

1 **MMS observations of energetic oxygen ions at the duskside magnetopause**
2 **during intense substorms**

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4 **Chen Zeng^{1,2}, Suping Duan¹, Chi Wang^{1,2}, Lei Dai¹, Stephen Fuselier^{3,4}, James Burch³,**
5 **Roy Torbert⁵, Barbara Giles⁶, Christopher Russell⁷**

6
7 ¹State Key Laboratory of Space Weather, National Space Science Center, Chinese Academy of Sciences,
8 China.

9 ²University of Chinese Academy of Sciences, China.

10 ³Southwest Research Institute, San Antonio, Texas, USA.

11 ⁴The University of Texas at San Antonio, San Antonio, Texas, USA.

12 ⁵University of New Hampshire, Space Science Center, Durham, New Hampshire, USA.

13 ⁶NASA, Goddard Space Flight Center, Greenbelt, MD, USA.

14 ⁷University of California Los Angeles, IGPP/EPSS, Los Angeles, California, USA.

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16 *Corresponding author:* Chi Wang (cw@spaceweather.ac.cn), Suping Duan(spduan@nssc.ac.cn)

17 **Abstract**

18 Energetic oxygen ions (1-40 keV) observed by the Magnetospheric Multiscale (MMS) satellites at
19 the duskside magnetopause boundary layer during phase 1 are investigated. There are 57 duskside
20 magnetopause crossing events during intense substorms ($AE > 500$ nT) are identified. These 57 events of
21 energetic O^+ at the duskside magnetopause include 26 events during the expansion phase and 31 events
22 during the recovery phase of intense substorms. It is found that the O^+ density in the duskside
23 magnetopause boundary layer during the recovery phase (0.081 cm^{-3}) is larger than that during the
24 expansion phase (0.069 cm^{-3}). The 26 events of energetic O^+ ion at the duskside magnetopause during
25 intense substorm expansion phase are all under the southward interplanetary magnetic field (IMF). There
26 are only 7 events under northward IMF and they all occurred during the intense substorm recovery phase.
27 The density of energetic O^+ at the duskside magnetopause ranges from 0.007 to 0.599 cm^{-3} . The
28 maximum density of O^+ occurred during the intense substorm recovery phase and under southward IMF.
29 When the IMF is southward, the O^+ density shows an exponential increase with the IMF B_z absolute
30 value. Meanwhile, The O^+/H^+ density ratio shows an exponential growth with the IMF B_y . These results

31 agree with previous studies in the near-Earth magnetosphere during intense substorm. It is suggested that
32 O^+ abundance in the duskside magnetopause boundary layer has a close relation with O^+ variations in the
33 near-Earth magnetosphere during intense substorms.

34 **1 Introduction**

35 Single charged oxygen ions (O^+) in the magnetosphere are exclusively from the ionosphere. They
36 are an important element in the mass and energy transport in the magnetospheric dynamic process,
37 especially during the expansion phase and recovery phase of intense substorms (e.g., Daglis et al., 1991,
38 1996; Duan et al., 2017; Fok et al., 2006; Ohtani et al., 2011; Ono et al., 2009; Nosé et al., 2000; Yau et
39 al., 1997, 2012; Kronberg et al., 2014). Processes in the magnetotail due to substorm can result in auroral
40 electrojet activity. This activity is generally caused by field-aligned currents increase and reflected by
41 the AE index (Tang and Wang, 2018). Previous studies have found that the density and energy density
42 of O^+ significantly increase with AE index in the near-Earth magnetosphere during the intense substorm
43 (e.g., Lennartsson and Shelley, 1986; Daglis et al., 1991, 1994; Duan et al., 2017). Lennartsson and
44 Shelley (1986) proposed that oxygen ions with energies less than 17 keV/e could provide 50% of the
45 density in the plasma sheet during disturbed geomagnetic activity. They found the increase in the O^+
46 energy density was strongest around local midnight where O^+ became the most abundant ion at AE~1000
47 nT. In the near-Earth plasma sheet (NEPS), The O^+ energy density has an explosively increases with AE
48 index in the range of larger than 500 nT during the intense substorm expansion phase (Daglis et al., 1994).
49 Previous studies reported that the O^+ from the nightside auroral region could rapidly feed in the near-
50 Earth magnetosphere during substorm expansion phase (e.g., Daglis and Axford, 1996; Duan et al., 2017;
51 Yu et al., 2013). Otherwise, the solar wind dynamic pressure also influences the oxygen content of ion
52 outflow from the ionosphere. Using the Thermal Ion Dynamics Experiment (TIDE) on the Polar satellite.
53 Elliott et al., (2001) found both the O^+ density and parallel flux increased with the solar wind dynamic
54 pressure.

55 The O^+ outflowing from the ionosphere with low energy of eV are accelerated to about 500 eV at
56 the high altitude polar region (e. g., Yau and André, 1997). Then they are convected tailward into the
57 lobe and the plasma sheet boundary layer. After O^+ enter the NEPS of magnetotail, they can be energized
58 up to tens of keV during intense substorm dipolarizations (e.g., Birn et al., 2013; Duan et al., 2017;

59 Fok, et al., 2006; Nosé et al., 2000; Ono et al., 2009; Yau et al., 2012). The inductive electric field
60 associated with substorm dipolarization is very significant for accelerating particles in the NEPS (e.g.,
61 Dai et al., 2014, 2015; Duan et al., 2011, 2016; Lui et al., 1999). Duan et al. (2017) reported that the O⁺
62 from the lobe or the plasma sheet boundary layer were efficiently accelerated by the kinetic Alfvén
63 eigenmode with significant unipolar electric field and rapidly feed in the NEPS during intense substorm
64 dipolarizations. These energetic O⁺ in the NEPS can be injected into the inner magnetosphere and drift
65 westward into the duskside outer magnetosphere (e.g., Ganushkina et al., 2005).

66 Oxygen ions decay from the ring current can leak into the dayside magnetopause boundary layer
67 (e.g., Li et al., 1993; Ebihara et al., 2011). Li et al. (1993) reported that the ring current O⁺ with tens of
68 keV energy interacted with the Pc 5 waves and then lost towards the dayside magnetopause. The solar
69 wind dynamic pressure enhancement plays a key role in the ring current particle loss into the outer
70 magnetosphere. This pressure enhancement pushing the magnetopause to move inward leads to a
71 reduction of the scale length of the magnetic field magnitude gradient along the magnetopause. The
72 magnetic gradient drift speed across the magnetopause will increase. So the ring current oxygen ions
73 along the magnetic gradient drift path can easily enter into the outer magnetosphere (Kim et al., 2005).
74 Ebihara et al. (2011) proposed that the field line curvature scattering was more effective on the loss of
75 energetic oxygen ions with its large gyro-radius. The energetic oxygen ions with pitch angle of ~90
76 degrees are more prone to leak into the dayside magnetopause.

77 The distribution of energetic oxygen ions density at the dayside magnetopause is asymmetric and it
78 has a close relationship with the interplanetary magnetic field (IMF) (e.g., Bouhram et al., 2005; Phan et
79 al., 2004; Luo et al., 2017). Bouhram et al. (2005) pointed out that the O⁺ density in the duskside (on
80 average 0.053 cm⁻³) magnetopause is higher than that in the dawnside (on average 0.014 cm⁻³). They
81 found O⁺ was the dominant contributor to the mass density (30%) on the duskside magnetopause in
82 comparison to 3% in the dawnside and 4% near the noon. The dawn-dusk asymmetries of the energetic
83 O⁺ (>~274 keV) distribution in three different regions (dayside magnetopause, near-Earth nightside
84 plasma sheet, and tail plasma sheet) are also observed by Luo et al., (2017). They found that the energetic
85 O⁺ distributions were mainly influenced by the dawn-dusk IMF directions and the enhancement of ion
86 intensity strongly related to the location of the magnetopause reconnection.

