#### **Reply to reviewer #1**

Dear reviewer:

We are very grateful for your comments about our manuscript. We acknowledge the reviewer's comments and suggestions very much, which are valuable in improving the quality of our manuscript. According to your advice, we amended the relevant part of the manuscript. The one-to-one responses are the following.

**Comment 1:** Line 138: The authors should clarify the information on which instrument dataset was used for each data product. Were the moments shown in Figure 2c-2e recalculated from the FPI distribution functions? Or are they the default moments calculated over the full FPI energy range?

**Response:** Thanks for the referee's kind advice. We should have clarified the dataset information in Figure 2 on which instrument was used. So I added detailed information about dataset what we used to the description of Figure 2. The plasma moments (e.g. Ion parallel and perpendicular temperatures, ion, and electron densities and ion velocity) from FPI shown in Figure 2c-2e are all from MMS L2 data products. They are default moments calculated over the full FPI energy range from 10 eV to 30 keV. But the O<sup>+</sup> density shown in Figure 2f is recalculated from HPCA distribution functions at energies from 1 to 40 keV. From the O<sup>+</sup> fluxes shown in Figure 2j, there still exists a large number of fluxes below 1 keV in the magnetosheath. This part of O<sup>+</sup> fluxes is fake and contamination from high proton fluxes. So we consider the number density of O<sup>+</sup> at energies from 1 to 40 keV is more appropriate to represent the true O<sup>+</sup> in the magnetopause. While the H<sup>+</sup> density (over the full HPCA energy range) from L2 data products are used in Figure 2f. The magnetic and electric fields in GSM are from FGM and EDP, respectively. The last four panels of Figure 2 show the omnidirectional differential fluxes of four individual ion species, H<sup>+</sup>, O<sup>+</sup>, He<sup>+,</sup> and He<sup>++</sup> measured by HPCA, respectively. See Line 198-207 in "Tracked change" manuscript.

**Comment 2:** Line 140 (Figure 2f): The calculations performed to derive the >1 keV O<sup>+</sup> density need to be described to inform the reader how the HPCA energy ranges were specified for those calculations. If a software package was used, then details of the software package and a citation to it should be included. The >1keV H<sup>+</sup> density could also be plotted in this figure panel.

**Response:** Thanks for the referee's good suggestion. As mentioned before, the >1 keV O<sup>+</sup> density (shown in Figure 2f) recalculated from the HPCA distribution functions at energies from 1 to 40 keV. As your suggestion, I also plotted the H<sup>+</sup> density over the full FPI energy range from 10 eV to 40 keV in Figure 2f for better comparison. Because of H<sup>+</sup> measurements from HPCA is accurate and the H<sup>+</sup> mean energy in the magnetosheath is typically 0.3 keV, so we used the H<sup>+</sup> density (over the full HPCA energy range) from L2 data products. These O<sup>+</sup> density calculations are used The Space Physics Environment Data Analysis System (SPEDAS) software package. More details about SPEDAS can be found in Angelopoulos et al. (2019) and cited as (Angelopoulos, V., Cruce, P., Drozdov, A. et al. Space Sci Rev (2019) 215: 9. https://doi.org/10.1007/s11214-018-0576-4). We also cited this paper in our revised manuscript, see Line 144-148 and Figure 2f in "Tracked change" manuscript.



**Figure 2.** The energetic  $O^+$  are observed at the magnetopause during an intense substorm on 03 October 2015 by MMS 4. From top to bottom are (a) the magnetic field three components, Bx (blue line), By (gree line), Bz (red line) and the total magnitude Bt (black line), (b) the electric field three components, Ex (blue), Ey (gree) and Ez (red), (c) Ion parallel (red) and perpendicular(black) temperatures, (d) The number density of ion (green) and electron (blue), (e) three components of the ion velocity, (f) number density of H<sup>+</sup> (over the full HPCA energy range) and O<sup>+</sup> (at energies from 1 to 40 keV), (g) electron omnidirectional differential energy fluxes, (h) ion omnidirectional differential energy fluxes, (i) to (l) present omnidirectional differential particle fluxes of H<sup>+</sup>, O<sup>+</sup>, He<sup>+</sup>, and He<sup>++</sup>, respectively. The Geocentric Solar Magnetospheric (GSM) coordinate system is adopted. The thick bars at the top of the panel represent different regions encountered on this magnetopause crossing event. The orange and blue bars represent the magnetosheath and the magnetosphere, respectively. The green bar represents the magnetopause boundary layer. The black horizontal line in figure 2j is at 1 keV and the O<sup>+</sup> contamination from high H<sup>+</sup> fluxes is indicated by the red box. The FPI data in Figure 2c-e and 2g-h are from FPI L2 data product and in the fast mode.

**Comment 3:** Line 158-164: The magnetopause identification criteria are not very convincing. Recommend carefully defining these criteria, as all statistics are derived based on the magnetopause identification. Recommend the authors review identification criteria used in previous works. For example, Haaland et al. (2016) and (2019) describe magnetopause observations by Cluster and THEMIS: Haaland, S., Reistad, J., Tenfjord, P., Gjerloev, J., Maes, L., DeKeyser, J., Maggiolo,R., Anekallu, C., and Dorville, N. (2014), Characteristics of the flank magnetopause:Cluster observations, J. Geophys. Res. Space Physics, 119, 9019–9037,doi:10.1002/2014JA020539.

Haaland, S., Runov, A., Artemyev, A., & Angelopoulos, V. (2019). Characteristics of the flank magnetopause: THEMIS observations. Journal of Geophysical Research: Space Physics, 124, 3421–3435. <u>https://doi.org/10.1029/2019JA026459</u>. Paschmann et al. (2018) describe magnetopause identification and observations by MMS:

Paschmann, G., Haaland, S. E., Phan, T. D., Sonnerup, B. U. Ö., Burch, J. L., Torbert, R. B., et al. (2018). Large-scale survey of the structure of the dayside magnetopause by MMS. Journal of Geophysical Research: Space Physics, 123, 2018–2033. <u>https://doi.org/10.1002/2017JA025121</u>.

**Response:** Thanks for the referee's kind suggestion and well recommend. We read the papers your recommended and found they did detailed work for magnetopause identification. It deepens my understanding of the flank magnetopause characteristics and helps me identifying the magnetopause more convincing. In this study, we mainly focus on the O<sup>+</sup> in the dusk flank magnetopause boundary layer. According to the magnetic field, B is about 40 nT and O<sup>+</sup> temperature, T is about 10 keV in this study, we can draw the O<sup>+</sup> gyroradius is about 1020 km. From Haaland et al. (2014), the flank magnetopause thickness varies from 150 to 5000 km with a median thickness of around 1150 at dusk. Thus the gyroradius of ten keV  $O^+$  is comparable to magnetopause thickness. In that situation,  $O^+$  will show the finite Larmor radius effects and the MMS detect partial gyro motion in the magnetopause. For acquiring complete  $O^+$  distribution functions, we need to measure  $O^+$  in more large spatial scales. So in this study, we focus on the magnetopause boundary layer judgment. The magnetopause boundary layers are identified here primarily through plasma fluxes and moments. The low-latitude boundary layer (LLBL) on the magnetospheric side of the magnetopause current layer and the magnetosheath boundary layer (MSBL) on the magnetosheath side of the magnetopause current layer can have densities and temperatures between that of the magnetosphere and magnetosheath. Ion jets are also signatures of passing through the magnetopause boundary layers. In this study, the separatrix between the magnetosheath and the magnetopause boundary layer is determined by the appearance of the magnetospheric electron. Similarily, the separatrix of the magnetosphere and the magnetopause boundary layer is determined by the magnetosheath electron disappearance. The revised details can be found in Line 221-233 in "Tracked change" manuscript.

**Comment 4:** Line 180: More details are needed to describe how the mean values of the H<sup>+</sup> and O<sup>+</sup> fluxes and densities were calculated.

**Response:** Thanks for the referee's kind advice. First, we determine the time interval of the magnetopause boundary layer crossings in each event. For example, on 03 October 2015 event, MMS 4 traversed the duskside magnetopause boundary layer from 15:25:10 to 15:36:50 UT judged by the typical characteristics in this region as mentioned before. Then, the H<sup>+</sup> and O<sup>+</sup> fluxes and densities were average during this time interval. We also give the error bars indicating 90% confidence intervals. We think these mean values represent the H<sup>+</sup> and O<sup>+</sup> fluxes and densities in the magnetopause boundary layer. See Line 232-233 in "Tracked change" manuscript.

**Comment 5:** Line 184: A more detailed description of how the substorm phase (i.e. expansion phase or recovery phase) was defined based on AE index is needed. The authors should use Figure 1 AE index to aid in their description.

**Response:** Thanks for the referee's suggestion. We should give more details to clarify how we define the substorm phase according to substorm indices, such as AU, AL, and AE index. First, we determined the time interval of the magnetopause boundary layer crossings in each event. Then, see how the substorm indices vary during that interval from the OMNI data. As Figure 1 shown, the time interval of the magnetopause boundary layer crossing is marked by the two blue dashed lines. As we know, the AE index is defined as AE=AU-AL. Generally, the substorm onset time is characteristic by the AL index starts to significantly decrease and the AE index significantly increase. During the substorm expansion phase, the AL index will decrease significantly. The interval of the AL index decrease from onset to its minimum is defined as the substorm expansion phase. Then it starts to increase and the interval of the AL index increase from the minimum to the quiet time level is regarded as the substorm recovery phase. In our event, the MMS4 crossed the magnetopause boundary layer from 15:25:10 to 15:36:50 UT on 3 October 2015. From Figure 1f, the AL index reached its minimum ~-750 nT and AE index reach the peak ~1000 nT at about 15:20 UT, then it started to increase to ~ -200 nT at the rest time of interest. The two blue dashed lines indicate the time interval of the magnetopause boundary layer crossing. According to the variation and peak value of the AU, AL and AE index in Figure 1e to 1g. The magnetopause boundary layer crossing occurred during the recovery phase of this intense substorm. The revised details can be found in Line 166-180 in "Tracked change" manuscript.



**Figure 1.** The three components IMF Bx, By, Bz, solar wind dynamic pressure, as well as AU, AL, and AE index from CDAweb OMNI data.

**Comments 6:** Line 202-209: Several narrow energy ranges used for comparing the  $O^+/H^+$  density ratio are noted. It is important to describe for the reader how these energy ranges were used in the density ratio calculations. In addition, a description of why these energy ranges were chosen should be included. Did the authors consider calculating the density ratio for all energies >1 keV instead of calculating the ratios over individual energy ranges? A comparison of density ratios using both methods may be helpful to decide which method to use. Such procedural information on which analysis methodology was chosen could be included in an appendix.

**Response:** Thanks for carefully evaluating this manuscript and kind suggestions. The description in Line 202-209 is not accurate and it appears that the referee has some misunderstanding on what we did. In this study, we calculate the  $O^+/H^+$  density ratio (as Figure 4b shown). The  $O^+$  density calculated at energies from 1 keV to 40 keV, but the H<sup>+</sup> density (over the full HPCA energy range) from L2 data products are used. In order to realize in which individual energy ranges the  $O^+$  abundance ( $O^+/H^+$ ) varies obviously on AE index and solar wind parameters. We calculated the particle fluxes ratio at several individual energy ranges (as Figure 4c shown). Since the energy channel range of HPCA for H<sup>+</sup>

and  $O^+$  is the same, so the fluxes ratio are defined as the ratio between their fluxes, We also give the error bars indicating 90% confidence intervals. See Line 280-292 in "Tracked change" manuscript.

**Comments 7:** Line 218-248: Figures 5, 6, 7 all show comparisons of the  $O^+/H^+$  density ratio. After addressing the previous comment on Line 202-209 on why separate narrow energy ranges were chosen instead of using a broad energy range, the authors may need to revise panel (b) of these three figures. For example on Line 240: Are the  $O^+$  and  $H^+$  densities referred to in this section calculated from one of the energy ranges discussed in Line 202-209? Greater detail and explanation are needed.

**Response:** It may be our inaccurate descriptions result in the referee's misunderstood. Figure 4b 5b, 6b, 7b show the  $O^+/H^+$  density ratio used the broad energy range (as mentioned in Response to comment 6). While Figure 4c, 5c, 6c, 7c show the  $O^+/H^+$  fluxes ratio at several individual energy ranges. We didn't calculate the  $O^+$  or  $H^+$  density from one of the energy ranges discussed in Line 202-209. So this relevant part of the description has been amended in my revised manuscript, see Line 277-285 in "Tracked change" manuscript.

**Comments 8:** Line 254-256: After addressing the above comments on how the ion densities were calculated, the authors should briefly address whether these comparisons of density across missions are relevant. For example, if the  $O^+$  density (calculated over defined HPCA energy range) is higher than seen by Cluster (calculated in what energy range and using which instrument?), what does this mean? Were the instrument energy ranges equivalent or similar? Otherwise, the direct comparison may not be meaningful.

