



Ducting of incoherent scatter radar waves by field-aligned irregularities

Michael T. Rietveld^{1,2}, and Andrew Senior³

¹ EISCAT Scientific Association, Ramfjordbotn, Norway.

5 ² University of Tromsø-The Arctic University of Norway, Norway.

³ Independent Researcher, Lancaster, UK.

Correspondence to: Michael Rietveld (mike.rietveld@eiscat.uit.no)

Abstract. We provide an explanation for a mysterious phenomenon that has been recognized in recent years in EISCAT
10 UHF incoherent scatter radar (ISR) measurement during many high power HF ionospheric pumping experiments. The
phenomenon is an apparent increase in electron density observed above the HF reflection altitude, extending up to the
observable limits usually in the range 400-650 km, as shown in several publications in recent years. It was shown by Senior
et al. (2013) that several examples of these enhanced backscatter could not be explained by increases in electron density. A
summary of characteristics of the backscatter enhancements is presented as well as the results of a survey of events. We
15 propose that medium- to large-scale HF-induced field-aligned irregularities (tens to hundreds of m scale) act to refract the
radar signals along the magnetic field, thereby acting as a guide so that the free-space r^2 spreading of the signals no longer
applies. The nature of the irregularities and the physical mechanism of their production by powerful HF waves is an exciting
topic for future research since, surprisingly, they appear to be preferentially excited by X-mode waves. The explanation
proposed here involving HF-induced irregularities may well apply to other ISR observations of the ionosphere in the
20 presence of specific natural irregularities.

1 Introduction

The HF facility (Rietveld et al., 2016) near Tromsø, Norway, consists of 12 transmitters of nominally 100 kW covering the
frequency band from 4 to 8 MHz which can be connected to one of three antenna arrays. Array 1 covers 5.5-8.0 MHz with a
mid-band gain of 30 dBi, Array 2 covers 4-5.5 MHz with a mid-band gain of 24 dBi, and Array 3 also covers 5.5-8.0 MHz
25 with a mid-band gain of 24 dBi. In many HF pumping experiments with the EISCAT facilities near Tromsø (69.6N 19.2E)
during the recent solar cycle maximum, when ionospheric F region critical frequencies were commonly in the range 5.5 to
more than 8 MHz, apparent increases in electron density were observed by the 933 MHz incoherent scatter radar above the
HF reflection height when the radar was pointed along the geomagnetic field. Some of these are shown in published papers
by Blagoveshchenskaya et al. (2011a, 2011b, 2013, 2015, 2017, 2018), Cheng et al. (2014), Borisova et al. (2016, 2017), Wu
30 et al. (2017) and Senior et al. (2013) and they have been presented by us in several conferences and workshops. They have
been observed rather often, but no systematic study has been made so far. Some characteristics as observed by these authors,
and which can be seen in several of the published examples listed above, are given here. The enhancements extend from
about the HF reflection height to as high an altitude as there is still enough backscatter to be detected, typically 400-500 km.
Importantly, they are observed only when the UHF radar is within 0.5° (Bazilchuk, 2019) of the geomagnetic field. They are
35 excited by both O and X-mode HF pumping but more commonly with X-mode (see survey below, and Blagoveshchenskaya
et al., 2018), with an enhancement factor reaching 2 for X-mode heating (Blagoveshchenskaya et al., 2018; Bazilchuk,
2019). They have usually been observed with HF pump frequencies of 5.4 MHz and above, although there seems to be one
published case near 4.9 MHz (Blagoveshchenskaya et al., 2011a) and we have found two cases at 4.544 MHz (see
supplement 1). They appear within some tens of seconds after HF turn on, but this is poorly documented and studied. The



40 decay times are similarly poorly studied but seems to be a few tens of seconds to a few minutes. When stepping in frequency
around a harmonic of an electron gyro frequency using an O-mode HF pump wave the enhancements are strongest when the
pump is near or above the gyroharmonic (Blagoveschenskaya et al., 2018). They appear to be anti-correlated with HF-
induced electron temperature increases as shown in Fig. 1 of this paper, Fig. 2 of Borisova et al. (2016), Fig. 3a of [Borisova et
al.](#) (2017) and Fig. 4 of Blagoveshchenskaya et al (2018).

45

The apparent density enhancements were interpreted as such because they result from a good theoretical spectral fit to the
measured ion-line spectra, which were not affected by enhancements caused by plasma instabilities, which occur in a narrow
altitude at and below the HF reflection height. To first order, the electron density is proportional to the backscattered power
of the ion line. Senior et al. (2013) however, showed a typical example which could not have been a real density
50 enhancement because the natural plasma frequency at the relevant heights showed no corresponding increase. Furthermore,
they showed another case from 2001, near the previous solar cycle maximum, the apparent density enhancement seen by the
Tromsø UHF radar near 300 km was not seen by the two remote receiving antennas in Kiruna and Sodankylä when pointed
to that common volume at 300 km. Because the enhancements are in the backscattered ion-line, they have been called Wide-
Altitude Ion Line Enhancements, or WAILEs for short. Subsequent examination of the plasma line backscattered power in
55 the first case studied by Senior et al. (2013) showed that it was also enhanced by approximately the same amount as the ion
line power, showing that the effect is caused by some increase in backscattered radar power at UHF frequencies, irrespective
of whether the scattering is from ion acoustic or electron acoustic plasma waves. Figure 3 of Borisova et al. (2017) confirms
this increase in plasma line power at 300 km for another event associated with the X-mode induced ion line power
enhancements and the apparent electron density increases, well above the reflection height of the HF wave at around 230
60 km.

We first present an example of the WAILE phenomenon pointing out some of the features. Next the results of a survey to
investigate some of the features occurrence of this phenomenon will be presented. Then we present an explanation of the
phenomenon in terms of refraction and guiding of VHF and UHF radar waves by field-aligned elongated electron density
65 irregularities. To support this model some raytracing results are presented. Finally, possible implications for other ISR radar
measurements of the natural ionosphere are presented.

1.1 Example of a WAILE

We present an example in Figure 1 which shows the apparent density enhancements during an experiment on 16 November
2011 where the O-mode HF pump wave was modulated in a 30-minute cycle of 20 minutes on and 10 minutes off starting at
70 11:15:05 UT and again at 12:50:05 UT. The HF on times are indicated by red lines above the top panel. The HF beam and
the 933 MHz ISR beam were pointed along the magnetic field line, 12° south of zenith, and the HF effective radiated power
was calculated to be approximately 660 MW O-mode and 11 MW X-mode. During the on period the frequency was first
constant at 6.7 MHz for 2 minutes after which it was stepped up every 10 s to 7 MHz in 108 steps of 2.778 kHz each. The
5th harmonic of the electron gyrofrequency is calculated to be 6.862 MHz at 200 km using the IGRF magnetic field model
75 for 2011. This example, from a standard analysis using 30s integration time, shows several of the typical features mentioned
above, but it is also a complicated example because of the frequency stepping and slightly atypical in that it is for O-mode
heating. No attempt was made to remove backscatter from satellites or ion lines enhanced by plasma instabilities near the
reflection height of the HF wave. The WAILEs are seen as the periodic apparent electron density increases between 300 and
600 km, in synchronism with the HF. The rise time of the WAILEs are not easy to determine from this experiment because
80 the initial enhancement after HF turn-on is very weak and barely noticeable in this plot until about 7 minutes after each start,



when the frequency has reached about 6.78 MHz. After that the enhancements are very clear until the end of the HF pulse, when they decay in about 2 minutes. A slight dip in the WAILE is apparent for about a minute in the first, second and fifth HF pulse at about 8 minutes before the end of each pulse when the frequency is near the gyroharmonic at 6.86 MHz. This interesting frequency dependence will not be investigated further in this paper, but is an important aspect for further studies.