87 There is ample evidence that magnetospheric ions could participate in the magnetopause
88 reconnection and directly escape along the reconnected open field lines (e.g., Sonnerup et al., 1981;

89 Fuselier et al., 1991, 2016; Slapak et al., 2012, 2015; Wang et al., 2014; Liu et al., 2015). The energetic
90 O^+ with energies larger than 3 keV in the reconnection jets at the duskside mid-latitude magnetopause
91 under steady southward IMF was reported by Phan et al. (2004). Zong et al., (2001) observed O^+ energy
92 dispersion due to time-of-flight (TOF) effects at the duskside magnetopause under southward IMF and
93 it was assumed that O^+ was escaping from the ring current along the reconnected field lines during steady
94 reconnection. However, Fuselier et al. (1989) reported that O^+ from the high latitude ionosphere were
95 not associated with any substorm cycle. O^+ from the high latitude ionosphere could form the O^+ rich
96 boundary layer in the low latitude magnetopause. When O^+ enter the reconnection jets, the reconnection
97 rate is likely reduced by the mass-loading but not suppressed at the magnetopause (Fuselier et al. 2019).

98 At present, variations of O^+ abundance (O^+/H^+) in the duskside magnetopause boundary layer
99 during intense substorm ($AE > 500$ nT) with AE index and solar wind conditions (e.g. IMF B_y , IMF B_z
100 and solar wind dynamic pressure) are still not clear. Previous studies of O^+ abundance variations during
101 substorms are mainly focused on the magnetotail or the near-Earth region (e.g., Duan et al., 2017; Nosé
102 et al., 2000; Ohtani et al., 2011). The Magnetospheric Multiscale (MMS) mission gives us an opportunity
103 to focus on the O^+ in the duskside magnetopause region. In this study, we investigate statistical features
104 of energetic O^+ at the duskside magnetopause and their relations with AE index and solar wind conditions
105 (e.g. IMF B_y , IMF B_z , and solar wind dynamic pressure) during the intense substorms ($AE > 500$ nT).

106 **2 Instrumentation and Data**

107 This study used data from the Magnetospheric Multiscale (MMS) mission. This mission comprises
108 four identical satellites that were launched on 2015 March 13 into an elliptical 28-inclination orbit with
109 perigee around $1.2 R_E$ and apogee around $12 R_E$ (Burch et al., 2016; Fuselier, et al., 2016b). The electric
110 field \mathbf{E} is from the electric double probe (EDP) (Ergun et al., 2016; Lindqvist et al., 2016), and magnetic
111 field \mathbf{B} is from the Fluxgate Magnetometer (FGM) (Russell et al., 2016). The plasma data are from the
112 Fast Plasma Investigation (FPI) and the Hot Plasma Composition Analyzer (HPCA). The FPI provides
113 plasma (electrons and ions) distribution functions at 32 energies from 10 eV to 30 keV. And it has a high
114 time resolution of 0.03 s for electrons and 0.15 s for ions in the burst mode and 4.5 s in the fast mode
115 (Pollock et al., 2016). The FPI does not discriminate between different ion species. While the HPCA
116 provides ion composition (H^+ , He^{++} , He^+ and O^+) measurements in the energy range from 1 eV/q to 40

117 keV/q (Young et al., 2016). Although the HPCA instrument employs radio frequency (RF) unit to
118 artificially reduce the proton fluxes in some areas where the proton fluxes are intense, there still exists a
119 low level of background that affects the O^+ fluxes in the magnetosheath. The majority of the O^+ fluxes in
120 the magnetosphere side of the magnetopause are at energies from 1 keV to 40 keV and that band below
121 1keV visible in the magnetosheath side are observations outside the RF operating range and
122 contamination from high proton fluxes. Due to this contamination, the O^+ at energies from 1 keV to 40
123 keV in the magnetopause boundary layer are considered in our study. The O^+ density recalculated from
124 the HPCA distribution functions at this energy range by using the Space Physics Environment Data
125 Analysis System (SPEDAS) software package. More details about SPEDAS can be found in
126 Angelopoulos et al. (2019). The solar wind parameters and AE index are available from the OMNI data
127 in CDAweb (<http://cdaweb.gsfc.nasa.gov/>). The data from the MMS 4 satellite are adopted in our
128 investigation since the data difference from other three spacecraft is negligible. This is due to spacecraft
129 separation and scales of particle motion.

130 **3 Results**

131 **3.1 Detailed event on 3 October 2015**

132 Figure 1 presents the three components of the IMF in Geocentric Solar Magnetospheric (GSM)
133 coordinates, solar wind dynamic pressure, as well as AU, AL, and AE index during the time of interest
134 from 14:30 to 16:30 UT on 3 October 2015. During this interval, the IMF B_x component is negative all
135 the time (Figure 1a). Its maximum value is about -1 nT at ~15:16 UT. The IMF B_y component is almost
136 negative except at ~14:32 and ~16:23 UT. The negative IMF B_z component is also observed during this
137 interval as shown in Figure 1c. The minimum value of the IMF B_z component is about -7.1 nT at ~14:30
138 UT. The solar wind dynamic pressure is only at the beginning of the time interval about 2 nPa. Then, it
139 increases sharply at 15:00 UT and reaches its maximum value about 4.4 nPa at ~15:12 UT. These solar
140 wind conditions led to an intense substorm ($AE > 500$ nT), as Figure 1g shown. The AE index is defined
141 as $AE = AU - AL$. Generally, the substorm onset time is characterized by the AL index starting to
142 significantly decrease and the AE index significantly increase. The interval of the AL index decreasing
143 from onset to its minimum is defined as the substorm expansion phase. The interval of the AL index
144 increasing from the minimum to the quiet time level is regarded as the substorm recovery phase. From

145 Figure 1e to 1g, the substorm onset time is about 14:45 UT marked by the AL index starting to sharply
146 decrease and AE index increase. After the AE index significantly increases and the AL index decreases
147 (Figure 1f), the AL and AE indexes reach their minimum and maximum values about -750 nT and 1000
148 nT at ~15:20 UT, respectively. This interval from ~14:45 to ~15:20 UT is regarded as the intense
149 substorm expansion phase. Then, the intense substorm enters the recovery phase as the AL index
150 gradually increases and AE index decreases after ~15:20 UT. The two blue dashed lines indicate the time
151 interval of the magnetopause boundary layer crossing. According to the above description, we can
152 identify this magnetopause boundary layer crossing occurred during the recovery phase of intense
153 substorm. The identification of the magnetopause boundary layer will be described later.

154 Figure 2 shows the overview of the magnetopause inbound crossing from 15:00 to 16:00 UT on 03
155 October 2015. During the magnetopause crossing, MMS 4 satellite was located at about (6.0, 8.8, -5.1)
156 R_E in GSM as shown in the bottom of Figure 2. From top to bottom, panels 2a and 2b show the magnetic
157 and electric fields in GSM from FGM and EDP, respectively. Ion and electron temperatures, plasma
158 density and ion velocity in GSM from FPI L2 data products are shown in Figures 2c-e. Figure 2f shows
159 the H^+ and O^+ densities, followed by the electron and ion omnidirectional differential energy fluxes from
160 FPI (Figure 2g-h). The last four panels present the differential fluxes of four individual ion species, H^+ ,
161 O^+ , He^+ and He^{++} measured by HPCA, respectively. The HPCA flux in panels 2i-l has artificial striping
162 every 4 energy bins due to way HPCA determines the count rate over 4 energy channels in survey mode.
163 It is noted that the differential fluxes (Figure 2i-l) and differential energy fluxes (Figures 2g-h) have
164 different units. To better identify the fluxes variations at specific energies, we choose the ion and electron
165 fluxes from FPI in the energy flux unit. The plasma moments (e.g. ion parallel and perpendicular
166 temperatures, ion and electron densities, and ion velocity) from FPI shown in Figures 2c-e are all from
167 MMS L2 data products. They are default moments calculated over the full FPI energy range from 10 eV
168 to 30 keV. Note that in the magnetosheath, O^+ measurements suffer from fake counts at energies below
169 1 keV which results from high proton fluxes contamination, as the red box in Figure 2j shown. So the
170 spurious counts should be excluded in the plasma moments calculation. The O^+ density shown in Figure
171 2f is recalculated from HPCA distribution functions at energies from 1 keV to 40 keV. Due to H^+
172 measurements from HPCA is accurate and the H^+ mean energy in the magnetosheath is typically 0.3 keV,
173 we adopted the default H^+ density from HPCA L2 data products which computed over the full HPCA
174 energy range from 1 eV to 40 keV, as the red line shown in Figure 2f.