**Response:** Thanks for the referee's good evaluation and kind suggestion. This comment is very important. From Line 254-256, we can't exclude the reason that Bouhram et al., 2005 used somewhat different energy range for O<sup>+</sup> observations result in lower O<sup>+</sup> density in their study than mine. The direct comparison can't be meaningful. In this study, the O<sup>+</sup> density calculated using HPCA distribution functions at energies from 1 to 40 keV, but Bouhram et al., (2005) used CODIF distribution functions at energies from 3 to 40 keV to avoid contamination from high H<sup>+</sup> fluxes. The composition and distribution function (CODIF) analyzer on the Cluster that measures 3-D distributions of the major ion species over the energy range 30–40000 eV. This contrast study is not rigid in this study, so we removed the relevant part in our revised manuscript.

**Comments 9:** Line 305: Since 31 events are not a large number, recommend the authors produce a table to list the dates and times of each of these events so that others in the space science community can also investigate the events for follow-on studies. Such a table could go in an appendix.

**Response:** Yes, this is a good suggestion. I have prepared such a table to list the dates and times of each of these events for follow-on studies in an appendix (see the appendix).

**Comments 10:** All the references in the manuscript need to be checked. For example, all the MMS instrument papers were referenced but do not appear in the references list. It is likely any other references have been missed. It is likely many other references have been missed.

Line 106: Pollock et al. (2016) is referenced but does not appear in the references list Line 105: Russell et al. (2016) is referenced but does not appear in the references list Line 104: Ergun et al. (2016) is referenced but does not appear in the references list Line 104: Lindqvist et al. (2016) is referenced but does not appear in the referenceslist Line 107: Young et al. (2016) is referenced but does not appear in the references list

**Response:** Thanks for the referee's kind suggestion and carefully evaluating this paper. This mistake should have avoided in the manuscript submission. We added the MMS instrument papers citations in the references list. We also checked carefully all the references in the manuscript to make sure all the citations in the references list. The other spelling and syntax errors have also been checked and corrected in the revised paper.

# **Reply to reviewer #2**

#### Dear reviewer:

We are very grateful to your comments for the manuscript and thanks for carefully evaluating this manuscript. According to your advice, we amended the relevant part of the manuscript. The one-to-one responses to your comments are the following.

#### Major comments

**Comments 1**:One of the conclusions of the manuscript is that particles are transported from the tail towards the dayside. To make such a conclusion more rigid one should show the anisotropy of the particle distributions, which would indicated that particles move from the tail towards the dayside. The oxygen ions could also come from other sources such as inner magnetosphere (filled directly from the nightside aurora into the ring current), from the diamagnetic cavities/cusp (e.g. Slapak et al., Ann. Geophys. 2013,10.5194/angeo-31-1005-2013).

**Response:** Thank you for pointing this out. Yes, the oxygen ions at the dayside LLBL have many sources such as the ring current in the inner magnetosphere, the high latitude auroral region and the cusp. Our paper focuses on the oxygen ions in the duskside magnetopause during intense substorms with AE >500nT. Previous research work has reported that the oxygen ions transferred faster into the ring current in the inner magnetosphere and then they are decayed at the dayside magnetopause under southward IMF or with their large gyroradius effect [e.g., Zong et al., 2001]. Under intense geomagnetic activities such as intense substorms and storms, the oxygen ions from the nightside aurora along the plasma sheet or plasma sheet boundary layer can be fast transferred into the near-Earth magnetotail and then injected into the ring current [e.g., Duan et al., 2017 JGR; Yu and Ridley, 2013 JGR]. Recently, Kronberg et al. [2014] reported that the oxygen ions distribution was really anisotropic at the dawn-dusk equator plane. Our observation result is consistent with their report. I have to admit making such a conclusion is not rigid. Because we can't exclude other origins. I removed this expression in my revised paper.

- Kronberg, E. A., Ashour-Abdalla, M., Dandouras, I., Delcourt, D. C., Grigorenko, E. E., Kistler, L. M.,...Zelenyi, L. M. (2014). Circulation of heavy ions and their dynamical effects in the magnetosphere: Recent observations and models. Space Science Reviews, 184(1-4), 173–235. https://doi.org/10.1007/s11214-014-0104-0
- Yu, Y., and A. J. Ridley (2013), Exploring the influence of ionospheric O+ outflow on magnetospheric dynamics: dependence on the source location, J. Geophys. Res. Space Physics, 118, 1711–1722, doi:10.1029/2012JA018411

Zong, Q.-G., B. Wilken, S. Y. Fu, T. A. Fritz, A. Korth, N. Hasebe, D. J. Williams, and Z.-Y. Pu (2001), Ring current oxygen ions escaping into the magnetosheath, J. Geophys. Res., 106(A11), 25,541–25,556.

**Comments 2:I** am not sure if one could make firm conclusions about dependence on the IMF Bz, if from 31 events only 6 events were observed during northward IMF. On my opinion the statistics are too poor for that.

**Response:** Thank you for the comment. The events of energetic oxygen ions at the duskside magnetopause during intense substorms in our studies are chosen from MMS Phase 1. Because there are limited number of events of intense substorms when MMS passes through the duskside magnetopause during the Phase 1. So

the O<sup>+</sup> abundance dependence on IMF Bz is not clear. In the revision processes, We added the 26 event satisfied with the criterion into our work. Our work presents 57 intense substorms events with 50 events under the sourthward IMF Bz and only 7 events under the northward IMF Bz. The intense substorms are usually occurring duing southward IMF Bz .This is consistent with the usually external condition of intense substorms [Lyons et al., 2005]. On the other hand, We have added 26 events satisfied with the criterion in the revised paper for better study. We found the O<sup>+</sup> density shows expontial decrease with IMF Bz from -10 to 0 nT. This conclusion will be substitute with a more rigid expression in my revised manuscript. -->"When the IMF is southward, the O<sup>+</sup> density shows an exponential increase with the IMF Bz absolute value.", as shown in "Tracked change" manuscript Line 31-32.

Lyons, L. R., D.-Y. Lee, C.-P. Wang, and S. B. Mende (2005), Global auroral responses to abrupt solar wind changes: Dynamic pressure, substorm, and null events, J. Geophys. Res., 110, A08208, doi:10.1029/2005JA011089.

**Comments 3:**The "intense substorms" are discussed in this study. Were these substorms associated with magnetic storms? Or these are pure substorm events? What is the reason for choosing intense substorms? Inlcuding other substorms may increase the statistics on the IMF dependence.

**Response:**Thanks for the referee's kind suggestion. In this statistical study, 31 magnetopause crossing events during intense substorm (AE>500 nT) were selected. Among them, there are 4 events during the non-storm time (Dst>-25 nT) and 27 events during the storm time (Dst<-25 nT). There are three resons that we focused on investigating the characteristics of energetic oxygen ions at the duskside magnetopause during intense substorms. Firstly, previous studies have reported that the density and energy flux of oxygen ions in the magnetosphere both increased during magnetic activities, such as intense substorm and storms (e.g., Daglis et al.,1994; Kronberg et al.,2014). During During disturbed times, oxygen ions can be energized due to duskward drift along the dawn- dusk eletric field in the course of their covection from the distant tail towards the Earth. And the oxygen duskward asymmetry is observed at the near-Earth nightside (eg. Nosé et al., 2000; Luo et al., 2014). So I want to realize the relations between magnetail processes with the Oxygen abundance in the duskside magnetpause. Second, Oxygen ions play a significant role in the energy and mass transport in the coupling process of the solar wind-magnetosphere-ionosphere during intense substorms. The responses of energetic oxygen ions at the duskside magnetopause boundary layer to the solar wind conditions during intense substorms have seldom be reported till now. Third, MMS project can provide a good chance to investigate the features of energetic oxygen ions in the dayside magnetopause boundary layer. The previous spacecraft observations provided significant results of oxygen ions mainly focusing on the tail plasma sheet or middle and high latitude region, such as Cluster [e.g., Nilsson et al., 2006; Slapak et al., 2011]. Thus, our investigation can provide new results in the duskside magnetopause.

Daglis, I. A., Livi, S., Sarris, E. T., & Wilken, B. (1994). Energy density of ionospheric and solar wind origin ions in the near-Earth magnetotail during substorms. Journal of Geophysical Research, 99(A4), 5691–5703. https://doi.org/10.1029/93JA02772

- Kronberg, E. A., Ashour-Abdalla, M., Dandouras, I., Delcourt, D. C., Grigorenko, E. E., Kistler, L. M.,...Zelenyi, L. M. (2014). Circulation of heavy ions and their dynamical effects in the magnetosphere: Recent observations and models. Space Science Reviews, 184(1-4), 173–235, doi:10.1007/s11214-014-0104-0.
- Ono, Y., M. Nosé, S. P. Christon, and A. T. Y. Lui (2009), The role of magnetic field fluctuations in nonadiabatic acceleration of ions during dipolarization, J. Geophys. Res., 114, A05209, doi:10.1029/2008JA013918.
- Luo, H., E. A. Kronberg, E. E. Grigorenko, M. Franz, P.W. Daly, G. X. Chen, A. M. Du, L. M. Kistler, and Y.Wei (2014), Evidence of strong energetic ion acceleration in the near-Earth magnetotail, Geophys. Res. Lett., 41, 3724–3730, doi:10.1002/2014GL060252.
- Nilsson, H., et al. (2006), Characteristics of high altitude oxygen ion energization and outflow as observed by Cluster: A statistical study, Ann. Geophys., 24, 1099–1112.
- Slapak, R., Nilsson, H., Waara, M., André, M., Stenberg, G., and Barghouthi, I. A.(2011), O+ heating associated with strong wave activity in the high altitude cusp and mantle, Ann. Geophys., 29, 931–944, doi:10.5194/angeo-29-931-2011

**Comments 4:** Introduction, first two paragraphs can be merged as they contain repeating information about acceleration during dipolarizations. The second paragraph is not completely logical. It would make more sense to describe acceleration of  $O^+$  starting from the polar region, then lobe, dipolarizations and then discuss drift. The sentence in lines 43-45 discussing acceleration of electrons during dipolarizations is not really needed as there is a number of references about acceleration of oxygen during dipolarizations in lines 29-47 and the whole text is about  $O^+$ .

**Response:** Thanks for the referee's kind advice. As you suggested, I adjusted the description order in seconde paragraph to make the introduction more logical and concise. The part of revision can be found in Line 40-65 in "Tracked change" manuscript.

**Comments 5:** lines 90-91, "At present,  $O^+$  near the dayside low-latitude magnetopause during substorm expansion phase and recovery phase are still not understood"  $\rightarrow$  What exactly do you mean under not understood? Which scientific questions are still open? Which questions do you try to answer?

**Response:** Thank you for these comments. Actually, what we want to know is how the O<sup>+</sup> abundance (O<sup>+</sup>/H<sup>+</sup>) in the duskside magnetopause varies on AE index and solar wind conditions (e.g. IMF By, IMF Bz, and solar wind dynamic pressure) during the intense substorm (AE >500 nT). The relevant description is revised in "Tracked change" manuscript Line 114-117.

**Comments 6:**lines 91-93, there is paper by Luo et al., JGR, 2017, 10.1002/2016JA023471, in which the energization of O<sup>+</sup> at the dayside is discussed. The study also discusses asymmetries of the energetic oxygen due to IMF By and Bz directions. Both IMF By and Bz influence the oxygen abundance at higher energies. However, this is large statistical study and not only cases for the intense substorms. This can be discussed.

**Response:** That would be great. We discussed Luo et al., (2017) results in my revised manuscript. The relevant part can be found in "Tracked change" manuscript Line 399-402. Recently, Using energetic ion

composition data at the low latitude dayside magnetopause measured by Magnetospheric Multiscale (MMS) satellites, we study the response of O<sup>+</sup>abundance (O<sup>+</sup>/H<sup>+</sup>) to the both IMF By and Bz and not only cases for the intense substorms. We found that they indeed influence the oxygen abundance even at lower energies (1-40 keV) and more significant duskside asymmetry of O<sup>+</sup> under southward IMF with positive IMF By. These results are consistent with those of Luo et al.,(2017).

**Comments 7:** lines 125-126, 130-131, 180-181, please provide a more precise definition of the substorm onset and recovery phase. For example in paper by Newell and Gjerloev, JGR, 2011, 10.1029/2011JA016779, is a nice example on how to define substorm onset, also using more precise SML index available at the SuperMAG. I do not think that definition when "AE index significantly increases" is a precise one. I do not think that one should provide twice the information about substorm onset in lines 125-126 and 130-131. I would remove the second sentence.

**Response:** Thanks for constructive comments and nice recommendation. We have added a more precise definition of the substorm onset, expansion phase and recovery phase in our revised manuscript. The second information about substorm phase description in lines 125-126 and 130-131 has been removed. We added AU, AL index in Figure 1 to help us identify the phase of a substorm. First, we determined the time interval of the magnetopause boundary layer crossings in each event. Then, we find out how the substorm indices change during that interval from the OMNI data. As Figure 1 shown, the time interval of the magnetopause boundary layer crossing is indicated by the two blue dashed lines. As we know, the AE index is defined as AE=AU-AL. Generally, the substorm onset time is characteristic by the AL index starts to significantly decrease and the AE index significantly increase. During the substorm expansion phase, the AL index will decrease significantly. The interval of the AL index decrease from onset to its minimum is defined as the substorm expansion phase. Then it starts to increase and the interval of the AL index increase from the minimum to the quiet time level is regarded as the substorm recovery phase. In our event, the MMS4 crossed the magnetopause boundary layer from 15:25:10 to 15:36:50 UT on 3 October 2015. From Figure 1f, the AL index reached its minimum ~-750 nT and AE index reach the peak ~1000 nT at about 15:20 UT, then it started to increase to  $\sim$  -200 nT at the rest time of interest. So the magnetopause boundary layer crossing occurred during the intense substorm recovery phase. (see Line 167-180 in "Tracked change" manuscript)



Figure 1. The three components IMF Bx, By, Bz, solar wind dynamic pressure, as well as AU, AL, and AE index from CDAweb OMNI data.