85

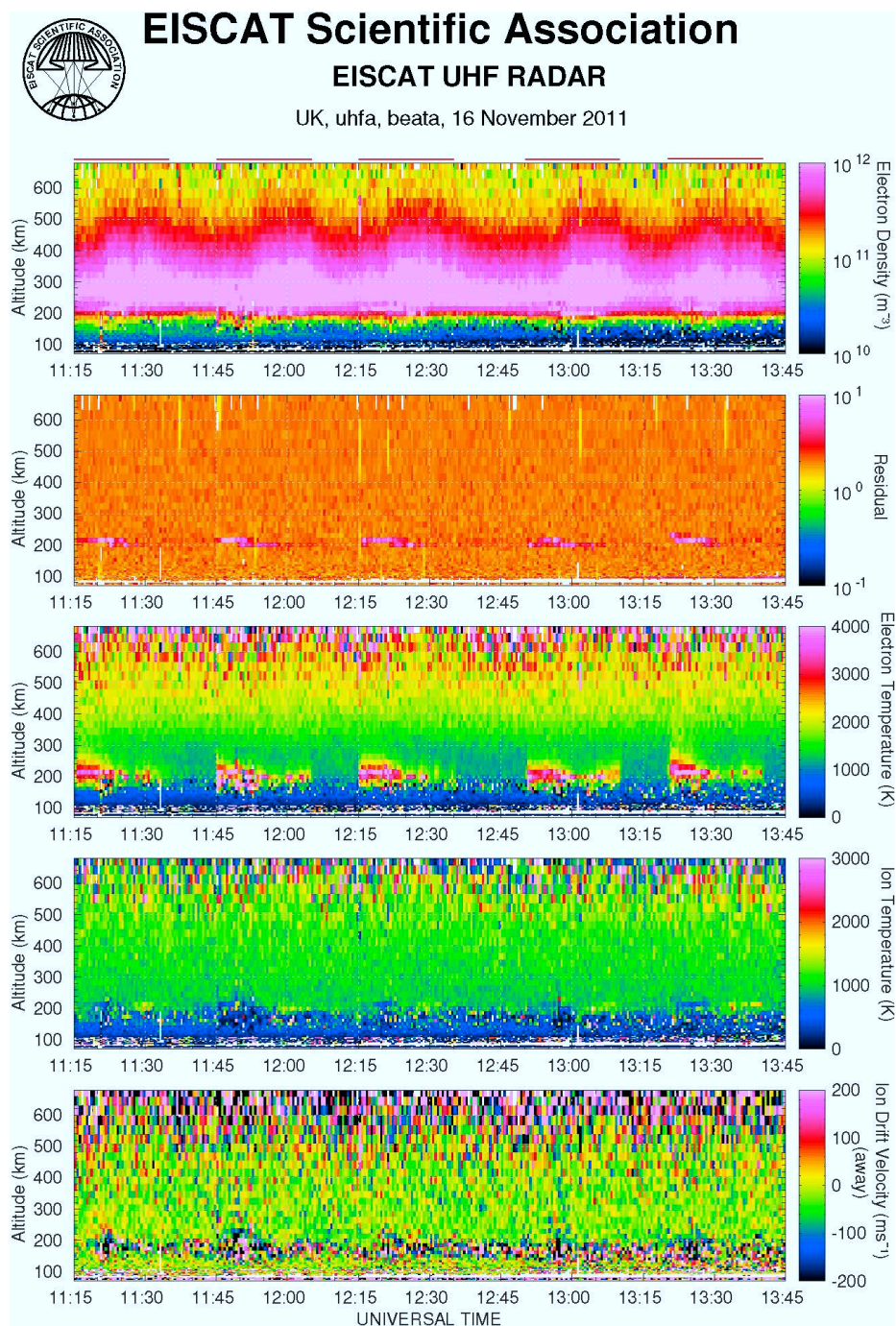


Figure 1: An example of the wide altitude ion line enhancement (WAILE). The top panel shows the enhancements interpreted by the analysis as an electron density increase between about 200 km and 600 km. The HF on times are shown by red lines above the top panel. The second panel (residual) shows how well the measured spectra agree with theoretical spectra for a Maxwellian plasma.



The second panel labelled 'residual', shows a parameter which should be close to 1 for a good fit of the measured spectrum to a theoretical spectrum for a Maxwellian plasma. It shows that the theoretical spectra were a very good fit everywhere so that the derived electron density and temperature and ion temperature and velocity are reliable except for a narrow region between mainly 190 and 220 km, around the HF pump reflection height near 200 km. There is some localised electron heating above 220 km extending up to 300 km and sometimes beyond which is most pronounced in the first 8 minutes of most of the HF pulses, when the WAILEs are weakest. This illustrates the observation made above that the WAILEs anticorrelate with HF-induced electron temperature increases. This electron heating is a well-known phenomenon associated with upper-hybrid plasma instabilities and the formation of decameter-scale striations when pumping away from a gyroharmonic (Honary et al. 1995). The observation that WAILEs are observed more frequently with X-mode pumping ties in with the anti-correlation with electron temperature enhancements since X-mode pumping results in rather weak Ohmic heating. A very similar observation to the example in Fig.1 is shown in Fig.6 of Borisova et al. (2016) for the 6th gyroharmonic.

2 Survey of Events

To provide some statistical evidence of the conditions under which the WAILE phenomenon is observed, a search was performed of all experiments carried out in the period 2001-01-01 to 2018-07-06 (the date when the search was begun). As anecdotally there is little evidence of WAILEs during experiments using antenna Array 2 (restricted to frequencies below 5.5 MHz) and since Array 3 (frequencies above 5.5 MHz, but lower gain than Array 1) was inoperative for part of the period, the search was confined to experiments using Array 1.

The search was performed in two parts. Firstly, an automated process was used to scan the computer-generated experiment log files to identify cases where Array 1 was used and where the pump frequency was below the F-region critical frequency. Then the resulting list of cases was used to manually inspect quick-look plots of the UHF radar data to determine whether or not WAILEs were observed for each case.

The automated process defined a "case" as a UT hour during which (a) at least two transmitters were active on Array 1; and (b) the lowest pump frequency used was below the highest foF2 measured by the EISCAT Dynasonde (Rietveld et al., 2008) during the hour plus a margin of 0.81 MHz. The margin is approximately the value by which fxF2 exceeds foF2 and means that experiments where the pump frequency was below the F-region critical frequency would be included regardless of the polarisation (O- or X-mode) used. If no Dynasonde data were available for the hour in question, condition (b) was not applied to avoid unnecessarily excluding useful cases at this stage.

The manual process inspected quick-look plots of the UHF radar data available in the EISCAT data archive together with the manual logs of HF Facility operations. For each case, it was first determined if the experiment mode was suitable for observation of WAILEs: the UHF radar must have been pointed field-aligned for at least part of the hour and running in a mode providing F-region coverage; the HF Facility must have used pump-on periods greater than about 30 s. Then for suitable cases, the UHF radar electron density range-time plots were inspected to decide whether WAILEs were observed with either O- or X-mode pumping during that hour.

