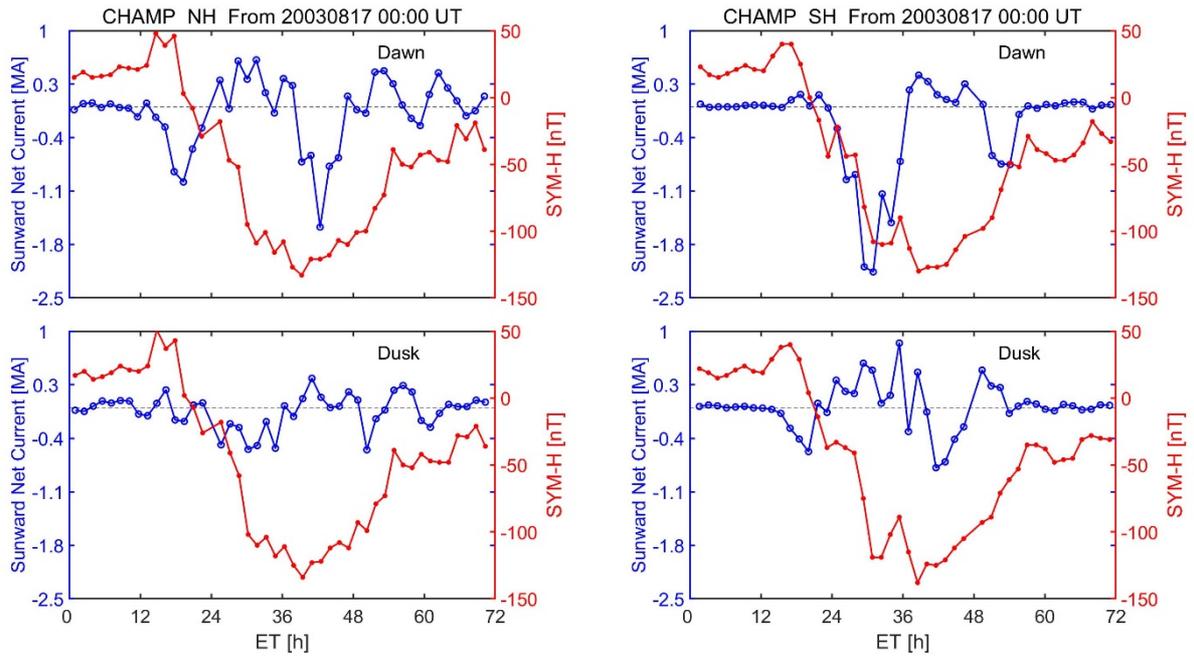
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Figure 6. The seasonal variation of eastward net currents. Presented are dawnside (top) and duskside (bottom) currents derived from high-latitude passes over the Northern (left) and Southern (right) Hemispheres. **Vertical bars represent the uncertainty of the monthly means.** Black curves are sinusoidal fits to the observations. In each panel the constant term, a_0 , annual amplitude, a_1 (both in kA) and the phases the peaks, θ_1 , (in month) are listed.



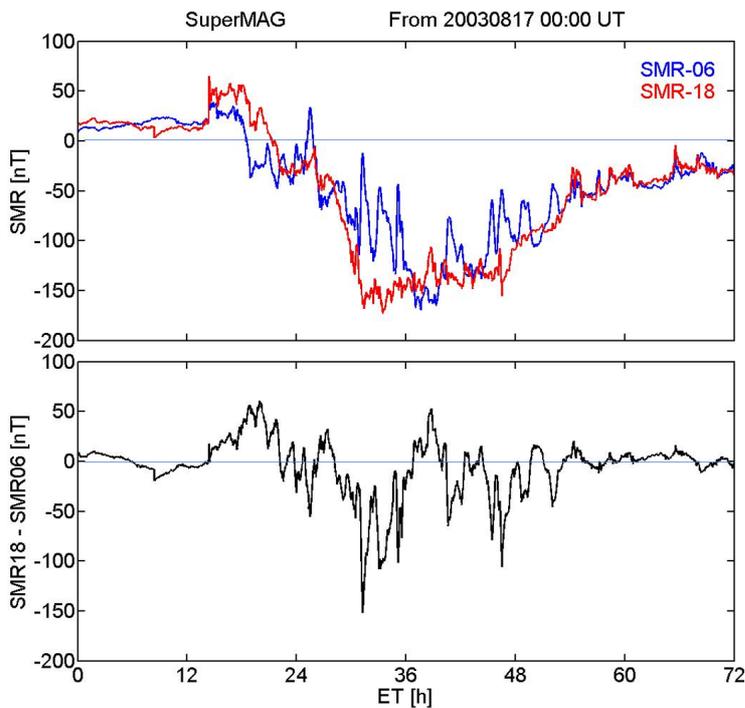
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Figure 7. (bottom) Solar wind velocity and interplanetary magnetic field components (GSM) variations for the storm starting on 17 August 2003. (top) The SYM-H index evolution during the storm and the total anti-sunward net current are shown for comparison.