**Comments 8:**lines 179-180, actual observations of the IMF and solar wind dynamic pressure could be used directly from the MMS observations at the magnetopause crossings. This would be much more precise.

**Response:** Getting the much more precise IMF and solar wind dynamic pressure would be better. When the IMF passes through the bow shock, its direction would be changed in the magnetosheath. We will compare

the IMF and solar wind dynamic pressure directly from the MMS-with those from OMNI data in the detailed events analysis. Then we will choose the more precise data.

**Comments 9:** lines 277-278, For higher energies the larger statistics one can clearly see that the stronger duskward asymmetry in the plasma sheet and the dayside magnetosphere is observed under the southward IMF, e.g. Luo et al., JGR, 2017. One should mention that no influence of IMF Bz is observed in case of the energies below 40 keV and for 31 intense substorm events.

**Response:** Thanks for your constructive suggestion.We agree with your comments. Recently, We used energetic ion composition data at the dayside magnetopause measured by Magnetospheric Multiscale (MMS) satellites, we study the response of O<sup>+</sup> abundance to IMF Bz. The O<sup>+</sup> abundance showing strong duskside asymmetry in the magnetopause boundary layer under southward IMF than that under northward IMF also be found in our study, which is consistent with Luo et al.2017 result. As you suggested, the influence of IMF Bz is not clear in our 31 intense substorm events. This description in Lines 277-278 has been removed.

**Comments 10:** lines 286-287, the energetic  $O^+$  occurs predominantly under southward IMF. Here I would say that it was chosen to be like this. Choosing the intense substorms one increases the probability of observing the southward IMF quite significantly. This also contradicts to statement in the lines 277-278, that IMF Bz does not influence abundance of  $O^+$  at the magnetopause. There is not enough provided data to conclude so.By increasing the number of events under the northward IMF one may see a different picture. One can see pretty nice trend in Figure 6b, that the abundance is increasing with the decrease of IMF Bz at least for the expansion phase. Generally on my opinion there is not enough statistics in this study to make conclusions about IMF dependence. One should expand the statistics

**Response:** Thanks for your valuable comments. We agree with you that choosing the intense substorms one increases the probability of observing the southward IMF significantly. To make convincing conclusions about IMF dependence, we expanded the statistic to 57 events. Some conclusions may be still not convincing due to not enough statistical events. As the MMS operate longer, more magnetopause crossing during intense substorm will be detected. It will be helpful. The sentences in Line 286-287 have been deleted.

**Comments 11:** lines 304-306, this conclusion is not supported by the observations. Just looking at the scatter points of the number density, I do not see a statistically significant difference between these two phases. One should either show fits to those points or bin them according to some parameters and show that the difference is significant.

**Response:** Thanks for your comments. At the beginning of this study, we focus on the response of  $O^+$  abundance on the geomagnetic activity and solar wind conditions during intense substorms. Because the magnetosphere has the different dynamics in the near-Earth space during the different phase of intense

substorms, especially in substorm expansion phase and recovery phase. So we want to investigat variations of energetic O<sup>+</sup> density at the duskside magnetopause boundary layer duing different phases of intense substorms. Due to the number of events are limited (only 26 events during expansion phase), we don't think it makes sense to fit those points or bin then according to some parameters. As the MMS operating longer, more magnetopause crossing during intense substorm will be detected. It will be helpful. Our selecting events we drawn our summary on the energetic O<sup>+</sup> density as description in the last part of our manuscript. In generally, the O<sup>+</sup> in the magnetophere are origin from the ionosphere and transferred into the different magnetophere region during magnetic activities. A excellent review paper of this aspect has been reported by Keika et al., (2013). Our new results from MMS data provide another support of previous studies.

Keika, K., L. M. Kistler, and P. C. Brandt (2013), Energization of O+ ions in the Earth's inner magnetosphere and the effects on ring current buildup: A review of previous observations and possible mechanisms, J. Geophys. Res. SpacePhysics, 118, 4441–4464, doi:10.1002/jgra.50371.

**Comments 12:** lines 313-315, energetic oxygen ions also indicate the transport at the dayside magnetosphere (e.g. Liao et. al, JGR, 2010, 10.1029/2010JA015613). These different transports are hard to distinguish (e.g. Luo et al., JGR, 2017).

**Response:** Thanks for your comments and paper recommendation. Liao et al., (2010) JGR and Luo et al., (2017) JGR are both cited in our revised manuscript (see see Line 391-397 and Line 399-402 in "Tracked change" manuscript). I agree with you, these different transports are hard to distinguish. Also, the different transports of oxygen ions from the ionosphere to different part of the magnetosphere are significant and interest. It is outside the focus of our manuscript. We will investigate this issue with conjunction observations by multiple spacecraft in different magnetosphere locations. This conclusion in Line 313-315 has been deleted.

**Comments 13:**Figures 4-7, just looking at the scatter plots it is hard to make certain conclusions. One should either bin the points to show the average trend or fit them with some dependences and increase the number of events.

**Response:** Thanks for your nice suggestions. I have binned the points to show the average trend before submitted this manuscript. As you said, the number of events is too low, so the trend is not obviously or has low credibility and we abandoned this method. As you suggested, we fit the oxygen density dependence on the IMF Bz and Psw, respectively. Recently, Using energetic ion composition data at the low latitude dayside magnetopause measured by Magnetospheric Multiscale (MMS) satellites, we study the response of H<sup>+</sup>, O<sup>+</sup>density and their ratio to the geomagnetic activity (indicated by SYM-H index) and solar wind conditions (including interplanetary magnetic field (IMF) By, IMF Bz and solar wind dynamic pressure). In this study, we bin the points due to enough events. Our new manuscript has been submitted to the JGR.

#### Minor comments:

1. **Line 19:** What is the energy range of the oxygen observations used in this study? Please indicate the upper energy limit in the abstract. This is important to know when assessing the number densities.

**Response:** In this study, only the O<sup>+</sup> at energies from 1 keV to 40 keV measured by HPCA are used. The upper energy limit of HPCA is 40 keV. This information is added to the abstract. (see Line 18 in "Tracked change" manuscript )

2. Line 45: I did not find the reference to Lui et al., 1999 in the reference list. (see Line 589-591 in "Tracked change" manuscript )

Line 47: "during activity geomagnetic disturbance" -> "during disturbed geomagnetic activity" (see Line 50-51 in "Tracked change" manuscript )

4.Line 55: "[e.g. Yau and Andre, 1997]. And then..." -> "[e.g. Yau and Andre, 1997] Then..." (see Line 62 in "Tracked change" manuscript )

5. Line 85: please remove one "However". (see Line 109 in "Tracked change" manuscript )

**Response to comments from 2-5:** The above expression errors have been checked and corrected. The missing reference has been added to the revised manuscript.

6. Lines 106-107: Does HPCA distinguish between O<sup>+</sup>, N<sup>+,</sup> and C<sup>+</sup>? Or what measures actually the CNO group?

**Response:** "the HPCA is a time-of-flight (TOF) mass spectrometer designed to measure the velocity distributions of the four ion species ( $H^+$ ,  $He^+$ ,  $He^+$  and  $O^+$ ) known to be important in the reconnection process. The measurement technique is based on a combination of electrostatic energy-angle analysis with time-of-flight velocity analysis. The result is an accurate determination of the velocity distributions of the individual ion species. In order to meet the stringent scientific requirements of the MMS mission, the HPCA incorporates three new technologies. The first extends counting rate dynamic range by employing a novel radio frequency mass filter that allows minor species such as  $He^{++}$  and  $O^+$  to be measured accurately in the presence of intense proton fluxes found in the dayside magnetopause. The second ensures that TOF processing rates are high enough to overlap with the low end of the RF dynamic range, while the third enhances ion mass resolution.

Fig. 20 TOF spectrum for four ion species and background  $(H_2^+)$  is a substitute for He<sup>++</sup> and N<sup>+</sup> is a substitute for O<sup>+</sup>). *Red areas* demarcate bins that define ion species and background. The *peak* at ~200 ns corresponds to N\_2<sup>+</sup>





During each energy scan a data set consisting of 63 TOF spectra  $\times$  512 TOF bins  $\times$  16 elevations is accumulated and histogrammed. The resulting TOF spectra are then parsed into five bins that define the ion species H+, He++, He+, O+ and background (Fig. 20). The red portions of the spectrum in Fig. 20 indicate typical species boundaries. Since ion times-of flight are both mass and energy-dependent the range of TOF limits for each species changes with energy (Fig. 31)". (the Figure and description are cited from Young, D. T., Burch, J. L., Gomez, R. G., De Los Santos, A., Miller, G. P., Wilson, P., et al. (2016). Hot Plasma Composition Analyzer for the Magnetospheric Multiscale Mission. Space Science Reviews, 199(1–4), 407–470, doi:10.1007/s11214-014-0119-6.).

As for this interesting question, I specially contacted the HPAC PI (Stephen Fuselier), he replied me "I'm working right now to see if we can see C<sup>+</sup> and possibly N<sup>+</sup> in the mass spectra. They would not appear as a separate mass peak because of straggling in the foil. I'm not sure if we can even tell if they are there. What we bring to the ground and call O<sup>+</sup> could contain substantial N<sup>+</sup>. The C<sup>+</sup> peak would probably be at a lower time-of-flight than what we bring to the ground, but you could safely say that what we call O<sup>+</sup> could be N<sup>+</sup>O<sup>+</sup>."

7. Line 122: "At the beginning of the time interval, the solar wind dynamic pressure..."->The dynamic pressure is only at the begging of the time interval about 2 nPa. (see Line 161-162 in "Tracked change" manuscript )

8. Lines 124-125: I would change to " These solar wind conditions led to the intense substorm (AE>500 nT). (see Line 166 in "Tracked change" manuscript )

9. Lines 148-150: Figure 2, I would say that the fluxes at energies below 2 keV in Figure 2j is also contamination. This should be mentioned also in Figure caption and even better when it is indicated on the plot itself. (see Figure 2j red box in revised manuscript )

**Response to comments from 7-9:** Thanks for the referee's valuable suggestion. The minor comments 7 and 8 have been corrected in my manuscript. The red box indicating the O<sup>+</sup> contamination from high proton fluxes was plotted in figure 2j and relevant description was mentioned in the Figure 2 caption (see Line 676-677).

10. Line 195:I would remove "On the other hand".

Response: This is a common colloquial expression. We removed it.

11. **Lines 221-224:** These results also agree with Kronberg et al., JGR, 2012,10.1029/2012JA018071 which showed for 10 keV O<sup>+</sup> strong increasing under the duskward IMF indicated by the clock angle in the inner magnetosphere.

**Response:** That would be great. We cited this paper in the relevant part of the revised manuscript to increase valid of our results. see Line 387-389 in "Tracked change" manuscript )

12. Lines 251-268: another reason can be that *Bouhram et al.*, 2005 have used somewhat different energy ranges for  $O^+$  observations.

**Response:** Yes, We agree with you. I can't exclude the reason that *Bouhram et al.*, 2005 used somewhat different energy range for  $O^+$  observations. In this study, the  $O^+$  density calculated using HPCA distribution functions at energies from 1 to 40 keV, but Bouhram et al., (2005) used CODIF distribution functions at energies from 3 to 40 keV to contamination from high H<sup>+</sup> fluxes. This contrast study is not rigid in this study. We removed the relevant part in my revisited manuscript.

13. Line 276:magenetopause -> magnetopause.(see Line 384 in "Tracked change" manuscript )

14. Line 279: have -> has. (see Line 390 in "Tracked change" manuscript )

15.Line 287: dominated occurring -> occurs predominantly.(removed)

17. **Lines 296-297:** I would change this sentence to "The reconnection rate is likely will be reduced by the mass-loading but not suppressed at the magnetopause [*Fuselier et al.*, 2019]. (see Line 429-430 in "Tracked change" manuscript )

17. **Figure 1**, caption, "The three components of the IMF, Bx, By, Bz..." (see Line 661 in "Tracked change" manuscript )

18. **Figure 2**, I would indicate on the plot contamination. In the caption, line 481 (k)–>(l). (see Line 671 in "Tracked change" manuscript )

**Response to comments from 13-18**: Thanks for referee'scarefully evaluating this paper and important suggestions. We have revised the above errors and plotted the red box indicating the  $O^+$  contamination from high proton fluxes in figure 2j. The other spelling and syntax errors have been checked and corrected. We acknowledge the reviewer's comments and suggestions very much, which are valuable in improving the quality of our manuscript.

#### **Reply to reviewer #3**

# Dear reviewer:

We are very grateful to your comments for the manuscript and thanks for carefully evaluating our manuscript. We acknowledge the reviewer's comments and suggestions very much, which are valuable in improving the quality of our manuscript. According to your advice, we amended the relevant part of the manuscript. Responses to your comments are below point by point.

