# **Reply to referee comments**

Wavevector spectral signature of decay instability in space plasmas

## H. Comişel, Y. Narita, and U. Motschmann

We thank very much the reviewer for carefully reading our manuscript and for providing valuable suggestions and comments. Here we give our reply. The changes in the manuscript are marked in boldface.

#### Reviewer:

The manuscript is presenting the results of analytical model describing parametric decay to explain earlier published results of 3D numerical simulations by the same authors. The manuscript requires some improvements before being publishable. Points to be addressed:

\* MAJOR. It is not clearly specified what are the new results in comparison with Viñas and Goldstein (1991) results and other recent studies. The reader should clearly see what has been done until now and what is the new addition to the field.

#### Reply:

In the two-dimensional analysis of Viñas and Goldstein 1991, the authors concluded the decrease of the oblique-decay growth rates into narrow band-profiles (in the wavenumber domain) when moderate oblique propagation angles of the daughter waves are considered. In the actual approach, we explore in more details the growth rates at beta values much smaller than the typical value ( $\beta$ =0.5) used in the previous study. The obtained solutions of the dispersion equations represented in a parallel- and perpendicular- wavenumber diagram reveal new spectral features of the decay instability in low beta plasmas which have never 20 been pointed out in former analytical investigations. Furthermore, this result strengthens the idea of multi-channel coupling of decay instability as a proof-of-concept proposed by Comisel et al. 2019 in a former study based on numerical 3-D hybrid simulations. By following the Hall MHD formalism of Viñas and Goldstein 1991, we constructed the  $k_{\parallel} - k_{\perp}$  spectrum, and the results are in good agreement between the numerical simulation and the analytic treatment.

Changes in manuscript:

Page 2, new paragraph added, Lines 42 to Line 51:

"In the two-dimensional analysis of Viñas and Goldstein 1991, the authors concluded the decrease of the oblique-decay growth rates into narrow band-profiles (in the wavenumber domain) when moderate oblique propagation angles of the daughter waves are considered. In the actual approach, we explore in more details the growth rates at beta values much smaller than the typical value ( $\beta$ =0.5) used in the previous study. The obtained solutions of the dispersion equations represented in a parallel- and perpendicular- wavenumber diagram reveal new spectral features of the decay instability in low beta plasmas which have never been pointed out in former analytical investigations. Furthermore, this result strengthens the idea of multi-channel coupling of decay instability as a proof-of-concept proposed by Comisel et al. 2019 in a former study based on numerical 3-D hybrid simulations. By following the Hall MHD formalism of Viñas and Goldstein 1991, we constructed the  $k_{\parallel} - k_{\perp}$  spectrum, and the results are in good agreement between the numerical simulation and the analytic treatment." 35

#### Reviewer:

\* MAJOR. Discussion section is missing. Currently the manuscript is not clearly demonstrating that the obtained results are important for something. In the abstract there is a statement "Growth-rate maps serve as a useful tool for predictions of 40

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the wavevector spectrum of density or magnetic field fluctuations in various scenarios for the wave-wave coupling processes developing at different stages in space plasma turbulence." It would be very useful if authors could discuss at least some scenario showing that their quantitative/qualitative results are applicable and important. Also it would be good to discuss the limitations of the method.

Reply: We introduce the section "Discussion".

Former studies based on 2-D hybrid simulations, e.g., Matteini et al. 2010, reported a transversal spectrum of daughter waves <sup>10</sup> observed at the decay of a large amplitude Alfvén wave with linear polarization in low beta plasmas. Matteini et al. 2010 discussed the observed result in terms of the finite oblique angle of propagation of the imposed Alfvén pump wave with respect to the mean magnetic field. Gao et al. 2013 confirmed the former 2-D numerical study and even noticed a similar perpendicular spectrum of waves excited at the decay of a field-aligned pump wave. Two decades before, the two-dimensional analytical studies of Viñas and Goldstein 1991 suggest, on basis of the beta dependence of the oblique-decay growth rates, that the decay <sup>15</sup> process of a field-aligned circullary-polarized Alfvén wave might become important at finite angles of propagation for the daughter waves in low beta plasmas. The mechanism controlling the oblique decay instabilities has been discussed in relation with the coupling between the electrostatic (dominant at parallel propagation) and electromagnetic (dominant at high oblique angles) terms in the Hall MHD non-linear equations. By representing the growth rates provided by Viñas and Goldstein 1991 formalism in the wave-vector domain, the perpendicular spectrum of daughter waves can be straightfully noticed in Fig. 2, <sup>20</sup> while this feature is not obvious in the usual representation in Fig. 1. The perpendicular decay of daughter waves have been

while this feature is not obvious in the usual representation in Fig. 1. The perpendicular decay of daughter waves have been therefore predicted by the Hall MHD theory well before the first observations provided by 2D hybrid simulations in low beta plasmas. On basis of this result, one may expect that growth-rate maps in the wave vector domain could reveal new properties of parametric instabilities which can be further investigated or diagnosed by analytical treatments or by numerical simulations.