The result of the search was a list of cases giving the UT date and hour of each case, the minimum pump frequency and maximum foF2 during that hour and a set of codes indicating whether the experiment was unsuitable, quick-look plots were



unavailable or if WAILEs could or could not be convincingly identified for O- and X-mode pumping if each polarisation was used. The results of the search are available in the supplement.

The main shortcoming of this search is the manual identification of the presence of WAILEs which is necessarily subjective. This, together with the inevitable risk of human error in inspecting the data and recording the results adds a degree of uncertainty to the results. Fully automating the process was deemed unfeasible. To do so would have required automatically analysing the electron density data to compare the density between pump-on and pump-off periods. The computer-generated log files give the transmitter status only at discrete time points and this would make reconstructing the pump cycle ambiguous.

The search resulted in a total of 1449 cases, of which 606 which were useful (experiment mode suitable, quick-look data available) and were used for further analysis.

Table 1 shows the observed occurrence of WAILEs with O- and X-mode pumping. A pooled two-proportion z-test comparing the proportion of cases with WAILEs between O- and X-mode rejects the null hypothesis that the proportion is the same with a z score of 7.7 indicating strong evidence that the proportion is higher with X-mode.

	O-mode	X-mode
WAILEs observed	214	183
Total cases	521	261
Proportion	0.41	0.70

Table 1: Counts and proportions of WAILE occurrence for O- and X-mode

3 Suggested Explanation

We now suggest a mechanism whereby the backscatter enhancements are explained in terms of UHF radio wave propagation along large scale irregularities, but we make no attempt to explain the more interesting problem of the creation of the postulated irregularities.

In standard incoherent scatter radar theory, it is usually assumed that at 933 MHz the radar transmissions propagate in straight lines as in free space because that frequency is very much larger than the maximum ionospheric plasma frequency which may reach 10-20 MHz at the most in the F region. A radio wave at 933 MHz, however, does experience some weak refraction from refractive index changes caused by electron density irregularities. An electromagnetic wave propagating at a small angle to a plane where the refractive index changes, will experience refraction. The ionosphere and magnetosphere is full of electron density irregularities at various scales which are aligned along the magnetic field. At the critical angle, given by $\sin^{-1}(n_2/n_1)$, where n_1 and n_2 are the refractive indices at an interface, a ray will propagate parallel to the interface and larger angles will be totally reflected if $n_2 < n_1$.

For simplicity let us assume that the HF pumping creates magnetic field aligned irregularities with a 5% electron density depletion spaced 100 m apart starting near the HF reflection height of typically 200 km and extending several tens of km along the magnetic field. With a background electron density of $8 \times 10^{11} \text{ [m}^{-3}\text{]}$ or plasma frequency of 8 MHz, this gives a critical angle for a 933 MHz wave of 89.889° , which is a grazing angle of 0.111° . This means that all rays within 0.111° of



field-aligned will be ducted by the irregularity and will not fall off with the usual r^2 dependence as rays outside this angle will. The one-way half-power beam width of the UHF radar is 0.6° , (Folkestad et al., 1983), which is also the value for the 'opening angle' used by the Guisdap incoherent scatter analysis software (Lehtinen and Huuskonen, 1996). For field-aligned pointing the ducted solid angle is therefore 14% of the total solid angle of the transmitted beam (Solid angle = $2\pi(1-\cos(\theta))$ where θ is half the apex angle of the subtended cone).

For a background plasma frequency of 5.4 MHz with a 5% density depletion the grazing angle for critical incidence is 0.074° , which corresponds to only 6% of the solid angle of the transmitted beam, which could largely explain why the ducting effect is smaller and has not been observed very often for heating frequencies less than 5.4 MHz.

The proposed mechanism also explains why backscatter enhancements were never seen at the remote receiving stations when they were still operating at 933 MHz. In the WAILE example from 11 November 2001 presented by Senior et al. (2013), the remote receivers in Kiruna and Sodankylä showed no enhancements while the backscatter signal received at Tromsø did. This is understandable from our ducting model for two reasons. Firstly, the volume of the scattering region above Tromsø as seen from the remote sites is much larger because of the greater distance to the receiver, so the horizontal redistribution of radar power from Tromsø is mostly contained within the receiver's field of view. Secondly, the scattered signals to the two stations are at such a large angle to the magnetic field (18° and 28° for Kiruna and Sodankylä respectively at 278 km above Tromsø) that they are effectively not refracted by the irregularities.

180

4 Raytracing model

To give a more detailed model of the ducting hypothesis, we performed some ray-tracing of 933 MHz radio waves in a medium with sinusoidal refractive index irregularities perpendicular to the direction of propagation and a limited spatial extent in the main direction of propagation. The ray-tracing equations were solved using the differential equation solver Isode.m in GNU Octave (version 4.4.1) or ode45.m in MATLAB.

The irregularities are modelled by a refractive index, $n = 1 + f(r)$ where n is the refractive index at r , the position vector (x,y) of a point on the ray and f represents "irregularities" with a sinusoidal variation in the y direction. We assume that the magnetic field is vertical, i.e. along x . In the x direction, the irregularity amplitude is tapered by $1+\tanh(x)$ functions to give gradual transitions into and out of the irregularity region. In the following example we used a maximum value for the refractive index perturbation, A , of 1×10^{-6} which, for a 933 MHz radio wave corresponds to 2.7% change in electron density at $8 \times 10^{11} \text{ [m}^{-3}\text{]}$ (8 MHz plasma frequency). The extent of the irregularities in the x direction, L , was 80 km starting at 220 and ending at 300 km and had a wavelength, W , of 50 m in the y direction. The lower height of 220 km is typically around or just above the reflection height of typical O-mode HF pump waves used in heating experiments, which is where we expect the HF-irregularities to be formed. Figure 2 shows the paths of a set of 30 rays uniformly spaced at angles between 0° and 0.3° launched from a point source at the transmitter $(0,0)$ out to a distance of 600 km along the x -axis. Clearly a focusing is seen for rays close to the x axis with regions of lower ray density further out.

190

195

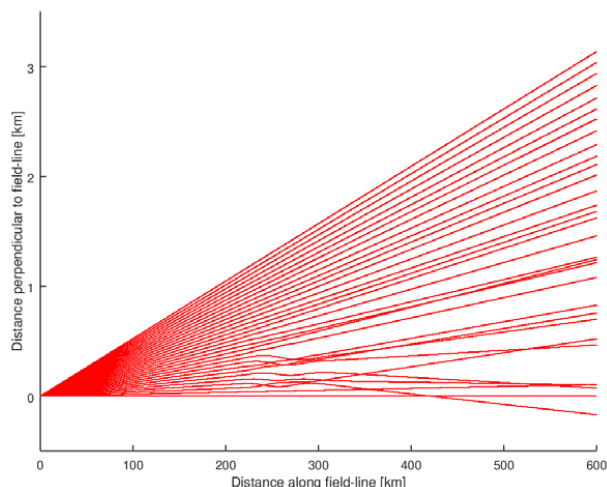


Figure 2. Ray paths of a set of 30 up-going 933 MHz waves from the transmitter at (0,0) with launch angles evenly spaced between 0 and 0.3° passing through a region of irregularities with $A=1 \times 10^{-6}$, $W=50$ and $L=80$ km between 220 and 300 km.