**Comments 1**: Lines 90-95: There is a lot of information leading up to this point in the introduction, however with the lines preceding and in this paragraph itself, it is unclear what is not well understood and how/what this paper will provide to answers to. Currently, the introduction reads as a quite thorough list of previous studies, but it is not readily apparent how they string together, and what they are necessarily building up to. I would suggest stating what the paper will study before this point and tailoring the introduction to build off of that somewhat, because at this point as a reader it is still unclear.

Response: Thanks for the referee's kind advice. As you suggested, we did some revisions in our revised manuscript. We adjusted the first two paragraphs to make the introduction more logical and concise (see Line 40-65 in "Tracked change" manuscript). The introduction is organized as following orders. First, we stress the importance of  $O^+$  during the intense substorms, describe acceleration of  $O^+$  starting from the polar region, then lobe, near-Earth plasma sheet and then discuss drift. Second, we describe the O<sup>+</sup> behavior in the magnetopause. Third, we referred to the  $O^+$  density dawn-dusk asymmetry in the magnetopause. Finally, we describe the questions what this paper tries to answer. As the following described: "At present, variations of O<sup>+</sup> abundance (O<sup>+</sup>/H<sup>+</sup>) in the dusk flank magnetopause during intense substorms (AE >500 nT) on AE index and solar wind conditions (e.g. IMF By, IMF Bz, and solar wind dynamic pressure) are still not understood. Previous studies of O<sup>+</sup> during intense substorms mainly focused on O<sup>+</sup> energizations in the NEPS in the magnetotail (e.g., Duan et al., 2017; Nosé et al., 2000; Ohtani et al., 2011). At present, The Magnetospheric Multiscale (MMS) mission gives us an opportunity to focus on the O<sup>+</sup> in the low latitude dayside magnetopause region. In this study, we mainly investigate statistical features of energetic O<sup>+</sup> in the dusk flank magnetopause and their relations with AE index and solar wind conditions (e.g. IMF By, IMF Bz and solar wind dynamic pressure) during the intense substorms. (see Line 114-124 in "Tracked change" manuscript)

**Comments 2:** Lines 128-150: HPCA & FPI fluxes are in differential flux and energy flux units. Is there a benefit in having their fluxes in different units? If they are to remain, a point should be included in the text that the units are different.

**Response:** Thanks for the referee's kind suggestion. We described the HPCA and FPI fluxes having different units in our revised manuscript. Figures 2g and 2h show the electron omnidirectional differential energy fluxes and ion omnidirectional differential energy fluxes, respectively. Figure 2i to 2l presents the differential particle fluxes of H<sup>+</sup>, O<sup>+</sup>, He<sup>+</sup>, He<sup>++</sup>, respectively. To better identify the fluxes variations at specific energies, we choose the ion and electron fluxes from FPI in the energy flux unit. The relevant description has been add into the Line 195-198 in "Tracked change" manuscript)



Figure 2. The energetic  $O^+$  is observed at the magnetopause during an intense substorm on 03 October 2015 by MMS 4. From top to bottom are (a) the magnetic field three components, Bx (blue line), By (gree line), Bz (red line) and the total magnitude Bt (black line), (b) the electric field three components, Ex (blue), Ey (gree) and Ez (red), (c) Ion parallel (red) and perpendicular(black) temperatures, (d) The density of ion (green) and electron (blue), (e) three components of the ion velocity, (f) the H<sup>+</sup>(overthe full HPCA energy range) and O<sup>+</sup> (at energies from 1 to 40 keV)densities, (g) electron omnidirectional differential energy fluxes, (h) ion omnidirectional differential energy fluxes, (i) to (1) present differential particle fluxes of H<sup>+</sup>, O<sup>+</sup>, He<sup>+</sup>, He<sup>++</sup>, respectively. The Geocentric Solar Magnetospheric (GSM) coordinate system is adopted. The thick bars at the top of the panel present different regions encountered

on this magnetopause crossing event. The orange and blue bars represent the magnetosheath and the magnetosphere, respectively. The green bar represents the magnetopause boundary layer. The black horizontal line in figure 2j is at 1keV and the  $O^+$  contamination from high H<sup>+</sup> fluxes is indicated by the red box. The FPI data in Figure 2(c-e) and (g-h) are from FPI L2 data products and in the fast mode.

**Comments 3:** Lines 128-150: The HPCA flux in panels i-l have artificial striping every 4 energy bins due to way HPCA determines the count rate over 4 energy channels in survey mode. It would be best to correct this, however, describing the artificial striping would also be sufficient. I am also not certain that these HPCA fluxes are actually omni-directional as they do not appear to be half-spin averaged, please verify.

**Response:** Thanks for your important comments. "The HPCA flux in panels i-l have artificial striping every 4 energy bins due to way HPCA determines the count rate over 4 energy channels in survey mode." The above sentence has been added into the Line 194-195 in "Tracked change" manuscript. Figure 2i to 2l presents the differential particle fluxes of H<sup>+</sup>, O<sup>+</sup>, He<sup>+</sup>, He<sup>++</sup>, respectively. They are actually not Omnidirectional and not half-spin averaged. We corrected this description in our revised manuscript. These differential particle fluxes of H<sup>+</sup>, O<sup>+</sup>, He<sup>++</sup> calculations are used The Space Physics Environment Data Analysis System (SPEDAS) software package. More details about SPEDAS can be found in Angelopoulos et al. (2019) and cited as (Angelopoulos, V., Cruce, P., Drozdov, A. et al. Space Sci Rev (2019) 215: 9. https://doi.org/10.1007/s11214-018-0576-4). We also cited this paper in our revised manuscript (see Line 478-479 in "Tracked change" manuscript).

**Comments 4:** Lines 134-137: Please describe where the FPI/HPCA moments shown come from. This is quite important since the majority of the results presented are dependent on these moments.

**Response:** Thanks for the referee's kind advice. We have clarified where the FPI/HPCA moments shown come from. We have added detailed information about moments into Line 198-207 in "Tracked change" manuscript. The plasma moments (e.g. Ion parallel and perpendicular temperatures, ion, and electron densities and ion velocity) from FPI shown in Figure 2c-e are all from MMS L2 data products. They are default moments calculated over the full FPI energy range from 10 eV to 30 keV. But the O<sup>+</sup> density shown in Figure 2f is recalculated from HPCA distribution functions in the range of energies from 1 to 40 keV. From the O<sup>+</sup> fluxes shown in Figure 2j, there still exist a large number of fluxes below 1 keV in the magnetosheath. This part of O<sup>+</sup> fluxes is fake and contamination from high proton fluxes. So we consider the number density of O<sup>+</sup> with energies from 1 to 40 keV. It is more appropriate to represent the true O<sup>+</sup> in the magnetopause. While the H<sup>+</sup> density which computed over the full HPCA energy range from 1eV to 40 keV from L2 data products are used in Figure 2f.

**Comments 5:** Figures 1-2: I would suggest using these two figures to establish the criteria for the statistical study. In my opinion, more text should be added that describes a greater context for these 2

figures inclusion. Either establishing points that lend themselves to the paper's conclusion and/or use the figure to establish conditions for the statistical study.

**Response:** Thanks for your nice comments. In this statistical study, First, we identified the magnetopause crossing event (complete magnetopause crossing from the magnetosheath to the magnetosphere, vice versa) during phase 1 from the summary plot in https://lasp.colorado.edu/mms/sdc/public/plots/. Then we plotted the more detailed overview of these events to identify the magnetopause boundary layers, as Figure 2 shown. Figure 2 was mainly used to determine the magnetopause boundary layer crossing interval. Only events that AE index larger than 500 nT during the magnetopause boundary layer crossings interval were selected. Finally, the mean value of the H<sup>+</sup>, O<sup>+</sup> density and their fluxes shown in Figure 2 were calculated in the magnetopause boundary layer. Correspondingly, the AE index, IMF By, Bz and solar wind dynamic pressure from the OMNI data system shown in Figure 1 were also averaged during that interval. Figure 1 mainly provided the corresponding solar wind conditions and AE index. The above expressions have been added into Line 245-256 in the "Tracked change" manuscript.

**Comments 6**: Lines 176-181: This is one of the more major comments on the paper. The current description of the event selection criteria is not sufficient. Interpretation of a statistical study is almost entirely dependent on understanding how the statistical study is conducted. It is currently not clear what the criteria for event selection is. Is it any MP crossing with AE > 500? Why was 500 chosen as a threshold in AE (i.e. stats are somewhat low, would AE > 300 or 400 provide more events and still be "intense"?).

**Response:** Thanks for your valuable comment. The magnetopause crossing event in our statistical study all during the intense substorm (AE > 500nT). How we chose these events is replied in the before comment. The reason why we choose intense substorm with AE >500 nT is based on the results from Daglis et al (1994) (Figure 6 in this reference, as shown in the below). They found that the O<sup>+</sup> energy density has a great correlation with the AE index in the near-Earth plasma sheet (NEPS). During the intense substorm expansion phase, O<sup>+</sup> energy density explosively increases with AE index in the range of larger than 500nT. Otherwise, Lennartsson and Shelley, (1986) pointed out that the ion composition had a large variance at substorm. During the intense disturbed conditions (AE~1000nT), the increase in the O+ energy density is strongest around local midnight where O<sup>+</sup> become the most abundant ion. The previous researches of oxygen ions during intense substorms are mainly focused on the nightside NEPS. Thus, we want to know whether the O<sup>+</sup> abundance in the dusk flank magnetopause varies on AE index and solar wind conditions during the intense substorm and how it changes with the above parameters. Characteristics of Oxygen ions in the high latitude polar region and near-Earth magnetosphere during

intense magnetic activities have been investigated deeply and widely. But O<sup>+</sup> abundance in the dayside magnetopause has seldom been reported during intense substorms.



Figure 6. Same format as Figure 2; all substorms with  $AE_{\text{exp}} \ge 500 \text{ nT}$  (encircled symbols) are excluded from the fit.

**Comments 7:** How exactly is the magnetopause boundary layer determined? Is there any consideration for if the substorm is during a storm or the 1st/2nd/3rd in a series of substorms? Specifically, how are substorm phases determined? What is meant by the mean value of the flux (over a range of energies, one energy)? How long were the average events? Please provide greater context for the choices of criteria used in this study.

Response: The magnetopause boundary layers are identified here primarily through plasma fluxes and moments. The magnetopause boundary layer can have densities and temperatures between that of the magnetosphere and the magnetosheath. Meanwhile, the magnetopause boundary layer shows the gradient of the energy flux of particles and number density and magnetic field obvious. Ion jets are also signatures of passing through the magnetopause boundary layers. In this study, the separatrix between the magnetosheath and the magnetopause boundary layer is determined by the appearance of the magnetospheric electron. Similarily, the separatrix of the magnetosphere and the magnetopause boundary layer is determined by the magnetosheath electron disappearance. The mean value of the H<sup>+</sup>, O<sup>+</sup> density and their fluxes are calculated in the magnetopause boundary layer. Correspondingly, the AE index, IMF By, Bz and solar wind dynamic pressure from the OMNI data system were averaged during the time interval of magnetopause boundary layer crossing. As Figure 2 shown, the time interval of the magnetopause boundary layer crossing is marked by the two blue dashed lines. As we know, the AE index is defined as AE=AU-AL. Generally, the substorm onset time is characteristic by the AL index starts to significantly decrease and the AE index significantly increase. During the substorm expansion phase, the AL index will decrease significantly. The interval of the AL index decrease from onset to its minimum is defined as the substorm expansion phase. Then it starts to increase and the interval of the AL index increase from the minimum to the quiet time level is regarded as the substorm recovery phase. In our event, the MMS4 crossed the magnetopause boundary layer from 15:25:10 to 15:36:50 UT on 3 October 2015. From Figure 1f, the AL index reached its minimum ~-750 nT and AE index reach the peak ~1000 nT at about 15:20 UT, then it started to increase to ~ -200 nT at the rest time of interest. The two blue dashed lines indicate the time interval of the magnetopause boundary layer crossing. According to the variation and peak value of the AU, AL and AE index in Figure 1e to 1g.The magnetopause boundary layer crossing occurred during the recovery phase of this intense substorm. The mean value of the flux is over two energy ranges close to the typical energy such as 1 keV, 10 keV and etc. We didn't consider if the substorm in during a storm or the 1st/2nd/3rd in a series of substorms. In this statistical study, 31 magnetopause crossing events during intense substorm (AE>500 nT) were selected. Among them, there are 4 events during the non-storm time (Dst> -25 nT) and 27 events during the storm time (Dst< -25 nT).These detailed contexts for choices of criteria used in this study are described Line 246-250 in the "Tracked change" manuscript.



Figure 1. The three components of the IMF Bx, By, Bz, solar wind dynamic pressure, as well as AU, AL, and AE index from CDAweb OMNI data.

**Comments 8:** Lines 179-180: One of the main points from this paper is that the high-density O<sup>+</sup> can be transported from the nightside tail to the magnetopause where it is observed. Please discuss any effect

(or lack there of) of using OMNI solar wind values at the bow shock to correlate with observations of high  $O^+$  density which is being driven by processes which invariably take some amount of time to occur.