175 The different regions encountered by MMS4 during the interval of 15:00 to 16:00 UT are marked
176 by the colored bar at the top of Figure 2, with the magnetosheath shown in orange, the outer
177 magnetosphere shown in blue, and the magnetopause boundary layer shown in green. From 15:00:00 to
178 15:25:10 UT, MMS4 was located in the magnetosheath. This region is characterized by the southward
179 magnetic field, low ion and electron temperatures (a few hundred eV for ions and tens of eV for electrons,
180 Figure 2c) with relatively high densities (on the order of $\sim 20 \text{ cm}^{-3}$, Figure 2d), and stable ion flow speed
181 of about 100 km/s. There are also very high fluxes at energies centered around 100 eV (nominal
182 magnetosheath energy) for electrons (Figure 2g) and at energies centered around 1 keV for ions (Figure
183 2h, also see H^+ fluxes in Figure 2i and He^{++} fluxes in Figure 2l) in the magnetosheath. While the O^+ and
184 He^+ fluxes above 1 keV nearly disappear in the magnetosheath (Figure 2j and 2k). From Figure 2j, the
185 majority of the O^+ fluxes at energies below 1 keV visible in the magnetosheath are the result of
186 contamination from the high proton fluxes, as the red box indicated.

187 The primary magnetopause crossing from the magnetosheath into the magnetosphere lasts about 12
188 min, from about 15:25:10 to 15:36:50 UT. Partial encounters of the magnetopause by MMS4 occurred
189 around 15:43:15, 15:47:10, 15:53:00 UT and etc. The magnetopause boundary layer is identified by the
190 plasma moments and the electromagnetic field. The plasma density and temperature at the magnetopause
191 are between the corresponding values of the magnetosphere and the magnetosheath, as shown in Figure
192 2d and 2c. The magnetopause boundary layer can also be identified by the significant increases in
193 electron fluxes at energies about several hundred eV and ion fluxes at energies around $\sim 10 \text{ keV}$, as shown
194 in Figure 2g and 2h, respectively. During this time of interest, the B_z component rotated from southward
195 to northward and back again several times before finally became northward when MMS 4 entered the
196 magnetosphere. The energetic O^+ density (1-40 keV) is around 0.018 cm^{-3} within the magnetopause
197 boundary layer as shown in Figure 2f. The corresponding H^+ and O^+ fluxes at specific energies and their
198 densities (shown in Figure 2f) were averaged in this region.

199 After 15:36:50 UT, MMS4 entered the magnetosphere which is identified by the observations of
200 the northward magnetic field (Figure 2a), much lower plasma densities (on the order of $\sim 1 \text{ cm}^{-3}$) with
201 respect to the densities in the magnetosheath (Figure 2d), higher plasma temperatures (Figure 2c, several
202 keV for ions and a few hundred eV for electrons), and a small bulk ion flow speed. Higher fluxes at
203 energies around several keV for electrons (Figure 2g) and at energies centered around $\sim 10 \text{ keV}$ for ions
204 (Figure 2h) also indicate that the MMS4 was in the magnetosphere. Finally, the presence of O^+ and He^+

205 at energies about ~ 10 keV is also used as a marker to verify that MMS4 was in the magnetosphere (Figure
206 2j and 2k).

207 **3.2 Statistical 57 events of energetic O^+ at the duskside magnetopause during intense substorms**

208 Based on the in-situ measurements of the dayside magnetopause crossings by MMS satellites in
209 phase 1, we identified the duskside magnetopause crossing event (complete magnetopause crossing from
210 the magnetosheath to the magnetosphere, vice versa) from the summary plot in
211 <https://lasp.colorado.edu/mms/sdc/public/plots/>. Then we plotted the more detailed overview figure of
212 these events to identify the magnetopause boundary layer, as Figure 2 shown. Only events with AE index
213 larger than 500 nT during the magnetopause boundary layer crossings interval were selected. There are
214 57 events of the duskside magnetopause boundary layer crossing during intense substorm satisfied
215 with the above criterion. In our statistical study, the mean values of the H^+ and O^+ fluxes at specific
216 energies and their densities are calculated in the magnetopause boundary layer. Correspondingly, the
217 solar wind dynamic pressure, IMF B_y , B_z and AE index from the OMNI data system were averaged
218 during the magnetopause boundary layer crossing time interval, as the two blue dashed lines shown in
219 Figure 1. The phase of the substorm is determined from the variations of AU, AL and AE indexes, as
220 mentioned before. For better follow-on studies, we add more detail information about 57 energetic O^+
221 events into an appendix. From the appendix, we can easily draw the conclusion that the O^+ density in the
222 duskside magnetopause during the recovery phase (0.081 cm^{-3}) of intense substorm is larger than that
223 during the expansion phase (0.069 cm^{-3}).

224 Figure 3 displays the locations of 57 energetic O^+ events at the duskside magnetopause ($-5.7 R_E <$
225 $Z_{GSM} < 1.7 R_E$) during intense substorms projected into the XY_{GSM} plane. The blue curve line represents
226 the nominal magnetopause, which is obtained by the magnetopause model of Shue et al., (1998) when
227 the IMF B_z is about -3.21 nT and solar wind dynamic pressure (P_{sw}) is ~ 2.87 nPa (averaged for the 57
228 events). The diamond and circle represent the event at the duskside magnetopause during the intense
229 substorm expansion phase and recovery phase, respectively. The O^+ density and the O^+/H^+ density ratio
230 are shown by the colored diamonds and circles at the corresponding magnetopause locations in Figures
231 3a and 3b, respectively. Among the 57 events of energetic O^+ at the duskside magnetopause during
232 intense substorms, there are 26 events that occurred during the expansion phase of intense substorms and

233 31 events occurred during the recovery phase. The maximum density of energetic O⁺ is found during the
 234 intense substorm recovery phase, as presented in Figure 3a.

235 Figure 4 presents the relationship between the energetic O⁺ at the duskside magnetopause and AE
 236 index during intense substorms. From top to bottom, panels show the O⁺ and H⁺ densities (Figure 4a),
 237 the O⁺/H⁺ density ratio (Figure 4b) and O⁺/H⁺ particle fluxes ratios at different energy ranges (Figure 4c),
 238 respectively. The energy channel ranges for O⁺ and H⁺ in the HPCA are the same. So the O⁺/H⁺ particle
 239 fluxes ratio is directly defined as the ratio between mean values of their fluxes, respectively. The particle
 240 fluxes are chosen at energies ~1 keV (energy range from 987.82 to 1165.21 eV), ~10 keV (energy range
 241 from 9.97 to 11.77 keV), ~20 keV (energy range from 19.31 to 22.78 keV) and ~35 keV (energy range
 242 from 31.69 to 37.39 keV). The error bars indicating the 90% confidence interval (CI) are also overplotted
 243 in each point. The confidence interval is based on the following formula:

$$244 \quad \bar{x} - k \frac{s}{\sqrt{n}} < \mu < \bar{x} + k \frac{s}{\sqrt{n}}$$

245 Where \bar{x} , s and n are the mean value, standard deviation and the sampling number of observations,
 246 respectively. k in the above formula can be determined by calculating a 90% confidence interval for each
 247 events (the k value is 1.65). Figure 4a shows that energetic O⁺ density at the duskside magnetopause
 248 during intense substorms is in the range from 0.007 to 0.599 cm⁻³. The maximum density value of
 249 energetic O⁺ at the duskside magnetopause during intense substorm recovery phase is presented at the
 250 higher AE index about 606 nT. The O⁺/H⁺ density ratio decreases with AE index from 900 to 1100 nT.
 251 The variations of O⁺ density and O⁺/H⁺ density ratio with AE index do not show obvious difference
 252 between during the expansion phase and the recovery phase.