- <sup>25</sup> There are still open questions started in early studies about the role played by parametric instabilities to the generation of turbulent cascades (see e.g., Hoshino and Goldstein 1989, Ghosh et al., 1994). The perpendicular spectrum of daughter waves triggered by the nonlinear wave-wave coupling processes is superposed to the perpendicular spectrum of plasma turbulence generated by the post-saturation processes of the decay instability. The plasma anisotropy index as a quantitative measure of the directional (perpendicular) turbulence evolution describes the wavevector anisotropy of a two-dimensional spectrum
- <sup>30</sup> of magnetic fluctuations (see e.g., Shebalin 1983). An increasing number of such 2-D or even 3-D wavevector spectra are observationally being reported from in-situ solar wind measurements, see e.g., Narita et al. (PRL, 2010), Narita (NPG, 2014), Narita et al. (Frontiers 2014). It would be interesting to compare the beta dependence of the wavevector anisotropy of solar wind turbulence (see e.g., Comisel et al. 2014) with the anisotropy determined from the growth-rate maps of decay instability in its dependence on plasma beta parameter as they are displayed in Fig. 2.

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To our knowledge, the current study is the first step to bring together the results of the Hall MHD analytic approach and kinetic simulations, in emphasizing the generation of oblique daughter waves through their nonlinear evolution described by the MHD model. One may also notice that the MHD description can be helpful to clarify observed properties of the parametric instabilities which could originate in the kinetic approach by non-physical fluctuations due to the particle discreteness expected <sup>40</sup> in hybrid or particle-in-cell simulations.

As limitations of the Viñas and Goldstein 1991 method used in this analysis, the linear dispersion was restricted to excitation of only fundamental side-band daughter waves (i.e., no harmonics are allowed). The electron-inertia effect is neglected; one may expect that time scales or length of interest are larger than the electron gyroperiod or electron inertial length, respectively. The

<sup>45</sup> validity of the MHD approach with included Hall term is thus limited up to higher frequencies close to the ion gyro-frequency. The Landau and cyclotron damping effects have also been neglected in the MHD approach.

#### Changes in manuscript:

## Page 5, New section, Lines 125 to Page 6, Line 163:

'Former studies based on 2-D hybrid simulations, e.g., Matteini et al. 2010, reported a transversal spectrum of daughter waves observed at the decay of a large amplitude Alfvén wave with linear polarization in low beta plasmas. Matteini et al. 2010 discussed the observed result in terms of the finite oblique angle of propagation of the imposed Alfvén pump wave with respect to the mean magnetic field. Gao et al. 2013 confirmed the former 2-D numerical study and even noticed a similar perpendicular spectrum of waves excited at the decay of a field-aligned pump wave. Two decades before, the two-dimensional analytical studies of Viñas and Goldstein 1991 suggest, on basis of the beta dependence of the oblique-decay growth rates, that the decay process of a field-aligned circullary-polarized Alfvén wave might become important at finite angles of propagation for the daughter waves in low beta plasmas. The mechanism controlling the oblique decay instabilities has been discussed in relation 10 with the coupling between the electrostatic (dominant at parallel propagation) and electromagnetic (dominant at high oblique angles) terms in the Hall MHD non-linear equations. By representing the growth rates provided by Viñas and Goldstein 1991 formalism in the wave-vector domain, the perpendicular spectrum of daughter waves can be straightfully noticed in Fig. 2, while this feature is not obvious in the usual representation in Fig. 1. The perpendicular decay of daughter waves have been therefore predicted by the Hall MHD theory well before the first observations provided by 2D hybrid simulations in low beta 15 plasmas. On basis of this result, one may expect that growth-rate maps in the wave vector domain could reveal new properties of parametric instabilities which can be further investigated or diagnosed by analytical treatments or by numerical simulations.

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To our knowledge, the current study is the first step to bring together the results of the Hall MHD analytic approach and kinetic simulations, in emphasizing the generation of oblique daughter waves through their nonlinear evolution described by the MHD model. One may also notice that the MHD description can be helpful to clarify observed properties of the parametric <sup>30</sup> instabilities which could originate in the kinetic approach by non-physical fluctuations due to the particle discreteness expected in hybrid or particle-in-cell simulations.

As limitations of the Viñas and Goldstein 1991 method used in this analysis, the linear dispersion was restricted to excitation of only fundamental side-band daughter waves (i.e., no harmonics are allowed). The electron-inertia effect is neglected; one may expect that time scales or length of interest are larger than the electron gyroperiod or electron inertial length, <sup>35</sup> respectively. The validity of the MHD approach with included Hall term is thus limited up to higher frequencies close to the ion gyro-frequency. The Landau and cyclotron damping effects have also been neglected in the MHD approach.'