We are simulating a 3-dimensional (3D) situation, of which the model is just a slice. In 3D, the density of the rays being proportional to power means that the solid angle subtended by a "bundle" of rays would be inversely proportional to radiated power. So in our 2D model the planar angle between rays is inversely proportional to the square-root of the power, or simply the wave electric field, E . To approximate the real situation better, we performed such raytracing with many more rays, but with their spacing varying to approximate the narrow transmitted beam of the radar. We used the formula for the radiation field from a circular aperture given by:

$$E = 2 \lambda J_1[(\pi D/\lambda) \sin\theta] / (\pi D \sin\theta) \quad (1)$$

where D is the diameter of the antenna's aperture in meters, λ is the radar wavelength in meters, θ is the angle of propagation from the boresight direction in radians, and J_1 is the first-order Bessel function. The rays were spaced approximately inversely proportional to E .

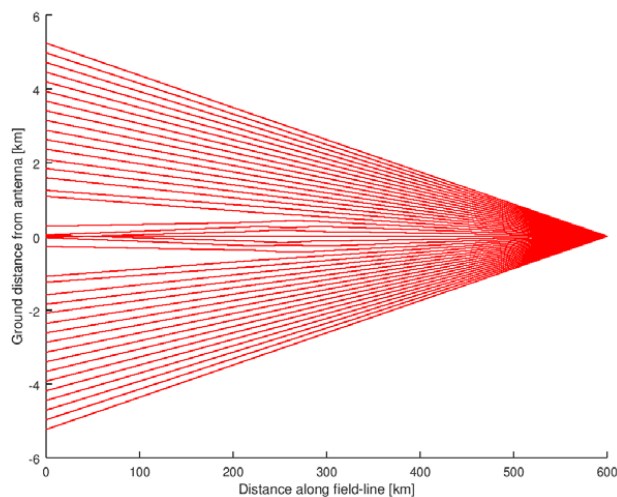


Figure 3. 50 backscattered down-going 933 MHz rays from one up-going ray arriving at 600 km. The 41 rays pass through a region of irregularities with $A=1 \times 10^{-6}$, $W=50$ and $L=80$ km between 220 and 300 km.



We next consider the received wave from the scattering region, which is simpler than transmitted beam because the incoherent scatter is isotropic. As an example of the scatter from a point along the field-aligned radar beam at 600 km, Fig. 3 shows a set of 50 rays backscattered downwards within a cone of 0.64° along the field line for a refractive index perturbation of 1×10^{-6} with a wavelength, W , of 50 m. The radar receiving antenna is at (0,0) and the perturbations are 220 to 300 km above the radar antenna. One can clearly see a higher density of rays at the antenna caused by the irregularities compared to the uniform distribution of rays one would have without irregularities. There are, off course, also regions where the density of rays is lower, in order to conserve the total energy. If most of the backscattering electrons are within a region of diameter ~ 6 km (the radar beam width at 600 km) then some of the ray paths will end up at the antenna with weaker intensity. So the total backscatter seen by the radar is more complicated than these illustrative simplified examples of ray tracing from a single point source. The resultant backscattered signal is the integrated effect of the ducting of the transmitted signal to a large volume and then the many backscattered and partly ducted signals from all the electrons within this relatively large volume.

To make a more realistic model, we used the position of each upward ray's arrival point at height, h , to re-transmit a set of randomly distributed rays with launch angles between $\pm 0.32^\circ$ of the reverse direction with which the up-going ray arrived downwards through the same irregularities as for the upward rays. This limited launch cone angle is made because we are mainly interested in only those rays that end up within the antenna aperture and within the main beam angle and to keep the number of rays at a manageable level. The number of backscattered rays, n , per ray arriving at h , is limited to typically 30 to make the total ray-tracing calculation time acceptable, but large enough to give representative and fairly robust results. Tests were made to ensure that the results were not critically sensitive to this number.

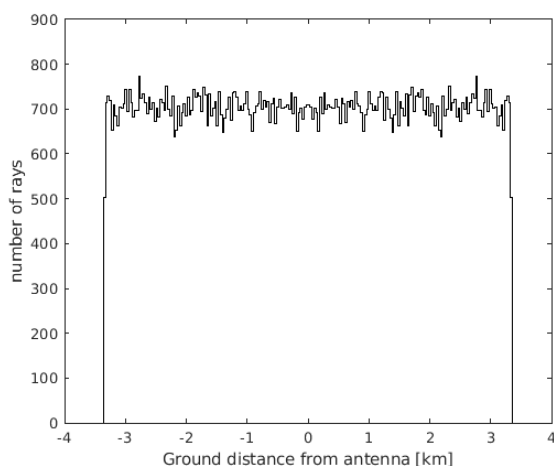


Figure 4. Histogram of rays arriving on the ground from backscatter at 600 km, in a limited cone of launch angles, with no irregularities. The bin width is 32 m, equal to the antenna diameter.

The intensity of the radar beam is proportional to the number of rays per unit distance in this 2-dimensional view, so we can show the density of rays in the form of histograms. Figure 4 shows a histogram of the number of rays received at the ground for the case of free space propagation, i.e. with no irregularities present. As expected, the distribution is uniform and the edges of the distribution are determined by the limited re-launch cone. The bin width in this histogram and that in Fig. 5 is 0.032 km, equal to the antenna diameter. Figure 5 shows a histogram for the case with irregularities with $A = 1 \times 10^{-6}$, $W = 50$ m, and $L = 80$ km. The number of backscattered rays at the 32-m diameter antenna are enhanced in Fig. 5 by about 20% compared to the smooth ionosphere case in Fig. 4, but there are also regions of slightly weaker signal within 1-2 km of



the antenna. But this is still too simplistic an estimate of the actual enhancement since some of the rays may reach the
245 antenna aperture at incident angles outside the main beam.

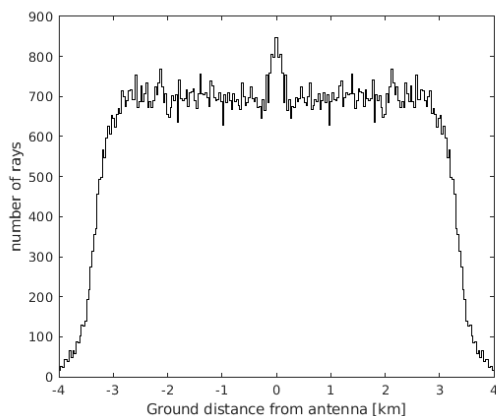


Figure 5. Histogram of rays arriving on the ground showing enhanced backscatter from 600 km with irregularities having parameters $A=1 \times 10^{-6}$, $W=50$ m, and $L=80$ km. The limited cone of launch angles and bin width are the same as in Fig. 4.

250 For the quantitative results presented below we count the number of rays arriving within a 0.032 km long segment at the
origin, corresponding to the UHF antenna aperture diameter, and with an incident angle of less than $\pm 0.35^\circ$ to approximate
the received signal within the 0.6° half-power beam-width of the main lobe of the antenna. The ratio of this number to the
number of rays received for the case of no irregularities is the measure of the enhanced backscatter strength (enhancement
factor) from that height. This ratio was squared to give the enhancement factors plotted in Figures 6-9 to convert the one
255 dimensional model results to the two dimensional reality.

The results of the raytracing are presented in Figure 6 for backscatter from heights of 300 km and from 600 km in Figure 7,
as plots of enhancement factor for four different values of A , the irregularity strength, and 21 different values of the
irregularity wavelength, W . Each point in Figs 6-9 is the result of 77370 rays being launched from the backscattered height,
260 h which took approximately 55 minutes to calculate using MATLAB on a Linux workstation. The smallest irregularity
wavelength modelled here is 20 m, 62 times the radar wavelength of 0.32 m, which means that the raytracing approach is
still valid.