253 Figure 5 shows the relationship between the energetic O⁺ at the duskside magnetopause and IMF
 254 B_y during intense substorms. The format of Figure 5 is the same as that of Figure 4. Figure 5a shows that
 255 the O⁺ and H⁺ densities decrease with IMF B_y from -6 to 0 nT and increase with IMF B_y from 4 to 8 nT.
 256 From Figure 5b, the O⁺/H⁺ density ratio shows an exponential growth with the IMF B_y. Based on the
 257 scatter plot in Figure 5b, we can define linear functional dependence between the logarithm of O⁺/H⁺
 258 density ratio and IMF B_y, as Eq. (1) shown. And the corresponding correlation coefficients is 94%. The
 259 correlation coefficient close to 100% indicates that there is a great correlation.

$$260 \quad \log \frac{n_{O^+}}{n_{H^+}} = 0.126 * \text{IMF } B_y - 5.174 \quad (1)$$

261 The dependency is constructed using a linear least-squares fit unless otherwise stated. The O⁺/H⁺
262 particle flux ratio at energy ~10 keV, ~20 keV and ~35 keV also show an obvious exponential increase
263 with IMF B_y. This dependency is consistent with Welling et al. (2011) simulation results found in the
264 ring current.

265 Figure 6 shows the relationship between the energetic O⁺ at the duskside magnetopause and IMF
266 B_z during intense substorms. The format of Figure 6 is the same as that of Figure 4. Figure 6a and 6b
267 both present that among 57 events of energetic O⁺ at the duskside magnetopause boundary layer during
268 intense substorm, there are 50 events under southward IMF and only 7 events under northward IMF. It
269 is noted that 26 events occurred during the expansion phase of intense substorms which are all under the
270 southward IMF conditions, as the blue points shown. Meanwhile, the events that occurred under
271 northward IMF are all during the intense substorm recovery phase, as the right red points with positive
272 IMF B_z shown. From -10 to 0 nT, the O⁺ density shows an obvious decrease with IMF B_z. To better
273 describe this variation trend, the empirical functional relation between the logarithm of O⁺ density and
274 IMF B_z (from -10 to 0 nT) is established in Eq.(2) and the corresponding correlation coefficient is 94%.
275 While the O⁺ density has a positive correlation with IMF B_z from 0 to 5 nT.

$$276 \quad \log n_{O^+} = -0.163 * IMF B_z - 3.737 \quad (2)$$

277 From Figure 6b, the O⁺/H⁺ density ratio during the recovery phase decrease with IMF B_z from about -2
278 to 2 nT. The maximum density of energetic O⁺ at the duskside magnetopause is under southward IMF.
279 Meanwhile, the maximum O⁺/H⁺ density ratio at the duskside magnetopause is also under southward
280 IMF.

281 Figure 7 displays the relationship between the energetic O⁺ at the duskside magnetopause and solar
282 wind dynamic pressure during intense substorms. The format of Figure 7 is the same as that of Figure 4.
283 Figure 7a presents that the O⁺ density at the duskside magnetopause during intense substorms has a
284 positive correlation with the solar wind dynamic pressure. The empirical functional relation between the
285 logarithm of O⁺ density and solar wind dynamic pressure (from 1 to 4.5 nPa) is also established in Eq.(3)
286 and the corresponding correlation coefficient is 94%.

$$287 \quad \log n_{O^+} = 0.325 * P_{sw} - 4.061 \quad (3)$$

288 From Figure 7b, the O⁺/H⁺ density ratio during recovery phase show a decrease from about 2.5 to 3 nPa.
289 It is noted that the O⁺/H⁺ density ratio increase with solar wind dynamic pressure from about 3 to 4 nPa.
290 The maximum density of energetic O⁺ at the duskside magnetopause, ~0.599 cm⁻³ take place at solar

291 wind dynamic pressure is about 3.9 nPa. While the maximum O⁺/H⁺ density ratio at the duskside
292 magnetopause appeared at solar wind dynamic pressure around 2.2 nPa. More details can be found in the
293 appendix.

294 **4 Discussion**

295 Energetic O⁺ (1-40 keV) with high density are observed by MMS satellites at the duskside
296 magnetopause during the expansion phases and recovery phases of intense substorms, as displayed in
297 Figure 3a. The density of energetic O⁺ is in range from 0.007 cm⁻³ to 0.599 cm⁻³ at the duskside
298 magnetopause boundary layer during intense substorms. In a companion paper from Zeng et al. (2019),
299 they study the O⁺ abundance variations on the solar wind conditions at the dayside magnetopause
300 boundary layer and not specific to the events that occurred during intense substorm. The mean value of
301 the O⁺ density at the duskside magnetopause boundary layer is 0.038 cm⁻³ in that paper. While during the
302 intense substorm, the O⁺ density increase to 0.075 cm⁻³ in this study. There are two reasons for this high
303 density of energetic O⁺ observed during the intense substorm. The first is the time interval for the
304 observations. Our observations are during intense substorms expansion phase and recovery phase. Daglis
305 et al. (1991) proposed that energetic O⁺ were significantly higher in the NEPS in the magnetotail after
306 intense substorms onset. The impulsive electric field accompanied by intense substorm dipolarization
307 plays a key role in the energization and sunward transfer of oxygen ions in the duskside of midnight
308 plasma sheet in the magnetotail (e.g., Fok et al., 2006; Nosé et al., 2000). These energetic O⁺ (tens of
309 keV) can be transported sunward into the duskside magnetopause boundary layer. The second reason for
310 the high densities is the locations of the observations. Our observations are near the duskside
311 magnetopause. This region is easily accessible by energetic O⁺ during intense geomagnetic activity
312 (Fuselier et al. 2016a). Phan et al. (2004) pointed out that energetic O⁺ with very high density 0.2-0.3
313 cm⁻³ in the reconnection jets at the duskside mid-latitude magnetopause were observed by Cluster.

314 During dynamic periods and intense substorms time, light ions yielded more symmetric patterns in
315 density than heavy ions and the O⁺ patterns in the active plasma sheet are a function of IMF conditions
316 (Winglee and Harnett 2011. Winglee et al. 2005). Welling et al. (2011) used multispecies MHD and the
317 PWOM to drive a ring current model and found that positive IMF B_y pushed the stronger O⁺
318 concentrations toward the duskside at a geocentric distance of about 6.6 R_E. This O⁺ density duskward

319 preference with positive IMF B_y in the NEPS is similar to our results. It may indicate that O^+ in the
320 magnetopause boundary layer enhancing with IMF B_y is due to the local time variations of O^+ in the
321 NEPS tied to IMF B_y . Our result of O^+ density increase with IMF B_y also agree with Kronberg et al.,
322 (2012). They showed for 10 keV O^+ strong increasing under the duskward IMF indicated by the clock
323 angle in the inner magnetosphere. It is suggested that the O^+ abundance at the duskside magnetopause
324 has a corresponding relation with the O^+ in the duskside near-Earth magnetosphere during intense
325 substorm. The O^+ path from the cusp to the magnetotail is asymmetric and it has the best correlation with
326 the IMF directions. This path asymmetry mainly controlled by the IMF B_y may influence the O^+
327 abundance at the duskside magnetopause. When the IMF B_y is positive, the O^+ from northern/southern
328 cusp tends to flow toward the dawnside/duskside. The transport path for negative IMF B_y is more
329 symmetric but shows some evidence for a reversed asymmetry when the negative IMF B_y is large enough.
330 While the IMF B_z has little influence on the asymmetry (Liao et al., 2010).

331 Due to not enough events occurred under northward IMF were observed, the influence of IMF B_z
332 on the O^+ abundance (1-40 keV) during intense substorm is not clear. While Luo et al. (2017) found that
333 the O^+ intensity ($> \sim 274$ keV) was significantly higher under southward IMF than that under northward
334 IMF, especially at the duskside magnetopause. Zeng et al. (2019) also showed that the duskside
335 asymmetry of O^+ density (1-40 keV) in the dayside magnetopause under northward IMF was less obvious
336 than under southward IMF when the IMF B_y was the same. Under the southward IMF, the interactions
337 between the solar wind and the magnetosphere become active. The inductive electric field or magnetic
338 field gradient related to magnetic reconfiguration will enhance with negative IMF B_z . So the large scale
339 dawn-dusk electric field drift along with the gradient-curvature drift can force oxygen ions convect to
340 the duskside magnetopause boundary layer (Kronberg et al., 2015; Luo et al., 2017).