**Response:** Thanks for the referee's good evaluation and kind suggestion. Making such a conclusion is not rigid. I didn't give direct evidence to prove that these O<sup>+</sup> are transported from the tail towards the dayside. So, I deleted this expression in my revised paper.

**Comments 9:** Lines 203-205: With the decimation of HPCA fluxes during survey mode, the count rate is recorded/distributed over 3-4 energy channels. With this in mind, is it appropriate to describe the comparisons of the flux as being over such a small energy range, since the flux/count rate could have been dominated by a nearby energy channel? Potentially, it would be more accurate to re-bin the HPCA flux into 16 energy channels instead of 63, and compare the >1 keV flux levels of these larger energy bins. Please discuss, currently it seems a bit misleading to describe the flux as being over such a narrow energy range.

**Response:** Thanks for the referee's nice comment and kind suggestion. The main purpose of calculating the  $O^+/H^+$  particle fluxes ratio is to study the  $O^+$  abundance at different energies on AE index and solar wind conditions (e.g. IMF By, IMF Bz and solar wind dynamic pressure) during the intense substorms. Since the energy range of  $O^+$  and  $H^+$  in the HPCA are the same. So we directly divide  $O^+$  particle fluxes by H+ particle fluxes and we mainly focus on the ratios and not the values of their fluxes at specific energies.

Comments 10: Lines 231-236: Here it is stated that, "the maximum number density of energetic O+ at the dusk flank magnetopause is during the intense substorms recovery phase under the southward IMF. But the maximum ratio of  $n(O^+)/n(H^+)$  at the dusk flank magnetopause is during intense substorm recovery phase under the northward IMF. IMF Bz seems play a minor role in O+ abundance at the dusk flank magnetopause during intense substorm." It is not clear from the data as it is presented that this is true. The density ratio is of course dependent on O+ and H+ (which can come from the ionosphere and the solar wind). Comparing Figures 4a and 5a, it is not clear to me by eye that n(O+) is more dependent on By than Bz. It very well may be, but it is not readily apparent. Thus, is the density ratio difference actually from O+ or H+? Additionally, only 6 of the events in the study have a Bz > 0. This is notable, as Bz not being random does have an impact on the events. Thus, from this study it appears that Bz does play a role in the events being studied.

**Response:** Thanks for your valuable comments. Some descriptions determined by eye are not convincing. So I add the detailed information about density and corresponding IMF conditions in the supplement materials. The sentences in Line 231-136 have been revised by "The maximum density of energetic O<sup>+</sup> at the dusk flank magnetopause is under the southward IMF. Meanwhile, the maximum O<sup>+</sup>/H<sup>+</sup> density ratio at the dusk flank magnetopause is under the southward IMF." The conclusion of "IMF Bz seems to play a minor role in O<sup>+</sup> abundance at the dusk flank magnetopause during intense substorm." In this manuscript is not rigid. So the relevant description has been removed. It noted that choosing the intense substorms one increase the probability of observing the southward IMF significantly. We found a nice trend that  $O^+$  abundance increase with the IMF By. From Figure 6b, the  $O^+/H^+$  density ratio show an obvious decrease with IMF Bz from -2 to 2 nT during the recovery phase (red crosses shown). Due to not enough statistical events (only 6 of the events in the study with northward IMF), some conclusions may be not convincing. As the MMS operating longer, more magnetopause crossing during intense substorm will be detected. It will be helpful. The relevant part "the IMF Bz seems play a minor role in  $O^+$  abundance at the dusk flank magnetopause during intense substorm." has been deleted.

**Comments 11:** Lines 241-242: "number density ratio at the dusk flank magnetopause during intense substorms have a weak correlation with the solar wind dynamic pressure." Can you quantify this correlation? In general, there are a lot of points currently that are driven from visual inspection of very scattered plots, when greater statistical rigor perhaps could be applied.

**Response:** Thanks for your suggestions. I agree with you, we need quantify this correlation and by eye is not rigid. So we fit the oxygen density dependence on pressure. Due to the number of events are limited (only 9 events during expansion phase) and distribution plot is very scattered, we don't think it makes sense to bin them according to some parameters. The sentence "number density ratio at the dusk flank magnetopause during intense substorms have a weak correlation with the solar wind dynamic pressure." has been removed. And substituted by more detailed description  $\rightarrow$  "Figure 7a present that the O<sup>+</sup> density at the dusk flank magnetopause during intense substorms has a positive correlation with the solar wind dynamic pressure. The empirical functional relation between the O<sup>+</sup> density and solar wind dynamic pressure (from 1 to 4.5 nPa) is also established in the Eq.(3) and the corresponding correlation coefficient is 94%. From Figure 7b, the O<sup>+</sup>/H<sup>+</sup> density ratio during recovery phase show a decrease from about 2.5 to 3 nPa. It is also noted that the O<sup>+</sup>/H<sup>+</sup> density ratio increase with solar wind dynamic pressure from ~3 to 4 nPa."

**Comments 12:** Figures 4-7: The captions of the figures mention that the 95% confidence intervals are shown. Please mention this in the text and describe how it is calculated.

**Response:** Thanks for your kind suggestion. In the revised manuscript, we add the error bars in each point indicating a 90% Confidence Interval (CI). How to calculate the CI is described as follows: Step 1: find the number of observations n in the magnetopause boundary layer. Then calculate their mean  $\overline{x}$  and standard deviation s. Step 2: Find the k value for 90% CI (the k value is 1.65). Step 3: use that k in this formula for the CI:

$$\overline{x} - k \frac{s}{\sqrt{n}} < \mu < \overline{x} + k \frac{s}{\sqrt{n}}$$

Where  $\bar{x}$ , *s* and *n* are the mean value, standard deviation and the sampling number of observations, respectively. *k* in the above formula can be determined by calculating a 90% CI for each events. See Line 286-292 in the "Tracked change" manuscript.

#### Very minor comments:

1. Lines 103-106: Please explicitly state that FPI does not discriminate between different ion species. **Response:** thanks for your kind suggestion. We added the "FPI does not discriminate between different ion species" in the Line 136-137 in the "Tracked change" manuscript.

2. Line 107: Strictly speaking, HPCA measures up to 40 keV/q (thus for He<sup>++</sup> this gets up towards 80 keV).

**Response:** Thanks for you carefully evaluate this manuscript. We agree with you, the HPCA maximum measurement for energy per charge is 40 keV/q. Line 138 in the "Tracked change" manuscript has been revised as you suggested.

**3.** Line 116: The authors might as well finish this thought, that this is due to spacecraft separation/scales of particle motion.

**Response:** Thanks for your nice suggestion. We added this sentence "this is due to spacecraft separation/scales of particle motion." into Line 150-151 in the "Tracked change" manuscript for finishing this thought.

# 4. Line 296: Fuselise et al. should be Fuselier.

**Response:** Thanks for you carefully evaluating this manuscript and giving important suggestions. We have revised this error. The other spelling and syntax errors have also been checked and corrected. See Line 430 in the "Tracked change" manuscript.

5. Lines 304-306: I would re-phrase this sentence. It is a minor distinction, but it currently reads as if you have studied energetic O+ across the entire magnetopause during substroms and found that the most prevalent region of O+ is the dusk flank during the recovery phase. Whereas, it should be more like, "Observations of energetic O+ at the dusk flank magnetopause during substorms are mainly found within the recovery phase."

**Response:** Thanks for referee's nice suggestion. We expand the statistical and found in our 57 events of energetic  $O^+$  at the dusk flank magnetopause, there are 26 events during the expansion phase of intense substorms. While there are 31 events during the recovery phase of intense substorms. So this sentence is not meaningful and has been removed in the revised manuscript.

1	MMS observations of energetic oxygen ions at the low-latitude duskside
2	magnetopause during intense substorms
3	
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# 17 Abstract

18 Energetic oxygen ions  $-0^+$  (>1keV1-40 keV) observed by the Magnetospheric Multiscale (MMS) 19 satellites at the dusk flankduskside magnetopause boundary layer during the phase 1 - la and 1b are 20 investigated. There are 31-57 dusk flankduskside magnetopause crossing events during intense substorms 21 (AE>500 nT) are identified. These 31-57 events of energetic O<sup>+</sup> at the dusk flankduskside magnetopause 22 include 9-26 events during the expansion phase and  $\frac{22-31}{2}$  events during the recovery phase of intense 23 substorms. It is found that the O<sup>+</sup> density in the duskside magnetopause boundary layer during the 24 recovery phase (0.081 cm<sup>-3</sup>) is larger than that during the expansion phase (0.069 cm<sup>-3</sup>). It is noted that 9 25 The 26 events of energetic O<sup>+</sup> ion at the dusk flankduskside magnetopause during intense substorms 26 expansion phase are all under the southward interplanetary magnetic field (IMF)-conditions. There are 27 only 7 events under northward IMF and they all occurred during the intense substorm recovery phase. 28 The number density of energetic O<sup>+</sup> at the dusk flankduskside magnetopause ranges from 0.010.007 cm<sup>-</sup> 29 <sup>3</sup>-to 0.20.599 cm<sup>-3</sup>. The maximum number density ratio of O<sup>+</sup>/H<sup>+</sup> is -0.055 occurred during the intense 30 substorm recovery phase and with AE index about 610 nT and under the norsouthward interplanetary 31 magnetic field (IMF). When the IMF is southward, the O<sup>+</sup> density shows an exponential increase with 32 the IMF  $B_z$  absolute value. Meanwhile, The number density ratio of  $O^+/H^+$  density ratio also shows an 33 exponential increase growth with the IMF By. These results agree with previous studies in the near-Earth 34 magnetosphere during intense substorm. It is suggested that O<sup>+</sup> abundance in the duskside magnetopause 35 boundary layer has a close relation with O<sup>+</sup> variations in the near-Earth magnetosphere during intense 36 substorms. While IMF Bz seems play a minor role in O<sup>+</sup> abundance at the dusk flank magnetopause 37 during intense substrom. Our observations suggest that energetic oxygen ions play a key role in the mass 38 and energy transferring from the tail to the dayside in the magnetosphere during intense substorms.

## 39 1 Introduction

40 Single charged oxygen ions  $(O^+)$  in the magnetosphere,  $\Theta^+$ , are exclusively from the ionosphere. 41 They are an important element in the mass and energy transport in the magnetospheric dynamic process 42 during intense substorms (AE >500 nT), especially during the expansion phase and recovery phase of 43 intense substorms (e.g., Daglis et al., 1991; 1996; Duan et al., 2017; Fok et al., 2006; Ohtani et al., 2011; 44 Ono et al., 2009; Nosé et al., 2000; Yau et al., 1997; 2012; Kronberg et al., 2014). Processes in the 45 magnetotail due to substorm can result in auroral electrojet activity. This activity is generally caused by field-aligned currents increase and reflected by the AE index (Tang and Wang, 2018). Previous studies 46 47 have found that Tthe number density and energy density of oxygen ions, O+, significantly increase with 48 AE index in the near-Earth magnetosphere during the intense substorm-expansion phase (e.g., 49 Lennartsson and Shelley, 1986; Daglis et al., 1991, 1994; Duan et al., 2017; Lennartsson and Shelley, 50 1986). Lennartsson and Shelley (1986) proposed that oxygen ions with energies less than 17 keV/e could 51 provide 50% of the density in the plasma sheet during disturbed geomagnetic activity. They found the 52 increase in the O<sup>+</sup> energy density was strongest around local midnight where O<sup>+</sup> became the most 53 abundant ion at AE~1000 nT. In the near-Earth plasma sheet (NEPS), The O<sup>+</sup> energy density has an 54 explosively increases with AE index in the range of larger than 500 nT during the intense substorm 55 expansion phase (Daglis et al., 1994). Previous studies reported that the O<sup>+</sup> from the nightside auroral 56 region could rapidly feed in the near-Earth magnetosphere during substorm expansion phase (e.g., Daglis 57 and Axford, 1996; Duan et al., 2017; Yu et al., 2013). Otherwise, the solar wind dynamic pressure also 58 influences the oxygen content of ion outflow from the ionosphere. Using the Thermal Ion Dynamics 59 Experiment (TIDE) on the Polar satellite. Elliott et al., (2001) found both the O<sup>+</sup> density and parallel flux
60 increased with the solar wind dynamic pressure.

61 The O<sup>+</sup> outflowing from the ionosphere with low energy of eV are accelerated to about 500 eV at 62 the high altitude polar region (e. g., Yau and André, 1997). Then they are convected tailward into the 63 lobe and the plasma sheet boundary layer. After O<sup>+</sup> enter the NEPS of magnetotail, they can be energized 64 up to tens of keV during intense substorm dipolarizations (e.g., Birn et al., 2013; Duan et al., 2017; 65 Fok, et al., 2006; Nosé et al., 2000; Ono et al., 2009; Yau et al., 2012). Daglis et al. (1991) reported that 66 the energy density of O<sup>+</sup> increased non-linearly with AE indices during substorms. The inductive electric 67 field associated with substorm dipolarization is very significant for accelerating particles in the near-68 Earth plasma sheet<u>NEPS</u> (e.g., Dai et al., 2014; 2015; Duan et al., 2011, ;2016; Lui et al., 1999). 69 Lennartsson and Shelley (1986) proposed that oxygen ions with energy of less than 17 keV/e can provide 70 50% of the number density in the plasma sheet during activity geomagnetic disturbance. Duan et al. 71 (2017) reported that the O<sup>+</sup> from the lobe or the plasma sheet boundary layer were efficiently accelerated 72 by the kinetic Alfven eigenmode with significant unipolar electric field and rapidly feed in the NEPS 73 during intense substorm dipolarizations. These energetic O<sup>+</sup> in the NEPS can be injected into the inner 74 magnetosphere and drift westward into the duskside outer magnetosphere (e.g., Ganushkina et al., 2005). 75 The oxygen ions from the nightside auroral region can rapidly feed in the near Earth magnetosphere 76 during substorm expansion phase (e.g., Daglis and Axford, 1996; Duan et al., 2017; Yu et al., 2013).