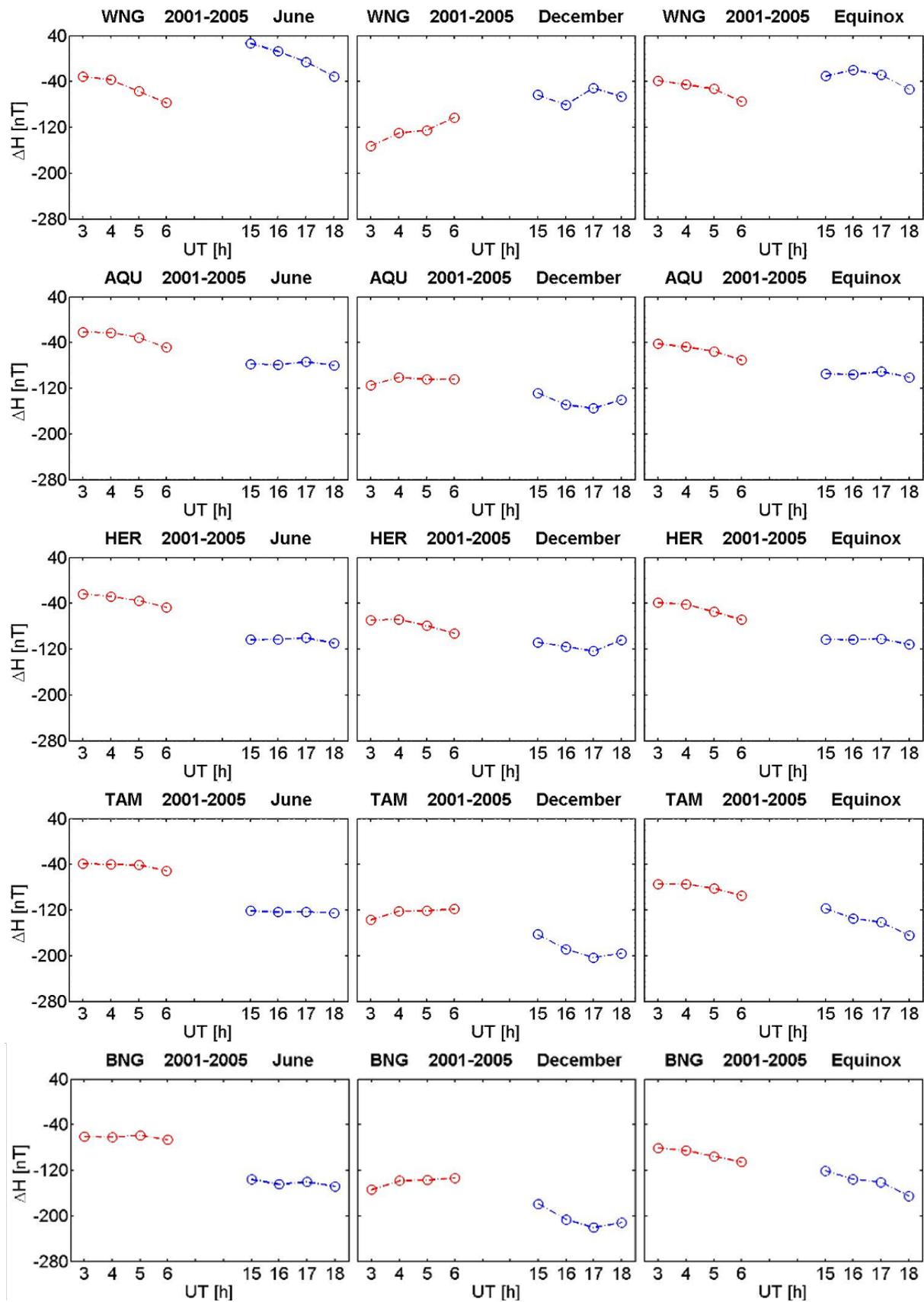
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Figure 8. Temporal evolutions of the SYM-H index and the net currents separately for both hemispheres and for dawn and dusk sides during the storm 17-20 August 2003. Magnetic local time ranges are, NH dawn passes: 06-09 MLT, NH dusk passes: 17-21 MLT, SH dawn passes: 03-10 MLT, SH dusk passes: 17-23 MLT.



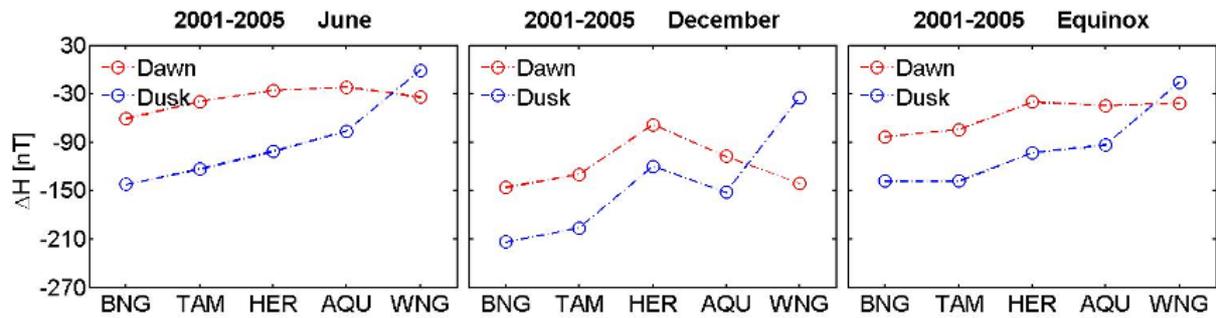
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Figure 9. (top) Temporal evolution of the SMR storm-time index from the 06 and 18 local time sectors. (bottom) Differences between the two time sectors (SMR-06 – SMR-18).