341 In this statistical study, there are 50 magnetopause boundary layer crossing events during intense
342 substorm under southward IMF with respect to 7 events under northward IMF. Choosing the intense
343 substorm may increase the probability of observing the events under southward IMF quite significantly.
344 Among 57 events of energetic O^+ near the duskside magnetopause, there are 26 events during intense
345 substorm expansion phase which are all under the southward IMF, as the blue circle shown in Figure
346 6b. There are only 7 events under northward IMF in our study and they all occurred during the intense
347 substorm recovery phase. But what relation between the IMF directions and phase of substorm is out of
348 scope for this article.

349 Previous studies demonstrated that the oxygen ions that originate from the aurora region could
350 rapidly feed in the NEPS during intense substorms expansion phase (e.g., Daglis and Axford, 1996; Duan
351 et al., 2017; Yu et al., 2013). Oxygen ions can be efficiently energized in the NEPS during intense
352 substorm dipolarization (e.g., Duan et al., 2017; Fok et al., 2006; Nosé et al., 2000). Under southward
353 IMF conditions, these energetic oxygen ions in the NEPS can be convected sunward and drift westward.
354 As a result, the energetic O^+ arrived near the duskside magnetopause can participate in the magnetopause
355 reconnection and escape along reconnected field lines during intense substorm expansion phase, as
356 reported by Wang et al. (2014) and Zong et al. (2001). When O^+ participate in the reconnection jets, the
357 reconnection rate will likely be reduced by the mass-loading but not suppressed at the magnetopause
358 (Fuselier et al. 2019). Whether these energetic O^+ at the duskside boundary layer could suppress the
359 intense substorm need further investigation.

360 **5 Summary and conclusions**

361 Using the measurements from MMS satellite during the phase 1, we have studied 57 events of the
362 energetic O^+ (1-40 keV) at the duskside magnetopause boundary layer and their variations on the solar
363 wind conditions (IMF B_y , IMF B_z and solar wind dynamic pressure) during intense substorm expansion
364 phases and recovery phases. According to the above analysis, we can draw our main conclusions as
365 follows. In our 57 events of energetic O^+ at the duskside magnetopause boundary layer, there are 26
366 events during the expansion phase of intense substorms and 31 events during the recovery phase. It is
367 noted that the mean values of the O^+ density during the expansion phase and recovery phases are 0.069
368 cm^{-3} and 0.081 cm^{-3} , respectively. And the maximum O^+/H^+ density ratio occurred during the intense
369 substorm recovery phase. It is found that 26 events of energetic O^+ at the duskside magnetopause during
370 intense substorms expansion phase are all under the southward IMF conditions, and only 7 events under
371 northward IMF which are all during the intense substorm recovery phase. The O^+ density shows an
372 exponential increase with IMF B_z absolute value under the southward IMF. Similarly, it also presents an
373 exponential growth with solar wind dynamic pressure, and the empirical functional relations are
374 established. Like previous studies during substorm in the near-Earth magnetosphere, The O^+/H^+ density
375 ratio in the duskside magnetopause boundary layer enhance with the IMF B_y . It is suggested that the O^+

376 abundance in the duskside magnetopause boundary layer has a close correlation with the O⁺ variations
377 in the near-Earth magnetosphere during intense substorm.

378 **Data availability**

379 All data used in this study are publicly accessible. MMS data are available at the MMS Science
380 Data Center (<https://lasp.colorado.edu/mms/sdc/public/>). The OMNI data can be downloaded from the
381 NASA Goddard Space Flight Center Coordinated Data Analysis Web
382 (CDAWeb:<http://cdaweb.gsfc.nasa.gov/>).

383 **Competing interests**

384 The authors declare that they have no conflict of interest.

385 **Author contribution**

386 C. Z. conducted the majority of the data processing, analysis and writing for this study. S.P.D, C.W,
387 L.D and S.F participated in the interpretation of the data and modified this paper. J.B, R.T, B.G and C.R
388 produced the data and controlled the data quality. All the authors discussed the results and commented
389 on the paper.