#### Reviewer:

\* lines 48-51. The description of k+- and w+- should clearly specify that those describe daughter waves. ko is vector but described in text as field-aligned wavevector, this is confusing.

### Reply:

The following comment is added in the text.

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Changes in manuscript:

Page 3, Line 60 to Line 62:

', where  $k^{\pm}$  and  $\omega^{\pm}$  describe the wavevectors and frequencies of the daughter waves, respectively,  $k_0 = k_0 e_{\parallel}$  with  $k_0$  - wavenumber of the Alfvén pump wave,  $e_{\parallel}$  is the unity vector parallel to the mean magnetic field, and  $\omega_0$  - frequency of the  ${}^{5}$  Alfvén pump wave.'

### Reviewer:

\* line 48. What is "state vector", it is never used in the manuscript.

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## Reply:

The state vectors indeed are not provided or discussed in the manuscript.

Changes in manuscript:

15 Page 3, Line 58:

'Each linear mode is specified by its frequency and wavevector.'

### Reviewer:

<sup>20</sup> \* line 53. It should be more clearly explained what authors mean by "limitation for the frequency domain".

Reply:

The sentence is rewritten at the end of the new section "Discussion".

Changes in manuscript:

Page 6, Line 159 to Line 162:

<sup>25</sup> 'The electron-inertia effect is neglected; one may expect that time scales or length of interest are larger than the electron gyroperiod or electron inertial length. The validity of the MHD approach with included Hall term is thus limited up to higher frequencies close to the ion gyro-frequency.'

#### Reviewer:

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\* line 87. The meaning of "domain (1,3) and (0,3)" not explained.

Reply:

Solving the dispersion equation to provide solutions in the  $(k_{\parallel}, k_{\perp})$  domain is a demanding numerical task and therefore solutions are searched in a limited wavenumber range where decay instability is expected. (for instance, the domain below <sup>25</sup>  $k_{\parallel}/k_0 = 1$  where modulational instability could be operational is omitted).

Changes in manuscript:

Page 4, Line 97 - Line 99:

'Because this procedure requires a demanding numerical task, the solutions of dispersion equations are searched in a limited wavenumber range where decay instability is expected (for instance, the domain below  $k_{\parallel}/k_0 = 1$  where modulational instability could be operational is omitted).'

### Reviewer:

\* line 104-105. Needs more clarification how one distinguishes Alfvén and sound daughter modes, as the figure shows only the magnetic field magnitude oscillations. For parallel propagation waves are not compressional so it is surprising to see 10 magnetic field compressions.

### Reply:

Indeed, the explanation of Fig. 3 right panel obtained from the hybrid simulation is not enough elaborated.

Changes in manuscript:

Page 4, Line 115 to Line 118:

'Magnetic field fluctuations are represented at negative wavenumbers corresponding to the backward propagating Alfvén daughter waves while the compressional forward propagating waves (acoustic-like) are shown as density fluctuations along the positive wavenumbers.'

### Reviewer:

\* Figure 2: caption does not specify the difference between top left and top right panels.

Reply: Figure 2 caption is modified as below.

Changes in manuscript:

Figure 2 caption:

' Map of the growth rates in the  $k_{\parallel} - k_{\perp}$  wavenumber domain for right- (panel left) and left- (panel right) handed polarization of the Alfvén pump wave corresponding to the analysis shown in Fig. 1 at plasma  $\beta$ =0.02.'

### Reviewer:

\* Figure 3: It would be good to mark the location of pump waves booth in analytical and simulation results.

right panel) are magnetic field and density fluctuations normalised to something, why are magnetic field and density fluctuations the same? It is not the case in other simulations.

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Reply:

Done. The pump wave is marked now in both panels.

Indeed, the magnetic field and density fluctuations look similar although not the same. The answer is that the 2-D wavenumber <sup>5</sup> spectrum is filtered in frequency domain such that only the frequencies expected for the Alfvén daughter waves at  $\omega = \omega_0 - \omega_s$ and density mode at  $\omega = \omega_s$  (here  $\omega_s$  is the frequency of the ion acoustic wave) are shown.

Changes in manuscript:

Page 4, Line 118 to Line 120:

<sup>10</sup> 'Furthermore, the 2D wavenumber spectra are filtered in frequency domain such that only the frequencies expected for the Alfvén daughter waves at  $\omega = \omega_0 - \omega_s$  and for density mode at  $\omega = \omega_s$  (here  $\omega_s$  is the frequency of the ion acoustic wave) are shown.'

Reviewer:

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\* Outlook section should be named "Summary" or "Summary and outlook".

Reply: Done.