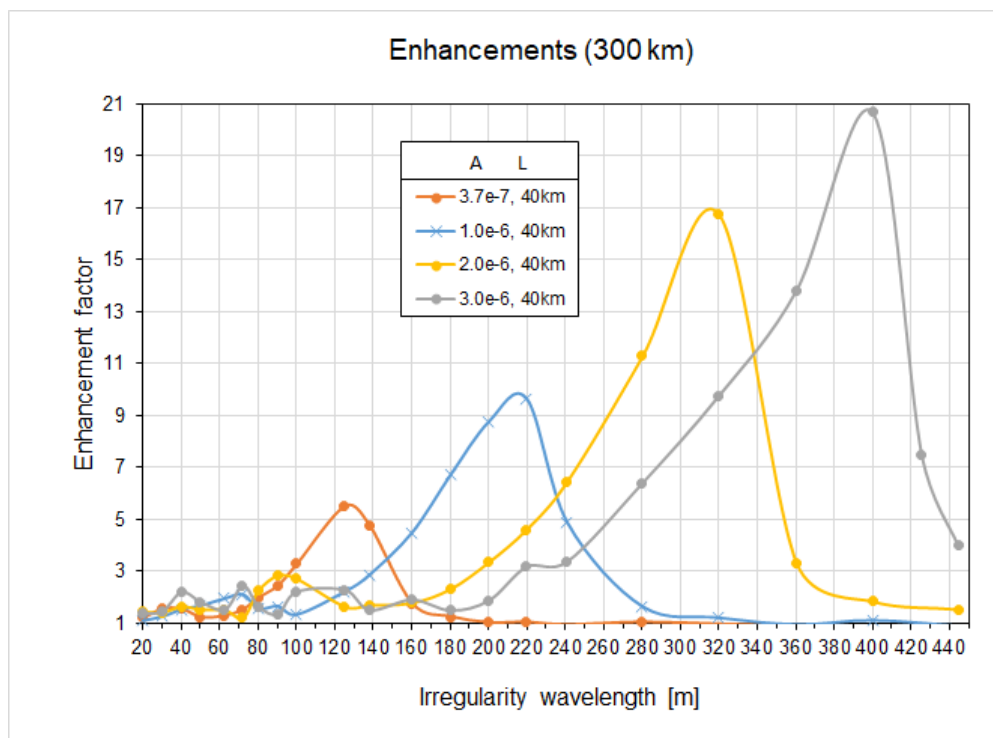


Figure 6. Enhancement factors for 933 MHz backscatter from 300 km calculated from the density of rays entering the antenna beam compared to the case of free space propagation, for various depths of irregularity, A, and wavelength, W.

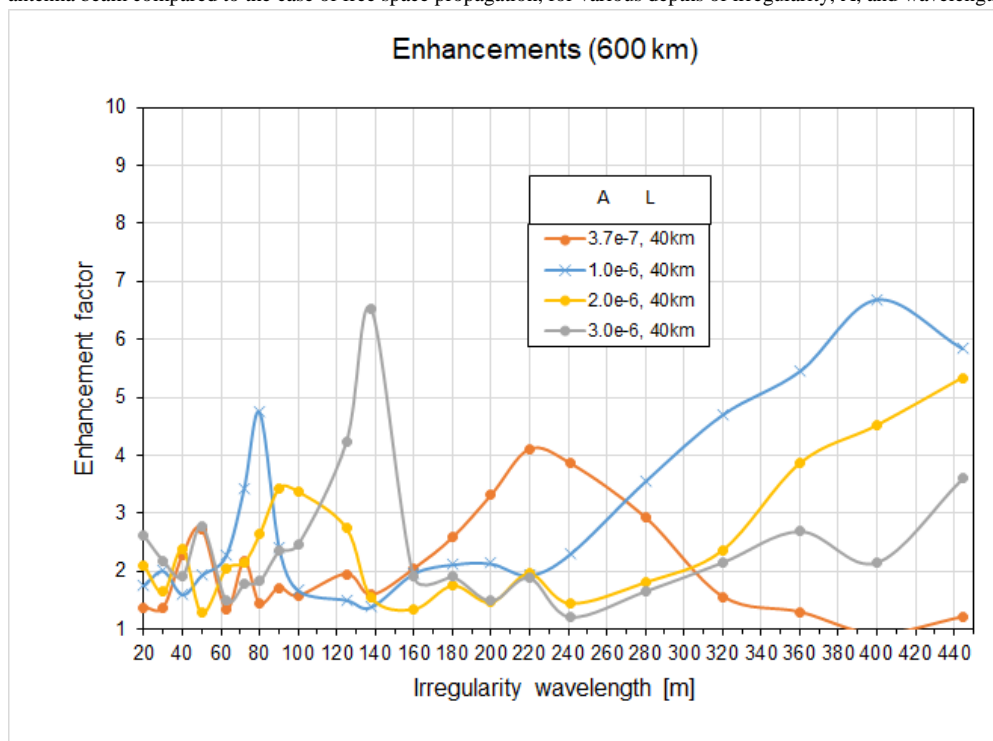


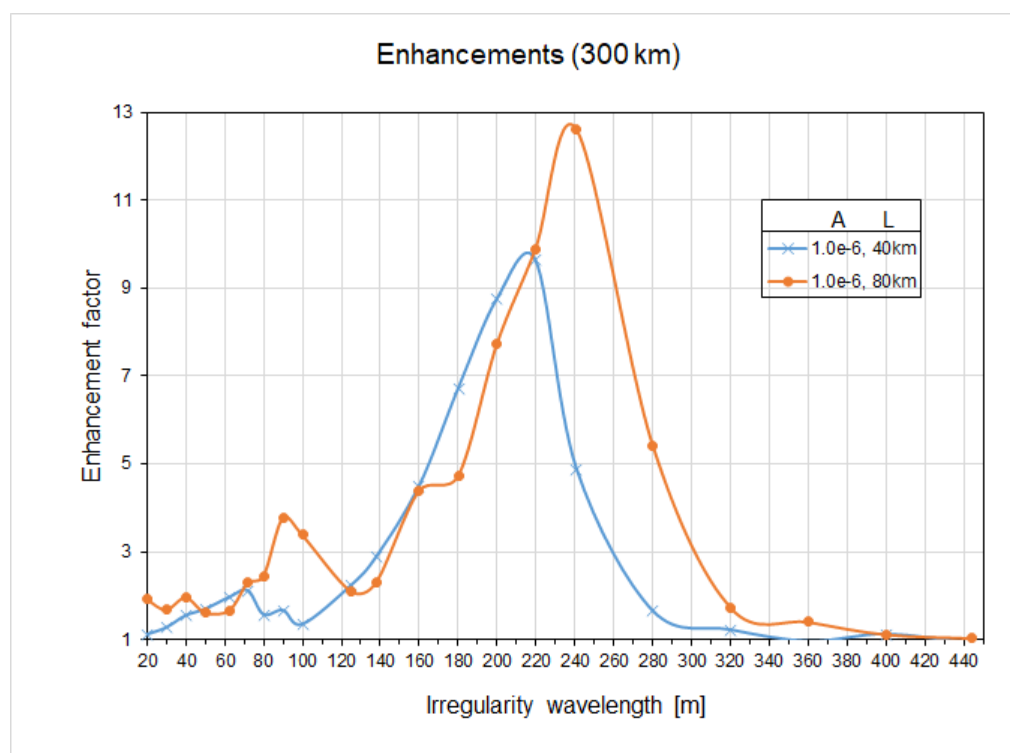
Figure 7. Enhancement factors for 933 MHz backscatter from 600 km calculated from the density of rays entering the antenna beam compared to the case of free space propagation, for various depths of irregularity, A, and wavelength, W.