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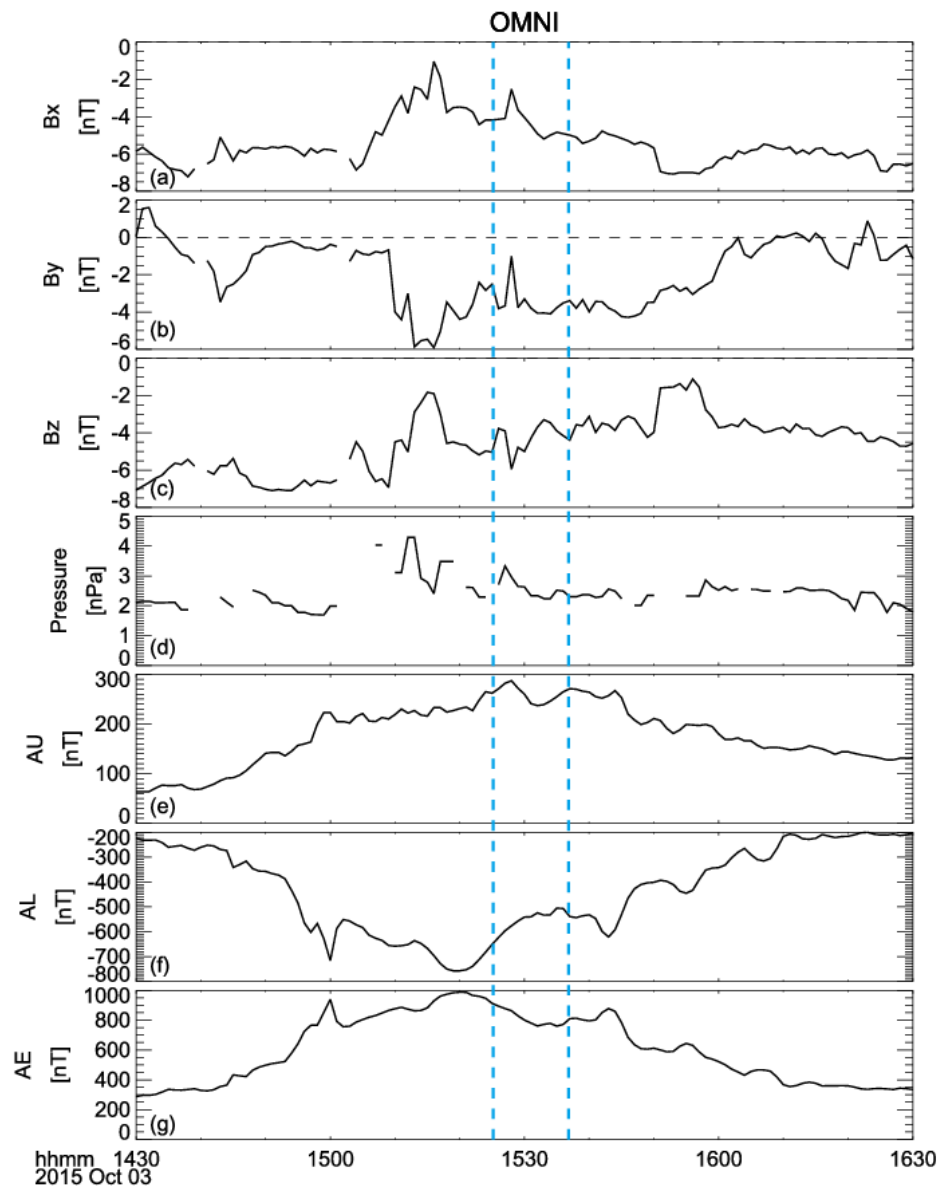
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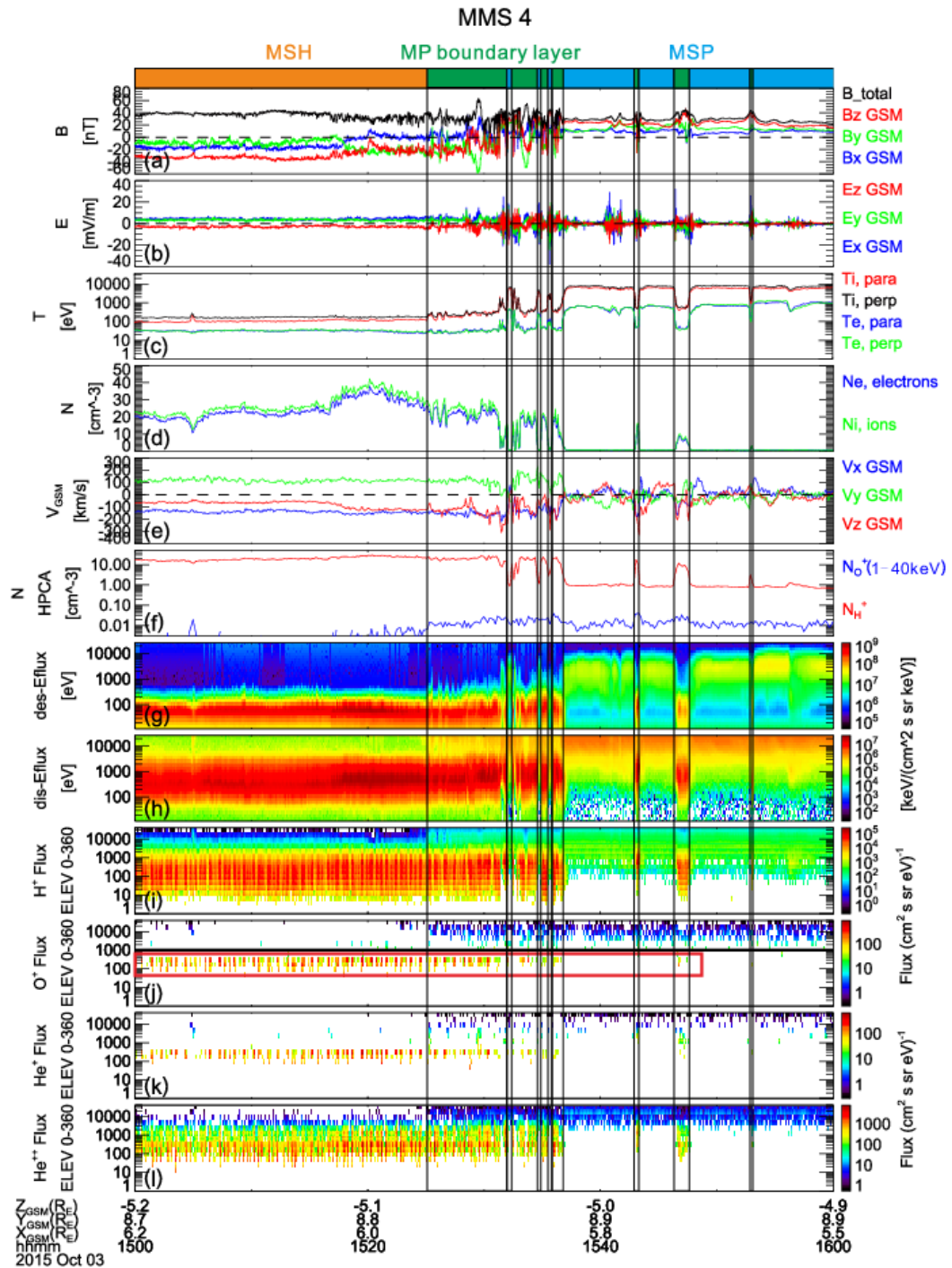
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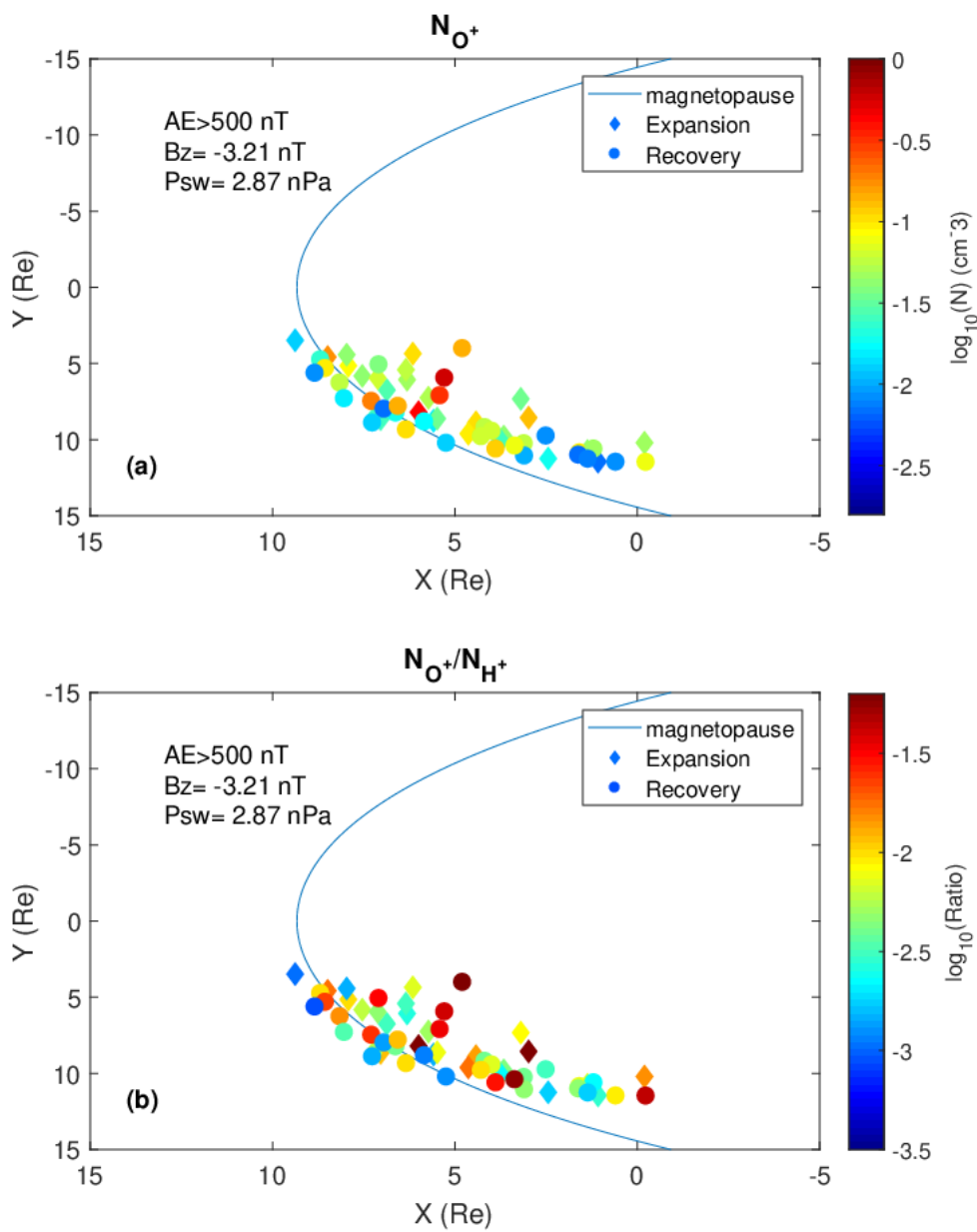
622 **Figure 1.** The three components of the IMF Bx, By, Bz in the Geocentric Solar Magnetospheric
 623 coordinates, solar wind dynamic pressure, as well as AU, AL and AE index from CDAweb OMNI data.
 624 The two blue dashed lines indicate the interval of the magnetopause boundary layer crossing.