77 Oxygen ions decay from the ring current and can leaked into the dayside low latitudemagnetopause 78 boundary layer (e.g., Li et al., 1993; Ebihara\_et al., 2011). Li et al. (1993) reported that the ring current 79  $O^+$  with tens of keV energy interacted with the Pc 5 waves and then lost towards the dayside 80 magnetopause. Kim et al. (2005) pointed out that The solar wind dynamic pressure enhancement playsed 81 a key role in the loss of the ring current particle loss into the outer magnetosphere. This pressure 82 enhancement pushing the magnetopause to move inward leads to a reduction of the scale length of the 83 magnetic field magnitude gradient along the magnetopause. The magnetic gradient drift speed across the 84 magnetopause will increase. So the ring current oxygen ions along the magnetic gradient drift path can 85 easily enter into the outer magnetosphere (Kim et al., 2005). Ebihara et al. (2011) proposed that the field 86 line curvature scattering is was more effective on the loss of energetic oxygen ions with its large gyro-87 radius. These energetic oxygen ions with pitch angle of ~90 degrees can are more prone to leak into the 88 dayside magnetopause. This process may be the key loss mechanism for the ring current ions.

89 The distribution of energetic oxygen ions density at the dayside magnetopause is asymmetry 90 distribution and it has a close relationship with the interplanetary magnetic field (IMF) and solar wind 91 dynamic pressure (e.g., Bouhram et al., 2005; Phan et al., 2004; Luo et al., 2017). Bouhram et al. (2005) 92 pointed out that the O<sup>+</sup> density in the duskside (on average 0.053 cm<sup>-3</sup>) magnetopause is higher in the 93 duskside than that in the dawn-side (on average 0.014 cm<sup>-3</sup>) magnetopause. They found  $O^+$  is was the 94 dominant contributor to the mass density (30%) on the dusk-side magnetopause in comparison to 3% in 95 the dawnside and 4% near the noon. The dawn-dusk asymmetries of the energetic O<sup>+</sup> (>~274 keV) 96 distribution in three different regions (dayside magnetopause, near-Earth nightside plasma sheet, and tail 97 plasma sheet) are also observed by Luo et al., (2017). They found that the energetic O<sup>+</sup> distributions were 98 mainly influenced by the dawn-dusk IMF directions and the enhancement of ion intensity strongly related 99 to the location of the magnetopause reconnection. Phan et al. (2004) reported that energetic O+ with 100 energy larger than 3 keV in the reconnection jets are observed by Cluster at the duskside mid latitude 101 magnetopause under steady southward interplanetary magnetic field (IMF) conditions.

102 There is ample evidence that magnetospheric ions could participate in the magnetopause 103 reconnection and directly escape along the reconnected open field lines (e.g., Sonnerrup et al., 1981; 104 Fuselier et al., 1991, 2016; Slapak et al., 2012, 2015; Wang et al., 2014; Liu et al., 2015). The energetic 105 O<sup>+</sup> with energies larger than 3 keV in the reconnection jets at the duskside mid-latitude magnetopause 106 under steady southward IMF was reported by Phan et al. (2004). -(Zong et al., (2001) observed O<sup>+</sup> 107 energy dispersion due to time-of-flight (TOF) effects at the dusk flankduskside magnetopause under 108 southward IMF and it was assumed that O<sup>+</sup> was escaping from the ring current along the reconnected 109 field lines during steady reconnection. However, However, Fuselier et al. (1989) reported that  $O^+$  from 110 the high latitude ionosphere were not associated with any substorm cycle. O<sup>+</sup> from the high latitude 111 ionosphere could form the O<sup>+</sup> rich boundary layer in the low latitude magnetopause. While When O<sup>+</sup> enter 112 the reconnection jets, the reconnection rate is likely reduced by the mass-loading, reconnection is but not 113 suppressed at the magnetopause\_(-Fuselier et al. 2019).

114 At present, <u>variations of O<sup>+</sup> abundance (O<sup>+</sup>/H<sup>+</sup>) in the duskside magnetopause boundary layernear</u> 115 the dayside low latitude magnetopause during intense substorm (AE >500 nT) on AE index and solar 116 wind conditions (e.g. IMF  $B_y$ , IMF  $B_z$  and solar wind dynamic pressure) expansion phase and recovery 117 phase-are still not <u>understood clear</u>. Previous studies of <u>O<sup>+</sup> abundance variationsoxygen ions</u> during 118 substorms are mainly focused on <del>O<sup>+</sup> energizations in the near Earth plasma sheet in</del> the magnetotail <u>or</u> 119the near-Earth region (e.g., Duan et al., 2017; Nosée et al., 2000; Ohtani et al., 2011). The120Magnetospheric Multiscale (MMS) mission gives us an opportunity to focus on the O<sup>+</sup> in the duskside121magnetopause region. In this study, Wwe investigate statistical features of energetic oxygen ionsO<sup>+</sup> at122the dusk flankduskside magnetopause and their relations with AE index and solar wind conditions (e.g.123IMF B<sub>y</sub>, IMF B<sub>z</sub>, and solar wind dynamic pressure)O<sup>+</sup> in the near Earth plasma sheet during the intense124substorms (AE >500 nT).

## 125 2 Instrumentation and Data

#### 126 - Instrumentation and data

#### 127 2.1 Instrumentation and data

128 This study uses-used\_data from the Magnetospheric Multiscale (MMS) mission. This mission 129 comprises four identical satellites that were launched on 2015 March 13 into an elliptical 28-inclination 130 orbit with perigee around 1.2 R<sub>E</sub> and apogee around 12 R<sub>E</sub> (Burch et al., 2016; Fuselier, et al., 2016b). 131 The electric field **E** is from the electric double probe (EDP) (Ergun et al., 2016; Lindqvist et al., 2016), 132 and magnetic field **B** is from the Fluxgate Magnetometer (FGM) (Russell et al., 2016). The plasma data 133 are from the Fast Plasma Investigation (FPI) (Pollock et al., 2016)-and the Hot Plasma Composition 134 Analyzer (HPCA)-(Young et al., 2016). The FPI provides plasma (electrons and ions) distribution 135 functions at 32 energies from 10 eV to 30 keV. And it has a high time resolution of 0.03 s for electrons 136 and 0.15 s for ions in the burst mode and 4.5 s in the fast mode (Pollock et al., 2016). The FPI does not 137 discriminate between different ion species. While T the HPCA provides ion composition (H<sup>+</sup>, He<sup>++</sup>, He<sup>++</sup>, He<sup>++</sup> 138 and O<sup>+</sup>) measurements in the energy range from -1 eV/q to 40 keV/q (Young et al., 2016). Although the 139 HPCA instrument employs radio frequency (RF) unit to artificially reduce the proton fluxes in some 140 areas where the proton fluxes are intense, there still exists a low level of background that affects the O<sup>+</sup> 141 fluxes in the magnetosheath. The majority of the  $O^+$  fluxes in the magnetosphere side of the 142 magnetopause are at energies above-from 1 keV to 40 keV and that band below 1keV visible in the 143 magnetosheath side are observations outside the RF operating range and contamination from high proton 144 fluxes. Due to this contamination, the  $-number O^+ density of O^+ with at - energy is from -above-1 keV to$ 145 40 keV is in the magnetopause boundary layer are considered at the magnetopause in our study. The O<sup>+</sup> density recalculated from the HPCA distribution functions at this energy range by using the Space
Physics Environment Data Analysis System (SPEDAS) software package. More details about SPEDAS
can be found in Angelopoulos et al. (2019). The solar wind parameters, IMF and AE index are available
from the OMNI data in CDAweb (http://cdaweb.gsfc.nasa.gov/). The data from the MMS 4 satellite are

adopted in our investigation since the data difference from other three spacecraft is negligible. <u>This is</u>

- 151 <u>due to spacecraft separation and scales of particle motion.</u>
- 152 2.2 Energetic O+ at the dusk flank magnetopause during intense substorms

### 153 2.2.1-3 Results

154 <u>3.1</u> Detailed event on 3 October 2015

155 Figure 1 presents the three components of the IMF in Geocentric Solar Magnetospheric (GSM) 156 coordinates, solar wind dynamic pressure and, as well as AU, AL, and AE index during the interval time 157 of interest from 14:30 UT to 16:30 UT on 3 October 2015. During this interval, the IMF B<sub>x</sub> component 158 is negative all the time (Figure 1a). Its maximum value is about -1 nT at ~15:16 UT. The IMF By 159 component is almost negative except at ~14:32 and ~16:23 UT. The negative IMF Bz component is also 160 observed negative-during this interval as shown in Figure 1c. The minimum value of the IMF  $B_z$ 161 component is about  $\frac{-6.77.1}{-6.77.1}$  nT at  $\frac{-14:30}{-0.00}$  UT. The solar wind dynamic pressure is <u>only at the beginning</u> 162 of the time interval about 2 nPa. Then, it increases sharply at 15:1000 UT and reaches its maximum value 163 about 4.4 nPa at ~15:12 UT. The maximum solar wind dynamic pressure is about 4 nPa. An intense 164 substorm (AE > 500 nT) occurred under the above IMF and solar wind dynamic pressure. The substorm 165 onset time is about 14:45 UT, when the AE index significantly increases. The maximum value of AE 166 index is  $\sim 1000$  nT at 15:20 UT. These solar wind conditions led to an intense substorm (AE > 500 nT), 167 as Figure 1g shown. The AE index is defined as AE=AU-AL. Generally, the substorm onset time is 168 characterized by the AL index starting to significantly decrease and the AE index significantly increase. 169 The interval of the AL index decreasing from onset to its minimum is defined as the substorm expansion 170 phase. The interval of the AL index increasing from the minimum to the quiet time level is regarded as 171 the substorm recovery phase. From Figure 1e to 1g, the substorm onset time is about 14:45 UT marked 172 by the AL index starting to sharply decrease and AE index increasing. After the AE index significantly 173 increases and the AL index decreases (Figure 1f), the AL and AE indexes reach their minimum and 174 maximum values about -750 nT and 1000 nT at ~15:20 UT, respectively. This interval from ~14:45 to ~15:20 UT is regarded as the intense substorm expansion phase. Then, the intense substorm enters the 176 recovery phase as the AL index gradually increases and AE index decreases after ~15:20 UT. The two 177 blue dashed lines indicate the time interval of the magnetopause boundary layer crossing. According to 178 the above description, we can identify this magnetopause boundary layer crossing occurred during the 179 recovery phase of intense substorm. The identification of the magnetopause boundary layer will be 180 described later.

181 Figure 2 shows that MMS 4 encountered the duskside magnetopause region the overview of the 182 magnetopause inbound crossing from 15:25:1015:00 to 15:36:5016:00 UT-during the intense substorm 183 recovery phase on 03 October 2015. This interval is marked by the two vertical blue lines in Figure 1. 184 This intense substorm onset time is around 14:45 UT with AE maximum value 1000 nT at 15:20 UT as 185 shown in Figure 1e. During the magnetopause crossing, MMS 4 satellites were was located at about ((6.0, 186 8.8, -5.1) Re  $R_E$  in GSM as shown in the bottom of Figure 2. From top to bottom, panels 2a and 2b show 187 that the magnetic and electric fields in GSM from FGM and EDP, respectively. Ion and electron 188 temperatures, ion and electron plasma density-number densities, and ion velocity in GSM from FPI L2 189 data products are shown in Figures 2e, 2d and 2e2c-e, respectively. Figure 2f shows the H<sup>+</sup> and O<sup>+</sup> 190 densities, followed by <u>-Tthe electron and ion omnidirectional differential energy fluxesomnidirectional</u> 191 electron and ion energy fluxes from FPI are shown in (Figure 2g-h and 2h)., respectively. The O+number 192 density at energies above 1keV is displayed in Figure 1f. The last four panels\_of Figure 2 show present 193 the omnidirectional differential fluxes of four individual ion species, H<sup>+</sup>, O<sup>+</sup>, He<sup>+</sup>, and He<sup>++</sup> measured by 194 HPCA, respectively. The HPCA flux in panels 2i-l has artificial striping every 4 energy bins due to way 195 HPCA determines the count rate over 4 energy channels in survey mode. It is noted that the differential 196 fluxes (Figure 2i-l) and differential energy fluxes (Figures 2g-h) have different units. To better identify 197 the fluxes variations at specific energies, we choose the ion and electron fluxes from FPI in the energy 198 flux unit. The plasma moments (e.g. ion parallel and perpendicular temperatures, ion and electron 199 densities, and ion velocity) from FPI shown in Figures 2c-e are all from MMS L2 data products. They 200 are default moments calculated over the full FPI energy range from 10 eV to 30 keV. Note that in the 201 magnetosheath, O<sup>+</sup> measurements suffer from fake counts at energies below 1 keV which results from 202 high proton fluxes contamination, as the red box in Figure 2j shown. So the spurious counts should be excluded in the plasma moments calculation. The O<sup>+</sup> density shown in Figure 2f is recalculated from 203