270 The results for the enhancements at 300 km show a systematic variation with irregularity wavelength with stronger enhancements peaking at larger scale irregularities and the peak being stronger for stronger irregularities. We have never observed such large enhancements as in these peaks with the strongest results being perhaps 2. For most of the irregularity strengths with a wavelength less than about 200 m we find enhancement factors greater than 1 and less than 2. For a given irregularity scale size one cannot simply conclude that stronger irregularities give stronger enhancements.

275

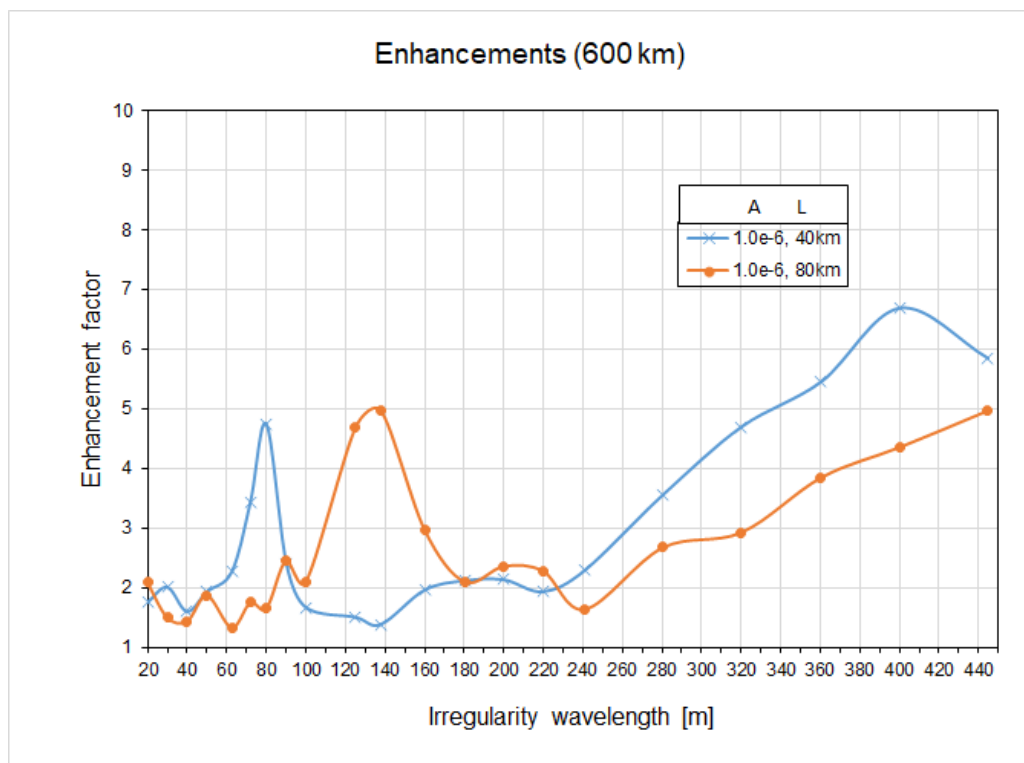
The results for the enhancements at 600 km show a more complex variation which is not so easy to understand. It would appear that the main peak enhancements seen at 300 km get weaker and move to much longer (about two times longer) wavelengths but other peaks appear or get stronger below about 150 m. As for the 300 km case there are enhancements greater than 1 and less than 2 for all irregularity strengths over a wide range of wavelengths less than about 300 m.



280

Figure 8. Enhancement factors for 933 MHz backscatter from 600 km for two lengths of the irregularity, L, along the field line.

285 Figures 8 and 9 show the effect on the enhancements at 300 km and 600 km respectively, of having longer irregularities along the magnetic field for the case of $A = 1 \times 10^{-6}$. In both cases the curve for 40 km long irregularities is moved to larger irregularity wavelengths. For many irregularity scales the enhancements are larger for the longer scale irregularities and the peak enhancements in Fig. 6 and perhaps in Fig. 7, are larger with increasing irregularity strength, A. Both these effects might be expected since stronger and longer irregularities should help guiding the rays.



290

Figure 9. Enhancement factors for 933 MHz backscatter from 600 km for two lengths of the irregularity, L, along the field line.

295 For the 12 October 2012 event studied by Senior et al. (2013), the enhancement factor was fairly constant with altitude with the value ~ 1.5 . If we were to take an amplitude of 10^{-6} and 40 km length, then Figs. 6 to 9 would suggest a transverse scale either below ~ 60 m or around ~ 100 m in order to get a similar enhancement at 300 and 600 km. We have no data to check this prediction, but the numbers do not seem unreasonable.

300 We note that the smallest scale of irregularities we have modeled here ($W = 20$ m) is close to the decameter scale of irregularities that produce backscatter on HF radars like SuperDARN. There does not seem to be an obvious connection between CUTLASS radar backscatter and WAILEs, although this should be studied more closely. HF pumping produces decameter scale irregularities that are often correlated with electron temperature increases so this would suggest that they might anti-correlate with WAILEs.

305 The raytracing modelling results suggest that WAILEs should be quite sensitive to the spectrum of irregularity produced by HF pumping and might thereby become a new diagnostic for these irregularities. We should warn, however, of taking the results of the raytracing here too literally. The model is rather idealised in that it incorporates only one irregularity scale at a time and in reality we certainly have a large spectrum of irregularities. The modelling here is only intended as a proof of the principle that ducting of the ISR waves can explain the backscatter enhancements. More detailed, independent modelling should be performed to verify our ideas.
310



5 Evidence of the irregularities

Experimental evidence for large scale (100's m to km scale) irregularities is lacking for most of the heating experiments we have examined at EISCAT. The EISCAT UHF radar cannot directly detect irregularities with scale sizes smaller than the beam width which is ca. 2.4 km at F region heights. There is of course evidence for large scale irregularities from earlier
315 dedicated campaigns, for example using radio-star scintillations (Frey et al., 1984) but these lacked simultaneous suitable field-aligned UHF radar data. Honary et al. (1993) deduced the existence of electron density and temperature irregularities along the field with horizontal scale sizes of 6-10 km. There is plenty of evidence for decameter-scale irregularities seen by HF coherent radars like CUTLASS (e.g. Senior et al., 2004) but there does not seem to be a correlation between the presence of this backscatter and WAILEs. The presence of such decameter scale irregularities does not seem to be sufficient to
320 produce WAILEs.

The width of the natural plasma line depends to a large extent on the range of electron densities found within the scattering volume, so this could give an indication of the presence of irregularities, from large to small scales, with the largest scale being limited by the width of the radar beam. There is a suggestion that the natural plasma line measured during the WAILEs
325 studied by Senior et al. (2013) was broadened (unpublished work by one of the authors, AS), but this is something to be examined in future work.