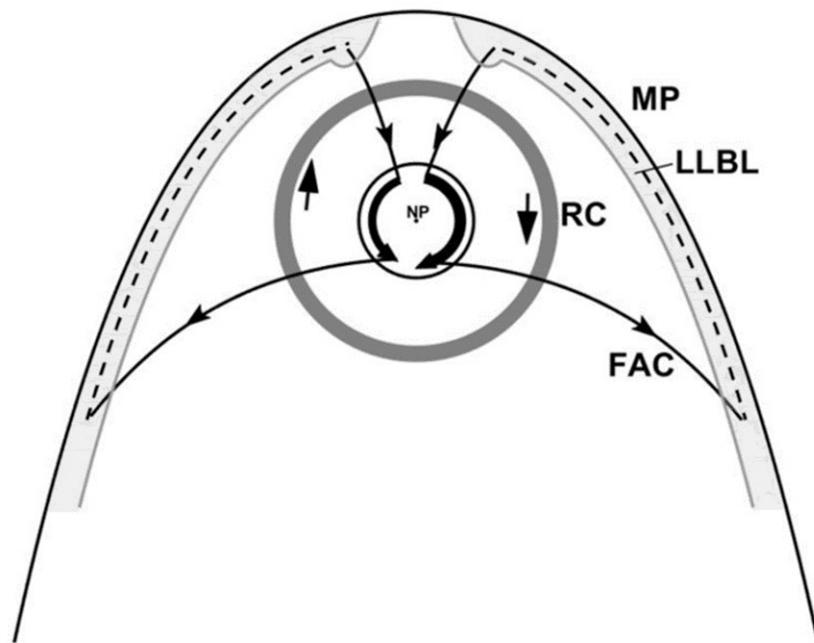
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Figure 10. Mean deflections of the H component in the dawn and dusk sectors at five considered observatories separately for different seasons during active times ($K_p \geq 6$).



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Figure 11. Comparison of the mean H component deflections between the dawn and dusk sectors at the five observatories separately for different seasons during active times ($K_p \geq 6$).



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Figure 12. Schematic drawing of the suggested 3D current circuits causing the storm-time disturbance asymmetries. Field-aligned currents flowing out of the dynamo region on the dayside feed the anti-sunward net currents in the polar region. On the nightside FACs lead the currents into the outer magnetosphere on the dawn and dusk flanks. Here the currents are assumed to flow sunward within the low-latitude boundary layer (LLBL) for closing the loops. Equivalent current circuits are expected in the southern hemispheres.

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Table 1. The E_m dependence of the net eastward currents during June and December solstice months for both the dawn and dusk sides.

Season	Local time sector	Northern Hemis.		Southern Hemis.	
		Slope (10^6Am/V)	Inters. (kA)	Slope (10^6Am/V)	Inters. (kA)
Months: 05-08	Dawn	-78	83	-75	44
	Dusk	49	16	73	-18
Months: 11-02	Dawn	-69	26	-80	82
	Dusk	67	8	29	16

Table 2. The mean deflections of H component (in nT) at five observatories for different seasons during active times ($K_p \geq 6$)

Station	Local time	June	December	equinox
WNG	Dawn	-34.2	-140.8	-41.8
	Dusk	-1.1	-34.7	-15.9
AQU	Dawn	-22.0	-107.3	-44.7
	Dusk	-76.4	-151.9	-93.3
HER	Dawn	-25.6	-68.3	-40.1
	Dusk	-101.6	-119.5	-102.7
TAM	Dawn	-39.4	-129.5	-74.3
	Dusk	-123.3	-196.0	-138.0
BNG	Dawn	-61.0	-146.2	-72.9
	Dusk	-142.3	-213.6	-138.2

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Table 3. Mean dawn – dusk differences of the H component (in nT) at four observatories for different seasons during active times ($K_p \geq 6$)

Station	DLat	June	December	Equinoxes
BNG	4.36°	81.3	67.4	65.3
TAM	24.81°	83.9	66.5	63.7
HER	-33.86°	76.0	51.2	62.6
AQU	42.45°	54.4	44.6	48.6

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Table 4. Mean a_P values (in nT) of the times considered for the study, separately for the three seasons and the two local times.

Local time/Season	June	December	Equinoxes
Dawn	110.8	154.5	129.3
Dusk	157.0	127.7	126.3

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