204 <u>HPCA distribution functions at energies from 1 keV to 40 keV. Due to H<sup>+</sup> measurements from HPCA is
205 accurate and the H<sup>+</sup> mean energy in the magnetosheath is typically 0.3 keV, we adopted the default H<sup>+</sup>
206 density from HPCA L2 data products which computed over the full HPCA energy range from 1 eV to
207 40 keV, as the red line shown in Figure 2f.
</u>

208 The different regions encountered by MMS4 during the interval of 15:00 to 16:00\_UT are marked 209 by the colored bar at the top of Figure 2, with the magnetosheath shown in orange, the outer 210 magnetosphere shown in blue, and the magnetopause region boundary laver shown in green. From 211 15:00:00 to 15:25:10 UT, MMS4 was located in the magnetosheath. This region is characterized by a the 212 southward magnetic field, low ion and electron temperatures (a few hundred a few hundred eV for ions 213 and tens of eV for electrons, Figure 2c) with relatively high number densities (on the order of  $\sim 20$  cm<sup>-3</sup>, 214 Figure 2d), and stable ion flow speeds of about one hundred 100 km/s. There are also very high fluxes at 215 energies centered around 100 eV (nominal magnetosheath energy) for electrons (Figure 2g) and at 216 energies centered around 1 keV\_-for ions (Figure 2h, also see H<sup>+</sup> fluxes in Figure 2i and He<sup>++</sup> fluxes in 217 Figure 21) in the magnetosheath. While the  $O^+He^+$  and  $O^+He^+$  fluxes above 1 keV nearly disappear in 218 the magnetosheath (Figure 2j and 2k). From Figure  $\frac{2k}{2}$ , the majority of the-O<sup>+</sup> fluxes at energies below 219 1 keV visible in the magnetosheath are the result of contamination from the high proton fluxes, as the 220 red box indicated.

221 The primary magnetopause crossing from the magnetosheath into the magnetosphere lastsed about 222 6-12\_min, from 15:31:20about 15:25:10 to 15:36:50 UT. Partial encounters of the magnetopause by 223 MMS4 occurred around 15:43:15, 15:47:10 and 15:53:00 UT and etc., respectively. The magnetopause 224 region-boundary layer is identified by the plasma momentsparticle fluxes and the electromagnetic field. 225 The plasma density and ion-temperature at the magnetopause are between the corresponding values of 226 the magnetosphere and the magnetosheath, as shown in Figure 2d and 2c, respectively. The 227 magnetopause region boundary layer can also be identified by the significant increases in electron fluxes 228 at energies about several hundreds of eV and ion fluxes at energies around  $\sim 10 \text{ keV}$ , as shown in Figure 229 2g and 2h, respectively. During the interval from 15:25:10 to 15:36:50 Ut this time of interest, the B<sub>z</sub> 230 component rotated from southward to northward and back again several times before finally became 231 northward when MMS 4 entered the magnetosphere. The energetic O<sup>+</sup> number density (1-40 keV) is 232 around 0.02018 cm<sup>-3</sup> within the magnetopause boundary layer as shown in Figure 2f. The corresponding 233 H<sup>+</sup> and O<sup>+</sup> fluxes at specific energies and their densities (shown in Figure 2f) were averaged in this region.

234 After 15:36:50 UT, MMS4 entered the magnetosphere which is identified by the observations of a 235 the northward magnetic field (Figure 2a), much lower number plasma densities (on the order of ~1 cm<sup>-</sup> 236 <sup>3</sup>) with respect to the densities in the magnetosheath (Figure 2d), higher ion-plasma temperatures (-Figure 237 <u>2c</u>, several a few thousands of keV for ions and a few hundred eV for electrons), and a small bulk ion 238 flow speed. Higher fluxes at energies-around several keV for electrons (Figure 2g) and at energies 239 centered around ~10 keV for ions (Figure 2h) also indicate that the MMS4 was in the magnetosphere. 240 Finally, the presence of  $O^+$ He<sup>+</sup>-and  $O^+$ He<sup>+</sup> at energies about ~10 keV is also used as a marker to verify 241 that MMS4 was in the magnetosphere (Figure 2j and 2k).

# 242 2.2.2.3.2 Statistical\_<u>-3157</u> events of <u>MMS first Phase 1 observations of energetic O<sup>+</sup> at the <del>dusk</del> flankduskside magnetopause-crossing during intense substorms(<u>AE>500nT</u>) </u>

244 Based on the in-situ measurements of the dayside magnetopause processes crossings by MMS 245 satellites in phase 1a and phase 1b, we identified the duskside magnetopause crossing event (complete 246 magnetopause crossing from the magnetosheath to the magnetosphere, vice versa) from the summary 247 plot in https://lasp.colorado.edu/mms/sdc/public/plots/. Then we plotted the more detailed overview 248 figure of these events to identify the magnetopause boundary layer, as Figure 2 shown. Only events with 249 AE index larger than 500 nT during the magnetopause boundary layer crossings interval were selected. 250 There are 31-57 events of the dusk flankdusksideside magnetopause boundary layer crossing during 251 intense substorm during intense substorms with AE index larger than 500 nTsatisfied with the above 252 criterion-are selected. In our statistical study, the mean values of the  $H^+$  and  $O^+$  fluxes at specific energies 253 and their densities are calculated in the magnetopause boundary layer. Correspondingly, the solar wind 254 dynamic pressure and IMF  $B_{zBy}$ ,  $B_{yz}$  and AE index from the OMNI data system was were averaged 255 during the magnetopause boundary layer crossing time interval, as the two blue dashed lines shown in 256 Figure 1. The phase of the substorm is determined from the variations of AU, AL and AE indexes, as 257 mentioned before. For better follow-on studies, we add more detail information about 57 energetic O<sup>+</sup> 258 events into an appendix. From the appendix, we can easily draw the conclusion that the O<sup>+</sup> density in the 259 duskside magnetopause during the recovery phase (0.081 cm<sup>-3</sup>) of intense substorm is larger than that 260 during the expansion phase (0.069 cm<sup>-3</sup>).

Figure 3 displays the locations of 57 energetic O<sup>+</sup> events at the duskside magnetopause (-5.7  $R_E$  <  $Z_{GSM} < 1.7 R_E$ ) during intense substorms projected into the XY<sub>GSM</sub> plane of 31 events of energetic O<sup>+</sup> at 263 the dusk flank magnetopause during intense substorms. The blue curve line represents the nominal 264 magnetopause, which is obtained by the magnetopause model of (Shue et al., (1998)) using when the IMF 265  $B_z$  is about ~ 1.96-3.21 nT and solar wind dynamic pressure (Psw) is ~ 2.372.87 nPa (averaged for the 266 31-57 events). The diamond and circle represent the event at the dusk flankduskside magnetopause during 267 the intense substorm expansion phase and recovery phase, respectively. The O<sup>+</sup> number density and the 268 number density ratios of  $O^+/H^+$  density ratio are shown by the colored diamonds and circles at the 269 corresponding magnetopause locations in Figures 3a and 3b, respectively. Among the  $\frac{31-57}{2}$  events of 270 energetic  $O^+$  at the dusk flank dusk side magnetopause during intense substorms, there are 9-26 events that 271 occurred during the expansion phase of intense substorms and 22-31 events occurred during 272 the recovery phase of intense substorms. These 9 and 22 events are shown in Figure 3 by 9 diamonds 273 and 22 circles, respectively. The maximum number density of energetic O<sup>+</sup> is found during the intense 274 substorm recovery phase as presented in Figure 3a. On the other hand, the maximum number density 275 ratio of O+/H+ is also found during the recovery phase of an intense substorm as shown by the red circle 276 in Figure 3b.

277 Figure 4 presents the relationship between the energetic O<sup>+</sup> at the dusk flankdusk side magnetopause 278 and AE index during intense substorms. From top to bottom, panels show that the  $O^+$  and  $H^+$  number 279 density densities, (Figure 4a), the  $O^+/H^+$  number density ratio of  $O^+/H^+$ , (Figure 4b), and  $O^+/H^+$  particle 280 fluxes ratios at different energy ranges - of O+/H+ (Figure 4c), respectively. The energy channel ranges 281 of for  $O^+$  and  $H^+$  in the HPCA are the same. So the  $O^+/H^+$  number density ratio ( $nO^+/ nH^+$ ) and  $O^+/H^+$ 282 particle fluxes ratio are directly defined as the ratio between mean values of their fluxes, respectively. 283 for The particle fluxes are chosen(fO+/fH+) at energies ~1\_keV (energy range from 987.82 to 1165.21 284 eV), ~10 keV (energy range from 99.73.98 to 11.7765.13 keV), ~20 keV (energy range from 19.310.05 285 to 22.777.82 keV and ~35 keV(energy range from 31.693.41 to 37.3985.04 keV). are defined as the 286 ratio between mean values of their number densities and fluxes, respectively. The error bars indicating 287 the 90% confidence interval (CI) are also overplotted in each point. The confidence interval is based on 288 the following formula:

289

 $\overline{x} - k \frac{s}{\sqrt{n}} < \mu < \overline{x} + k \frac{s}{\sqrt{n}}$ 

290 Where  $\overline{x}$ , s and n are the mean value, standard deviation and the sampling number of observations, 291 respectively. k in the above formula can be determined by calculating a 90% confidence interval for each 292 events (the k value is 1.65). Figure 4a shows that the number density of energetic O<sup>+</sup> density at the dusk 293 flankduskside magnetopause during intense substorms is in the range from  $\frac{-0.01-0.007}{-0.010}$  em<sup>3</sup>-to  $\frac{0.20}{-0.20}$ 294  $cm^{-3}$ . The maximum<u>-number</u> density value of energetic O<sup>+</sup> at the <u>dusk flankduskside</u> magnetopause 295 during intense substorm expansion-recovery phase is presented at the higher AE index - about 870-606 296 nT-as displayed in Figure 4a the last right blue circle. While the maximum value of the energetic O+ 297 number density, the ratio of nO+ / nH+, and fO+ /fH+ at the dusk flank magnetopause during intense 298 substorm recovery phase are all around AE index 600 nT. The O<sup>+</sup>/H<sup>+</sup> density ratio decreases with AE 299 index from 900 to 1100 nT. The variations of O<sup>+</sup> density and O<sup>+</sup>/H<sup>+</sup> density ratio with AE index do not 300 show obvious difference between during the expansion phase and the recovery phase.

301 Figure 5 shows the relationship between the energetic  $O^+$  at the dusk flank dusk side magnetopause 302 and the-IMF By during intense substorms. The format of the Figure 5 is the same as that of Figure 4. 303 Figure 5a shows that the O<sup>+</sup> and H<sup>+</sup>-number densities increase decrease with IMF B<sub>y</sub>-when it is larger 304 than 4 nT from -6 to 0 nT and increase with IMF By from 4 to 8 nT. From Figure 5b, the O<sup>+</sup>/H<sup>+</sup>-number 305 density ratio shows an exponential growth with increase in the\_-IMF By is associated with an exponential 306 growth by a factor 0.07. Based on observations the scatter plot in Figure 5b, we can define linear 307 functional dependence between the logarithm of  $O^+/H^+$  - number density ratio and IMF B<sub>ya</sub> as Eq. (1) 308 shown.: And the corresponding correlation coefficients is 94%. The correlation coefficient close to 100% 309 indicates that there is a great correlation.

311

$$\frac{\log(n^{0^{\pm}}/_{H^{\pm}}) = 0.07 * IMF By - 5.14}{\log \frac{n_{0^{+}}}{n_{H^{+}}} = 0.126 * IMF By - 5.174}$$
(1)

The dependency is constructed using a linear <u>least-least-squares fit unless otherwise stated</u>. The O<sup>+</sup>/H<sup>+</sup> particle flux ratio at energy <u>about ~10\_keV</u>, ~20\_keV and ~<del>3035\_keV</del> also show an obvious exponential increase with IMF B<sub>y</sub>. This dependency is consistent with Welling et al. (2011) simulation results found in the ring current.

Figure 6 shows the relationship between the energetic O<sup>+</sup> at the <u>dusk flankduskside</u> magnetopause and\_<u>the</u>IMF B<sub>z</sub> during intense substorms. The format of <u>the</u>Figure 6 is the same as <u>that of</u> Figure 4. Figure 6a and Figure 6b both present that among <u>our 31–57</u> events of energetic O<sup>+</sup>-near <u>at</u> the <u>dusk</u> flankduskside magnetopause boundary layer during intense substorm, there are 50 events under southward IMF and only 7 events under northward IMF.there are 9-It is noted that 26 events occurred 321 during the expansion phase of intense substorms which are all under the southward IMF conditions, as 322 the blue points shown the blue circle in Figure 6. While Meanwhile, there are only 6 \_events that 323 occurred under the northward IMF are all during the intense substorm recovery phase, as the right circle 324 red points with positive IMF  $B_z \rightarrow 0$  shown. From -10 to 0 nT, the O<sup>+</sup> density shows an obvious decrease 325 with IMF Bz. To better describe this variation trend, the empirical functional relation between the 326 logarithm of O<sup>+</sup> density and IMF  $B_z$  (from -10 to 0 nT) is established in Eq.(2) and the corresponding 327 correlation coefficient is 94%. While the  $O^+$  density has a positive correlation with IMF B<sub>z</sub> from 0 to 5 328 nT.