Frolov et al. (2016) show evidence, from the Demeter satellite, of field-aligned large scale irregularities (km scale) extending from about 230 to well over 400 km in height, which act as ducts for very low frequency (VLF) waves. The ducts had
330 dimensions of 80-100 km at heights around 660 km, and were produced exclusively by O-mode heating. In many cases smaller structuring of ~10 km was seen within the main duct. Such large scale plasma structuring may well be related to the phenomena we observe. At EISCAT there have been fewer heating experiments performed with DEMETER or other satellites, and then they were sometimes made without ISR observations, so we cannot claim to have seen similar irregularities at Tromsø.

335 6 Predictions for other radars

The refractive index irregularity strength of 10^{-6} used in the above examples for a 931 MHz radar wave corresponded to an electron density irregularity of 2.68% for $N_e = 8 \times 10^{11} \text{ m}^{-3}$ or an 8 MHz plasma frequency. For the EISCAT 224 MHz VHF radar that is co-located with the UHF radar, this would correspond to a 0.155% density irregularity so that for a given
340 irregularity strength the focusing effect would be stronger at VHF than at the UHF frequency, neglecting other differences such as antenna beam widths. So for the commonly used ISR frequency of around 430-440 MHz such as used in the American radars and for the EISCAT Svalbard Radar at 500 MHz the effect would also be larger than at 931 MHz. Not all the radars can point along the magnetic field however. The Arecibo 430 MHz and the Jicamarca 50 MHz cannot point field-aligned, and the EISCAT VHF antenna, while physically able to point along the field, is restricted from doing so for operational reasons. So this effect should be stronger for lower frequency radars, especially in the 220-250 MHz range like
345 the EISCAT VHF radar and the new EISCAT-3D radar (McCrea et al., 2015) being built. Of the radars mentioned here only Arecibo has an HF facility that could produce these irregularities. Unfortunately, the geographical separation of the new EISCAT-3D radar transmitter being built at Skibotn from the present heating facility and the very stringent field-aligned nature of the phenomenon is likely to prevent the WAILE phenomenon being observed directly by the new radar in monostatic mode or with the presently planned bi-static receiver sites. The HF facility cannot tilt its beam in the east-west
350 plane towards the field-line at Skibotn. The multi-beam and higher time and spatial resolution expected with this new radar



should, however, allow the postulated irregularities to be observed from the side. There are two other lower frequency ISR's at around 158 MHz: one in Kharkiv, Ukraine and one in Irkutsk, Russia. Only the Irkutsk radar can point near field-aligned.

7 Implications of the ducting hypothesis

355 The presence of ducting irregularities described here should also affect signals from radio stars or satellite beacon signals received on the ground. The problem with testing this at the high latitude of Tromsø is that there are few if any signal sources in the field-aligned direction. An increase in scintillation of beacon signals in the field-aligned direction has been seen in very early studies, e.g. Figure 1 of Singleton and Lynch (1962). These authors also discuss the same mechanism proposed here, that of reflection at angles larger than the grazing angle to explain some scintillation effects. Rush and Colin
360 (1958) show some raytracing examples of the effect of long (with respect to a wavelength) cylindrical columns of electrons on various VHF/UHF waves. These papers are concerned mainly with the scintillation phenomena and do not explicitly predict intensity enhancements. Although we have concentrated on the backscatter intensity enhancements, enhanced scintillation of ISR signals may also have occurred. Such scintillation of the received signal is probably masked by the typical integration times of 30 or 60s used in the incoherent scatter analysis, but would be another interesting thing to
365 investigate.

Do natural irregularities affect ISR returns in the same way as the HF-induced irregularities? It is possible that naturally occurring irregularities, known to occur often in the auroral zone, may affect ISR measurements in the same way as the artificial irregularities discussed here. It remains to be seen whether there is something special about the HF-induced
370 irregularities that might make this phenomenon unique to HF-pumping experiments. A naturally-occurring counterpart might well go unnoticed since electron density will often show variable enhancements due to varying soft or energetic particle precipitation. This same particle precipitation may also produce irregularities which duct the ISR waves contributing to distortions of the measured density profiles. We emphasize the importance of using plasma line measurements when available with ISR's for determining electron density rather than relying on the ion line due to the possible invalidation of
375 the assumptions concerning the interpretation of the backscatter power.

Total electron content (TEC) measurements derived from field-aligned GNSS differential phase measurements should be examined to see whether irregularities could produce errors in the derived densities, since such waves will also be affected by field-aligned irregularities, to a greater or lesser extent depending on frequency.

380

The fact that X-mode pumping preferentially produces WAILEs over O-mode pumping is an important addition to the variety of other phenomena that have been found in recent years to be excited by X-mode waves and which were expected to be impossible (Blagoveshchenskaya et al. 2017, 2018). The interesting question is how do X-mode waves produce the irregularities that we postulate produce the WAILEs. One possibility is through self-focusing of the HF wave, but why this
385 should be more efficient for X-mode than O-mode is a mystery since the O-mode generally produces stronger electron temperature enhancements through upper-hybrid resonance instabilities and thereby stronger irregularities for the self-focusing. These and other questions relating to the effects of X-mode pumping are questions at the forefront of HF active experiments in the future.



8 Conclusions

390 We have provided a qualitative explanation for the mysterious phenomenon of apparent electron density enhancements seen
in magnetic field aligned UHF radar data during many HF-pumping experiments. The mechanism is refraction leading to
ducting of the incoherent scatter radar waves by large scale density irregularities. This mechanism explains why the
enhancements were not observed in bistatic measurements. A simple raytracing model explains ISR backscatter
enhancements greater than 1 for a wide range of irregularity scale sizes. The more interesting problem lies in the unknown
395 nature of the irregularities, and their excitation by both O- and, preferentially, X-mode pump waves. The postulated
irregularities causing the WAILE phenomenon and other effects of X-mode pumping are poorly understood and are at the
forefront of HF active experiment research, both experimentally and theoretically. Another intriguing question is whether
natural irregularities can produce the same enhancements. Because of this uncertainty, plasma line measurements should be
used whenever possible for determining the electron density in ISR measurements when measuring close to or along the
400 magnetic field of the ionosphere. The width of the natural plasma line is another parameter which should be exploited to
determine the scale of ionospheric irregularities.

Data availability. Plots of the analysed UHF radar data used in this study are available from <https://portal.eiscat.se/madrigal/>
and the heating facility log files are available on-line with access details available from the first author.

Supplement. Results of a survey of WAILE events between 2001 and 2018 are listed.

405 *Author contributions.* MTR developed the explanation, performed the raytracing and prepared the paper. AS developed the
raytracing model and made the survey.

Competing interests. The authors declare that they have no conflict of interest.

Acknowledgments. EISCAT is an international association supported by research organisations in China (CRIRP), Finland
(SA), Japan (NIPR and STEL), Norway (NFR), Sweden (VR), and the United Kingdom (NERC). We thank Craig
410 Heinselman, Ingemar Häggström, Nataly Blagoveshchenskaya, Wu Jun, Björn Gustavsson and Juha Vierinen for discussions
on the data and interpretation of this mysterious new phenomenon.

References

- Bazilchuk, Z.: Angular dependence of wide altitude ion line enhancements (WAILEs) during ionospheric heating at the
EISCAT Tromsø Facility, Faculty of Science and Technology Department of Physics and Technology, Masters thesis,
415 <https://munin.uit.no/handle/10037/15663>, 2019.
- Blagoveshchenskaya, N. F., Borisova, T. D., Yeoman, T. K., Rietveld, M. T., Ivanova, I. M., and Baddeley, L. J.: Artificial
small-scale field-aligned irregularities in the high latitude F region of the ionosphere induced by an X-mode HF heater wave,
Geophys. Res. Lett., 38, L08802, doi:10.1029/2011GL046724, 2011a.
- Blagoveshchenskaya, N. F., Borisova, T. D., Rietveld, M. T., Yeoman, T. K., Wright, D. M., Rother, M., Lühr, H., Mishin,
420 E. V., and Roth, C.: Results of Russian experiments dealing with the impact of powerful HF radio waves on the high-latitude
ionosphere using the EISCAT facilities, Geomagnetism and Aeronomy, ISSN 0016-7932, 51(8), 1109-1120, 2011b.