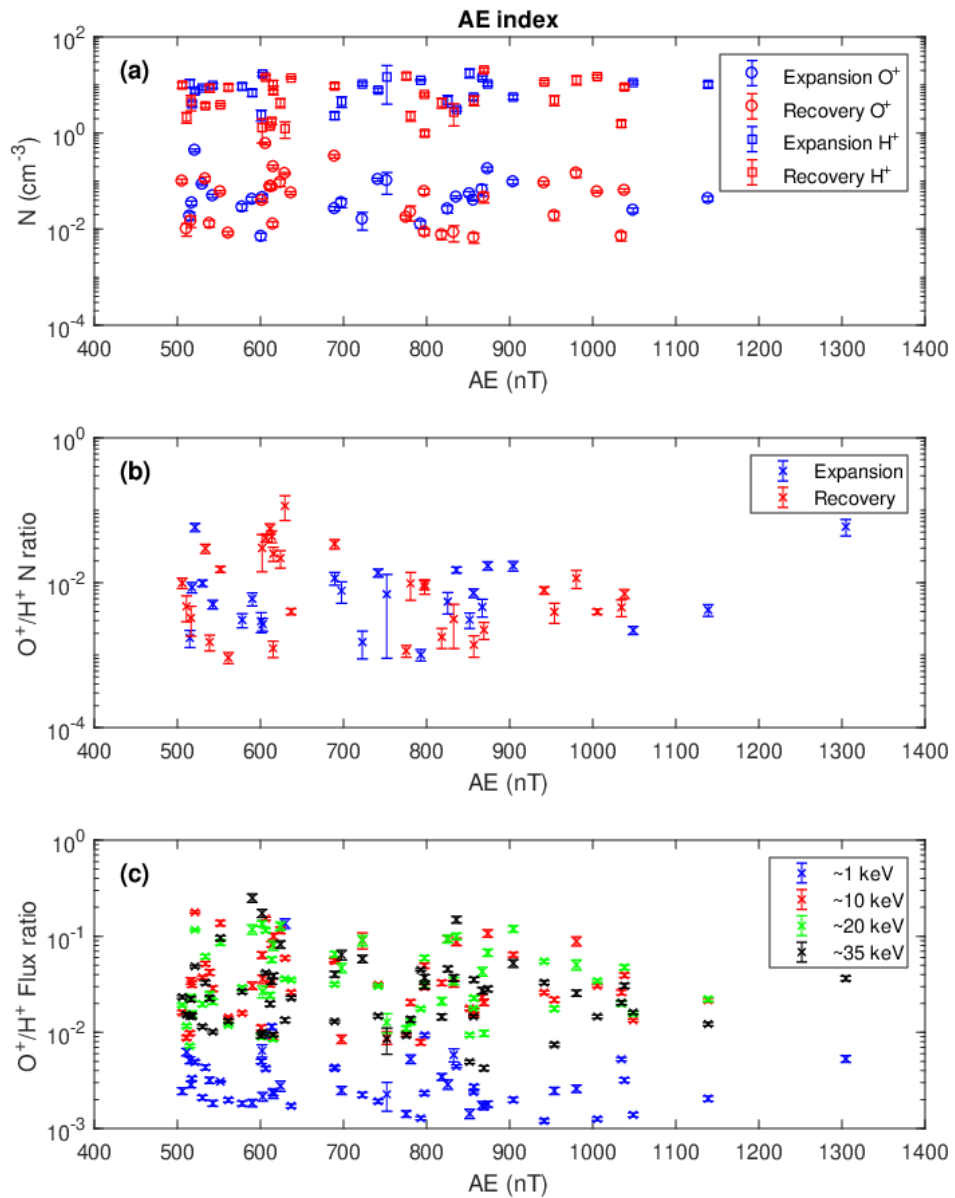
625

626 **Figure 2.** The energetic O⁺ is observed at the magnetopause during an intense substorm on 03 October
 627 2015 by MMS 4. From top to bottom are (a) the magnetic field three components, B_x (blue line), B_y
 628 (green line), B_z (red line) and the total magnitude B_t (black line), (b) the electric field three components,
 629 E_x (blue), E_y (green) and E_z (red), (c) ion parallel (red) and perpendicular (black) temperatures, as well

630 as electron parallel (blue) and perpendicular (green) temperatures, (d) The density of ion (green) and
 631 and electron (blue), (e) three components of the ion velocity, (f) the H^+ (over the full HPCA energy range
 632 from 1 eV to 40 keV) and O^+ (at energies from 1 keV to 40 keV) densities, (g-h) electron and ion
 633 omnidirectional differential energy fluxes ($\text{keV}/(\text{cm}^2 \text{ sr KeV}^{-1})$), (i) to (l) present differential particle
 634 fluxes ($\text{cm}^2 \text{ s sr eV}^{-1}$) of H^+ , O^+ , He^+ , He^{++} , respectively. The Geocentric Solar Magnetospheric coordinate
 635 system is adopted. The thick bars at the top of the panel present different regions encountered on this
 636 magnetopause crossing event. The orange and blue bars represent the magnetosheath and the
 637 magnetosphere, respectively. The green bar represents the magnetopause boundary layer. The black
 638 horizontal line in figure 2j is at 1 keV and the O^+ contamination from high H^+ fluxes is indicated by the
 639 red box. The FPI data in Figure 2c-e and g-h are from FPI L2 data products and in the fast mode.



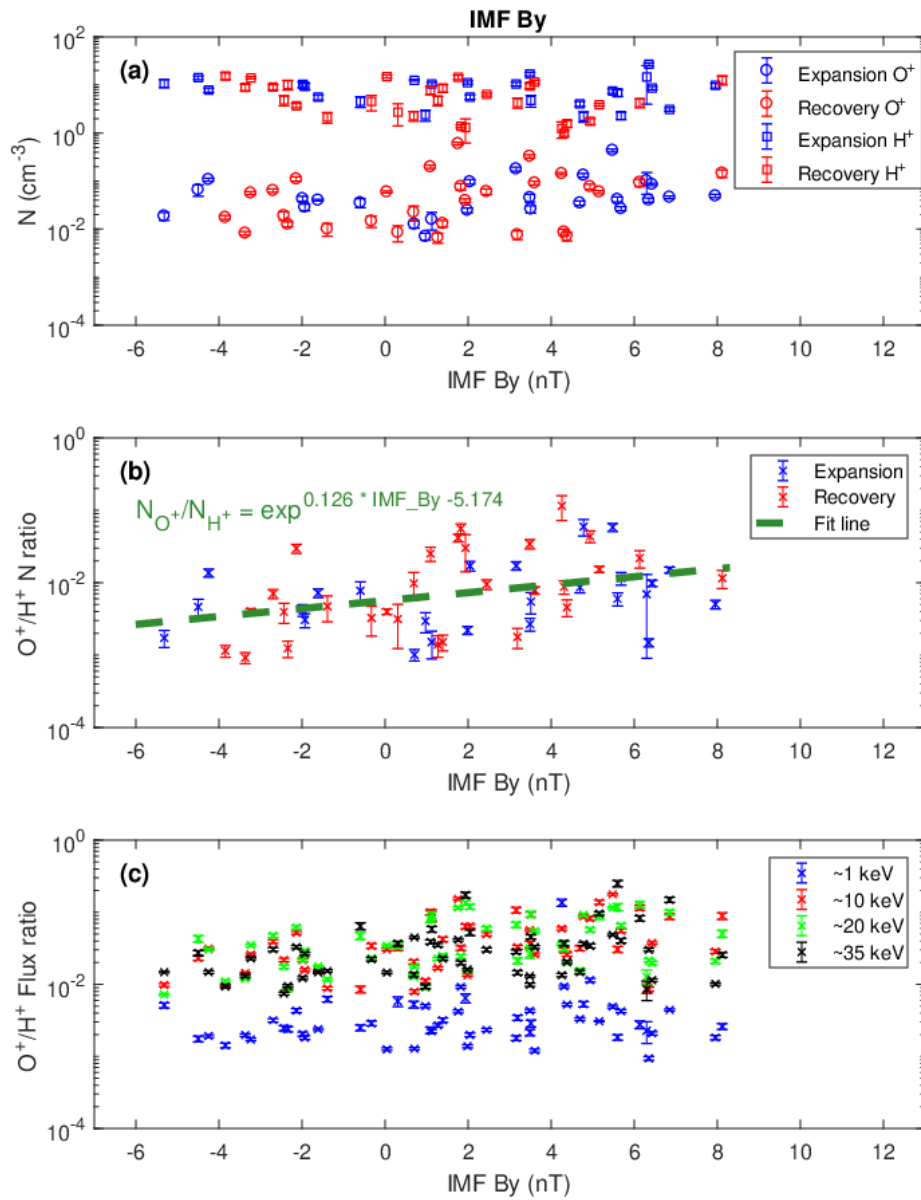
641 **Figure 3.** Maps of 57 events of energetic O^+ at the duskside magnetopause during intense substorms with
 642 AE index larger than 500 nT in XY_{GSM} plane. The O^+ density and the density ratio of O^+/H^+ are shown
 643 by the color signatures at the corresponding magnetopause location in Figure 3a and 3b, respectively.
 644 The blue curve line represents the nominal magnetopause. The diamond and circle represent the event at
 645 the magnetopause during the intense substorm expansion phase and recovery phase, respectively.



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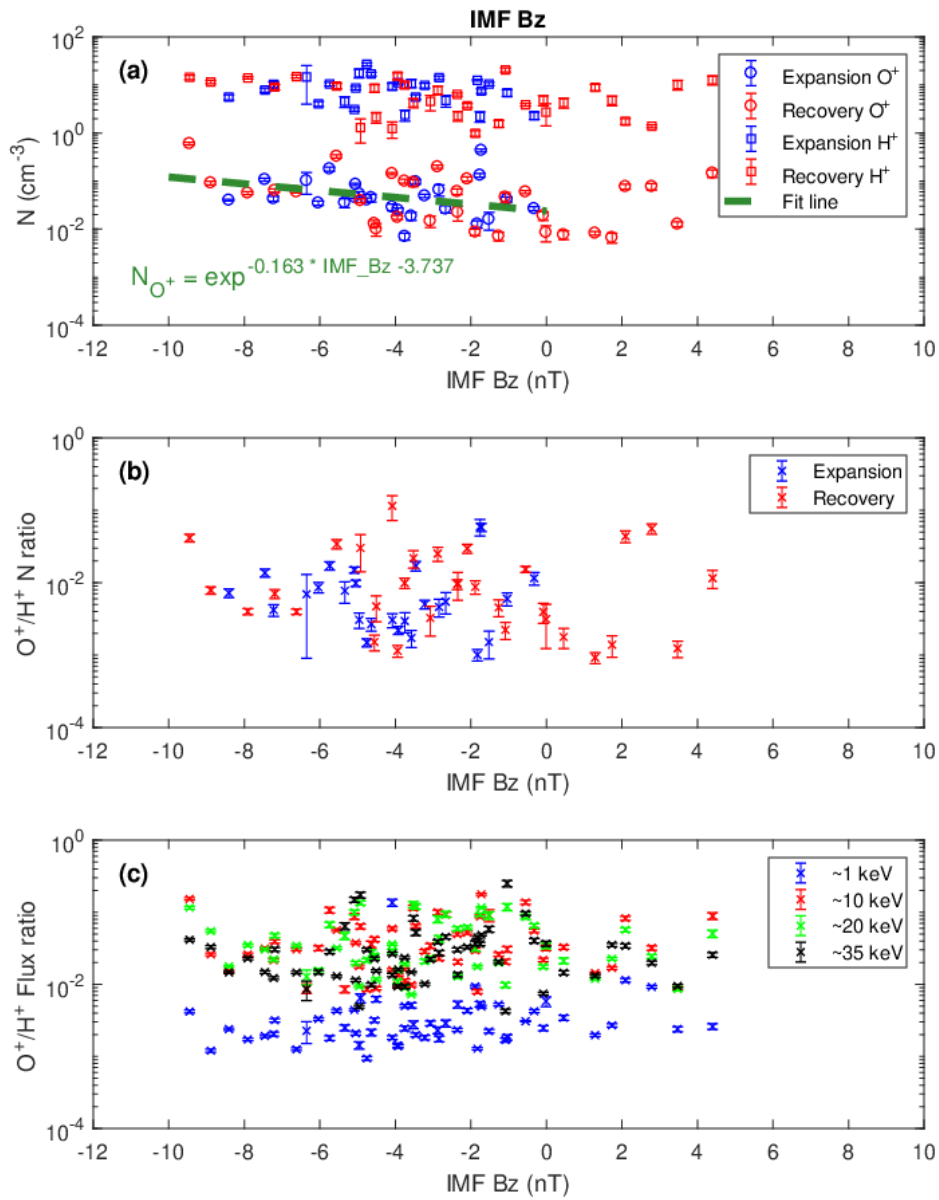
647 **Figure 4.** The relationship between the energetic O^+ at the duskside magnetopause and AE index during
 648 intense substorms. From top to bottom, panels show the O^+ and H^+ densities (Figure 4a), the O^+/H^+

649 density ratio (Figure 4b), and the O^+/H^+ particle flux ratio (Figure 4c), respectively. Error bars indicate
650 90% confidence intervals.

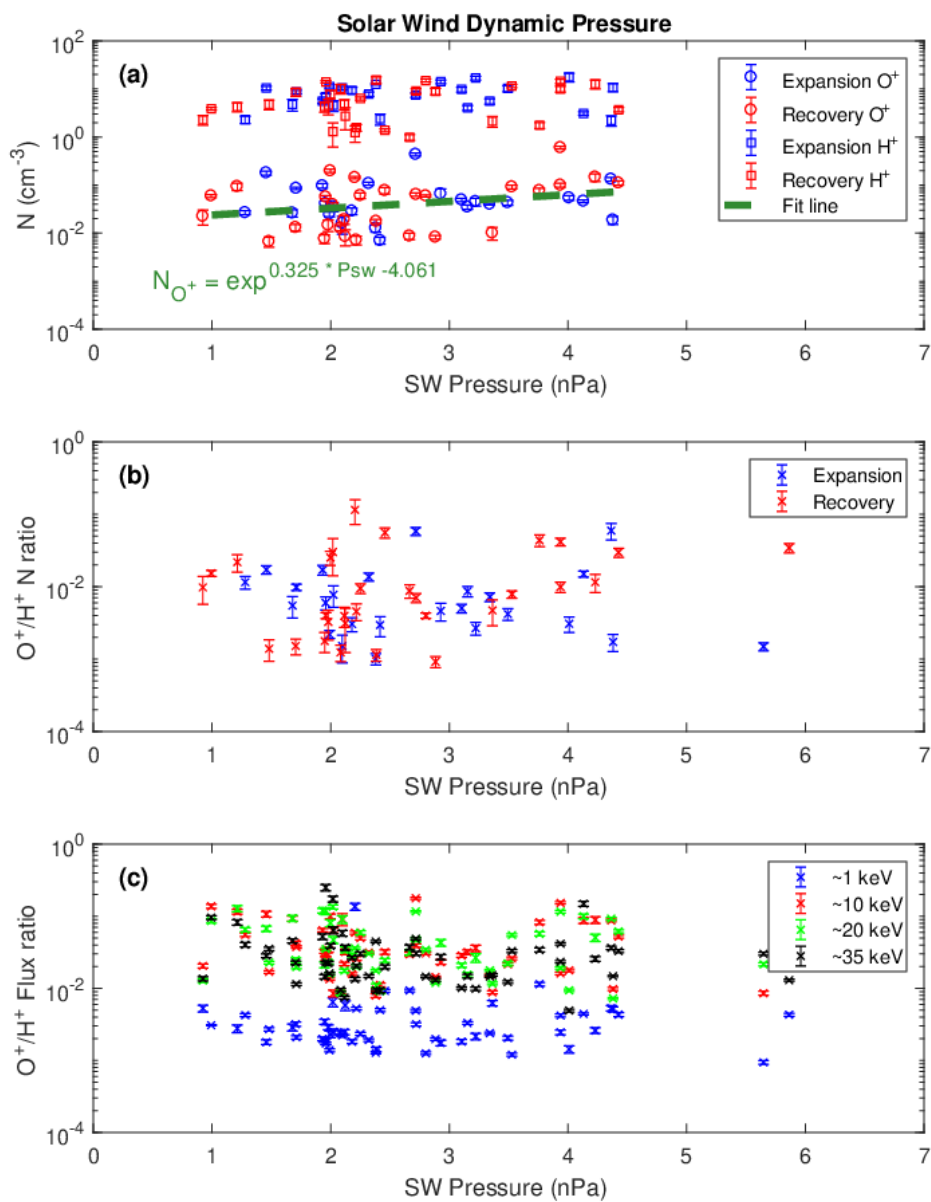


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652 **Figure 5.** The relationship between the energetic O⁺ at the duskside magnetopause and IMF B_y during
 653 intense substorms. The format is the same as that of Figure 4.



655 **Figure 6.** The relationship between the energetic O⁺ at the duskside magnetopause and IMF Bz during
 656 intense substorms. The format is the same as that of Figure 4.



657

658 **Figure 7.** The relationship between the energetic O⁺ at the duskside magnetopause and solar wind
 659 dynamic pressure during intense substorms. The format is the same as that of Figure 4.