- Blagoveshchenskaya, N. F., Borisova, T. D., Yeoman, T. K., Rietveld, M. T., Häggström, I., Ivanova, I. M.: Plasma modifications induced by an X-mode HF heater wave in the high latitude F region of the ionosphere, *Journal of Atmospheric and Solar-Terrestrial Physics*, 105-106, 231-244, 2013.
- 425 Blagoveshchenskaya, N. F., Borisova, T. D., Yeoman, T. K., Häggström, I., Kalishin, A. S.: Modification of the high latitude ionosphere F region by X-mode powerful HF radiowaves: Experimental results from multi-instrument diagnostics, *J. Atmos. Sol.-Terr. Phys.*, 135, 50-63, doi:10.1016/j.jastp.2015.10.009, 2015.
- Blagoveshchenskaya, N. F., Borisova, T. D., Kalishin, A. S., Yeoman, T. K., Häggström, I.: First observations of electron gyro-harmonic effects under X-mode HF pumping the high latitude ionospheric F-region, *Journal of Atmospheric and Solar-*
- 430 *Terrestrial Physics*, 155, 36-49, 2017.
- Blagoveshchenskaya, N. F., Borisova, T. D., Kalishin, A. S., Kayatkin, N. V., Yeoman, T. K., Häggström, I.: Comparison of the effects induced by the ordinary (O-Mode) and extraordinary (X-Mode) polarized powerful HF radio waves in the high-latitude ionospheric F region, *Cosmic Research*, 56, 1, 11–25, 2018.
- Borisova, T. D., Blagoveshchenskaya, N. F., Kalishin, A. S., Rietveld, M. T., Yeoman, T. K., and Häggström, I.:
- 435 Modification of the high-latitude ionospheric F region by high-power HF radio waves at frequencies near the fifth and sixth electron gyroharmonics, *Radiophysics and Quantum Electronics*, 58(8), (Russian Original 58(8), August, 2015), DOI 10.1007/s11141-016-9629-2, 2016.
- Borisova, T. D., Blagoveshchenskaya, N. F., Yeoman, T. K., and Häggström, I.: Excitation of artificial ionospheric turbulence in the high-latitude ionospheric F region as a function of the EISCAT/Heating effective radiated power,
- 440 *Radiophysics and Quantum Electronics*, 60(4), DOI 10.1007/s11141-017-9798-7, 2017.
- Cheng, M.-S., Xu, B., Wu, Z.-S., Li, H.-Y., Xu, Z.-W., Wu, J. Wu, J.: A large increase in electron density in ionospheric heating experiment, *Chinese J. Geophys. (in Chinese)*, 57(11), 3633-3641, doi:10.6038/cjg20141117, 2014.
- Folkestad, K., Hagfors, T., Westerlund, S.: EISCAT: An updated description of technical characteristics and operational capabilities, *Radio Sci.*, 18, 867-879, 1983.
- 445 Frey, A., Stubbe, P., Kopka, H.: First experimental evidence of HF produced electron density irregularities in the polar ionosphere diagnosed by UHF radio star scintillations, *Geophys. Res. Lett.*, 11(5), 523-526, 1984.
- Frolov V. L., Rapoport, V. O., Schorokhova, E. A., Belov, A. S., Parrot, M., and Rauch, J.-L.: Features of the electromagnetic and plasma disturbances induced at the altitudes of the Earth's outer ionosphere by modification of the ionospheric F2 region using high-power waves radiated by the SURA heating facility, *Radiophys. Quant. Electron.*, 59(3),
- 450 177-198, doi:10.1007/s11141-016-9688-4, 2016.
- Honary, H., Stocker, A. J., Robinson, T.R., Jones, T.B., Wade, N. M., Stubbe, P., and Kopka, H.: EISCAT observations of electron temperature oscillations due to the action of high power HF radio waves, *J. Atmos. Terr. Phys.*, 55, 10, 1433-1448, 1993.
- Honary, H., Stocker, A. J., Robinson, T.R., Jones, T.B., and Stubbe, P.: Ionospheric plasma response to HF radio waves
- 455 operating at frequencies close to electron gyroharmonics, *J. Geophys. Res.*, 100, 21489-21501, 1995.
- Lehtinen, M. and Huuskonen, A.: General incoherent scatter analysis and GUISDAP, *Journal of Atmospheric and Terrestrial Physics*, 58(1-4), 435–452. [https://doi.org/10.1016/0021-9169\(95\)00047-X](https://doi.org/10.1016/0021-9169(95)00047-X), 1996.
- McCrea, I., Aikio, A., Alfonsi, L., Belova, E., Buchert, S., Clilverd, M., Engler, N., Gustavsson, B., Heinselmann, C., Kero, J., Kosch, M., Lamy, H., Leyser, T., Ogawa, Y., Oksavik, K., Pellinen-Wannberg, A., Pitout, F., Rapp, M., I. Stanislawski,
- 460 I., and Vierinen, J.: The science case for the EISCAT_3D radar, *Progress in Earth and Planetary Science*, 2:21, DOI 10.1186/s40645-015-0051-8, 2015.
- Rietveld, M. T., Wright, J. W., Zobotin, N., and Pitteway, M. L. V.: The Tromsø Dynasonde, *Polar Science*, 2, 1, 55-71, doi:10.1016/j.polar.2008.02.001 2008.



- Rietveld, M. T., Senior, A., Markkanen, J., and Westman, A.: New capabilities of the upgraded EISCAT high-power HF
465 facility, *Radio Sci.*, 51, 1533-1546, doi:10.1002/2016RS006093, 2016.
- Rush, S. and Colin, L.: The effects on radio astronomical observations due to longitudinal propagation in the presence of
field-aligned ionization, *Proc. I. R. E.*, 46, 356- 357, 1958.
- Senior, A., Borisov, N. D., Kosch, M. J., Yeoman, T. K., Honary, F., and Rietveld, M. T.: Multi-frequency HF radar
measurements of artificial F-region field-aligned irregularities, *Ann. Geophys.*, 22, 3503-3511, 2004.
- 470 Senior, A., Rietveld, M. T., Haggstrom, I., and Kosch, M. J.: Radio-induced incoherent scatter ion line enhancements with
wide altitude extents in the high-latitude ionosphere, *Geophys. Res. Lett.*, 40(9), 1669-1674, DOI: 10.1002/grl.50272, 2013.
- Singleton, D. G. and Lynch, G. J. E.: The scintillation of the radio transmissions from Explorer VII - II Some properties of
the scintillation producing irregularities, *J. Atmos. Terr. Phys.*, 24, 363-374, 1962.
- Wu, J., Wu, J., Rietveld, M. T., Haggstrom, I., Zhao, H., and Xu, Z.: The behavior of electron density and temperature
475 during ionospheric heating near the fifth electron gyrofrequency, *J. Geophys. Res. Space Physics*, 122,
doi:10.1002/2016JA023121, 2017.