329

$$\log n_{0^+} = -0.163 * \text{IMF Bz} - 3.737 \tag{2}$$

From Figure 6b, the O<sup>+</sup>/H<sup>+</sup> density ratio during the recovery phase decrease with IMF B<sub>z</sub> from about -2
 to 2 nT. The maximum number density of energetic O<sup>+</sup> at the dusk flankduskside magnetopause is during
 the intense substorms recovery phase \_\_under the southward IMF. But Meanwhile, the maximum O<sup>+</sup>/H<sup>+</sup>
 density ratio of nO+ / nH+ at the dusk flankduskside magnetopause is also during intense substorm
 recovery phase under the northward southward IMF. Comparing with IMF By, IMF Bz seems play a
 minor role in O+ abundance at the dusk flank magnetopause during intense substrom.

Figure 7 displays the relationship between the energetic O<sup>+</sup> at the <u>dusk flankduskside</u> magnetopause and the solar wind dynamic pressure during intense substorms. The format of the Figure 7 is the same as <u>that of</u> Figure 4. Figure 7a and 7b both presents that the <u>number O<sup>+</sup></u> density of energetic O + and the O + to H<sup>+</sup> number density ratio at the <u>dusk flankduskside</u> magnetopause during intense substorms have has a a weak positive correlation with the solar wind dynamic pressure. The empirical functional relation between the logarithm of O<sup>+</sup> density and solar wind dynamic pressure (from 1 to 4.5 nPa) is also established in Eq.(3) and the corresponding correlation coefficient is 94%.

343

$$\log n_{0^+} = 0.325 * \text{Psw} - 4.061 \tag{3}$$

From Figure 7b, the O<sup>+</sup>/H<sup>+</sup> density ratio during recovery phase show a decrease from about 2.5 to 3 nPa.
It is noted that the O<sup>+</sup>/H<sup>+</sup> density ratio increase with solar wind dynamic pressure from about 3 to 4 nPa.
The maximum number density of energetic O<sup>+</sup> at the dusk flankduskside magnetopause, ---0.2 --0.599
cm<sup>-3</sup> take place at solar wind dynamic pressure <u>is 1.9 about 3.9</u> nPa during an intense substorm recovery
phase. While the maximum O<sup>+</sup>-to-/H<sup>+</sup> number density ratio at the dusk flankduskside magnetopause is
around appeared at solar wind dynamic pressure around 2.5 2.2 nPa. More details can be found in the
appendix.

352 **43** Discussion

353 Energetic O<sup>+</sup> (1-40 keV) with high number density are observed by the MMS satellites at the dusk 354 flankduskside magnetopause during the expansion phases and recovery phases of intense substorms, as 355 356 0.010.007 cm<sup>-3</sup> to 0.2599 cm<sup>-3</sup> at the duskside flank-magnetopause boundary layer during intense 357 substorms. This O+ number density is larger than the energetic O+ near the duskside mid latitude 358 magnetopause observed by Cluster, which was 0.011 to 0.053 cm 3 (e.g., Bouhram et al., 2005). 359 Bouhram et al. (2005) demonstrated that the O+ number density is much higher in the dusk side than in 360 the dawn side magnetopause. In a companion paper from Zeng et al. (2019), they study the O<sup>+</sup> abundance 361 variations on the solar wind conditions at the dayside magnetopause boundary layer and not specific to 362 the events that occurred during intense substorm. The mean value of the O<sup>+</sup> density at the duskside 363 magnetopause boundary layer is  $0.038 \text{ cm}^{-3}$  in that paper. While during the intense substorm, the O<sup>+</sup> 364 density increase to 0.075 cm<sup>-3</sup> in this study. The significantly higher number density of energetic O+ is 365 obtained from MMS at the dusk flank magnetopause during intense substorms in our paper. There are 366 two reasons for this high number density of energetic  $O^+$  observed during the intense substorm. The first 367 is the time interval for the observations. Our observations are during intense substorms expansion phase 368 and recovery phase. Daglis et al. (1991) proposed that energetic O<sup>+</sup> were significantly higher in the near-369 Earth plasma sheetNEPS in the magnetotail after intense substorms onset. The impulsive electric field 370 accompanied by intense substorm dipolarization plays a key role in the energization and sunward transfer 371 of oxygen ions in the duskside of midnight plasma sheet in the magnetotail (e.g., Fok et al., 2006; Nosée 372 et al., 2000). These energetic O<sup>+</sup> (tens of keV) can be transported sunward into the low latitudeduskside 373 magnetopause boundary layer. The second reason for the high densities is the locations of the 374 observations. Our observations are near the dusk flankduskside magnetopause. This region is easily 375 accessible by energetic O<sup>+</sup> during intense geomagnetic activity (Fuselier et al. 2016a). Phan et al. (2004) 376 pointed out that energetic O<sup>+</sup> were observed with very high number density 0.2-0.3 cm<sup>-3</sup> in the 377 reconnection jets at the dusk-side mid-latitude magnetopause were observed by Cluster.

351

378 During dynamic periods and intense substorms time, light ions yielded more symmetric patterns in 379 density than heavy ions and the O<sup>+</sup> patterns in the active plasma sheet are a function of IMF conditions 380 (Winglee and Harnett 2011. Winglee et al. 2005). Welling et al. (2011) used multispecies MHD and the 381 PWOM to drive a ring current model, and found that positive IMF By pushing pushed the stronger O<sup>+</sup> 382 concentrations toward the duskside at a geocentric distance of about 6.6 R<sub>E</sub>. This O<sup>+</sup> density duskward 383 preference with positive IMF  $B_y$  in the near Earth plasma sheet<u>NEPS</u> is similar to our results. It may 384 indicate that the O+ in the dusk flank magenetopause magnetopause boundary layer O+-enhancing with 385 IMF  $B_v$  is due to the local time variations of  $O^+$  in the <u>near Earth plasma sheetNEPS</u> tied to IMF  $B_v$ . 386 While IMF Bz seems play a minor role in O<sup>+</sup> abundance at the dusk flank magnetopause during intense 387 substrom. These results showOur result of O<sup>+</sup> density increase with IMF B<sub>y</sub> also agree with Kronberg et 388 al., (2012). They showed for 10 keV O<sup>+</sup> strong increasing under the duskward IMF indicated by the clock 389 angle in the inner magnetosphere. It is suggested that the O<sup>+</sup> abundance at the dusk flankduskside 390 magnetopause have has a corresponding relation with the O<sup>+</sup> in the duskside near-Earth plasma 391 sheetmagnetosphere during intense substorm. The O<sup>+</sup> path from the cusp to the magnetotail is asymmetry 392 and it has the best correlation with the IMF directions. This path asymmetry mainly controlled by the 393 IMF  $B_v$  may influence on O<sup>+</sup> abundance at the duskside magnetopause. When the IMF  $B_v$  is positive, the 394 O<sup>+</sup> from northern/southern cusp tends to flow toward the dawnside/duskside. The transport path for 395 negative IMF  $B_y$  is more symmetric but shows some evidence for a reversed asymmetry when the 396 negative IMF  $B_y$  is large enough. While the IMF  $B_z$  has little influence on the asymmetry (Liao et al., 397 2010). 398 Due to not enough events occurred under northward IMF were observed, the influence of IMF Bz 399 on the O<sup>+</sup> abundance (1-40 keV) during intense substorm is not clear. While in Luo et al. (2017) study 400 that not only the cases for the intense substorm were considered, they found the  $O^+$  intensity (>  $\sim$  274 keV) 401 was significantly higher under southward IMF than that under northward IMF especially at the duskside 402 magnetopause. Zeng et al. (2019) also showed that the duskside asymmetry of O<sup>+</sup> density (1-40 keV) in 403 the dayside magnetopause under northward IMF was less obvious than under southward IMF when the 404 IMF By was the same. Under the southward IMF, the interactions between the solar wind and the 405 magnetosphere become active. The inductive electric field or magnetic field gradient related to magnetic

406 reconfiguration will enhance with negative IMF B<sub>z</sub>. So the large scale dawn-dusk electric field drift along

407 with the gradient-curvature drift can force oxygen ions convect to the duskside magnetopause boundary
408 layer (Kronberg et al., 2015; Luo et al., 2017).

409 In this statistical study, there are 50 magnetopause boundary layer crossing events during intense 410 substorm under southward IMF with respect to 7 events under northward IMF. Choosing the intense 411 substorm may increase the probability of observing the events under southward IMF quite significantly. 412 Among our 31-57 events of energetic O<sup>+</sup> near the dusk flankduskside magnetopause, there are 9-26 events 413 during intense substorm expansion phase which are all under the southward IMF, as the blue circle shown 414 in Figure 6b. Hsu, T. - S., and McPherron, R. L. (2003) found that approximately 60% of all substorms 415 expansion phase are correspond to the northward IMF Bz. But,tThere are only 6-7 events under the 416 northward IMF in our study and they all occurred during the intense substorm recovery phase. But what 417 relation between the IMF directions and phase of substorm is out of scope for this article. Thus our 418 statistical results demonstrate that the energetic O+ near the dusk flank magnetopause dominated 419 occurring under the southward IMF.

420 Previous researches demonstrated that the oxygen ions origin from the aurora region can could 421 rapidly feed in the near Earth plasma sheetNEPS -in the magnetotail during intense substorms expansion 422 phase (e.g., Daglis and Axford, 1996; Duan et al, 2017; Yu et al., 2013). Oxygen ions can be efficiently 423 energized in the near Earth plasma sheet NEPS during intense substorm dipolarization (e.g., Duan et al., 424 2017; Fok et al., 2006; Nosée et al., 2000). Under southward IMF conditions, these energetic oxygen 425 ions in the NEPS can be convected sunward and drift westward. As a result, the energetic O<sup>+</sup> arrived near 426 the dusk flankduskside magnetopause during intense substorm expansion phase could can participate in 427 the magnetopause reconnection and escape along reconnected field lines<u>during intense substorm</u> 428 expansion phase, as reported by Wang et al. (2014) and Zong et al. (2001). But from Fuselise et al. (2019), 429 whileWhen O<sup>+</sup> participate in the reconnection jets, the reconnection rate <u>-is will</u> likely be reduced by the 430 mass-loading, reconnection is but not suppressed at the magnetopause (Fuselier et al. 2019). Whether 431 these energetic O<sup>+</sup> at the duskside boundary layer could suppress the intense substorm need further 432 investigation.

433 **<u>5</u>** Summary and conclusions

434 Using the measurements from FPI, HPCA and FGM on MMS mission satellite during the phase 1a 435 and 1b, we have studied 31-57 events of the energetic oxygen ions  $O^+$  (1-40 keV) - distributions at the 436 dusk flankduskside magnetopause boundary layer and their variations on the solar wind conditions (IMF 437  $B_y$ , IMF  $B_z$  and solar wind dynamic pressure) during intense substorm expansion phases and recovery 438 phases. According to the above analysis, we can draw our main conclusions as follows. During intense 439 substorms, the energetic oxygen ions are mainly observed at the dusk flank magnetopause within the 440 substorm recovery phase. In our 31-57 events of energetic O<sup>+</sup> at the dusk flankduskside magnetopause 441 boundary layer, there are only 926 events during the expansion phase of intense substorms and 31 events 442 during the recovery phase. While there are 22 events during the recovery phase of intense substorms. It 443 is noted that the mean values of the  $O^+$  density during the expansion phase and recovery phases are 0.069 444  $cm^{-3}$  and 0.081 cm<sup>-3</sup>, respectively. And the maximum O<sup>+</sup>/H<sup>+</sup> density ratio occurred during the intense 445 substorm recovery phase. We find out It is found that 9-26 events of energetic O<sup>+</sup> at the dusk flank duskside 446 magnetopause during intense substorms expansion phase are all under the southward IMF conditions, 447 and only 7 events under northward IMF which are all during the intense substorm recovery phase. The 448  $O^+$  density shows an exponential increase with IMF B<sub>z</sub> absolute value under the southward IMF. 449 Similarly, it also presents an exponential growth with solar wind dynamic pressure, and the empirical 450 functional relations are established. Like previous studies during substorm in the near-Earth 451 magnetosphere, The  $O^+/H^+$  number density ratio in the duskside magnetopause boundary layer 452 exponential increase enhance with the IMF By and the maximum number density ratio of oxygen ion to 453 proton is ~0.055 during intense substorm recovery phase. It is suggested that the O<sup>+</sup> abundance in the duskside magnetopause boundary layer has a close correlation with the O+ variations in the near-Earth 454 455 magnetosphere during intense substorm. The high energetic O<sup>+</sup> number density at the dusk flank 456 magnetopause during intense substorms is in the range from 0.01 cm<sup>-3</sup> to 0.2 cm<sup>-3</sup>. Our observations 457 suggest that energetic oxygen ions are a key indicator in the mass and energy transport from the tail to 458 the dayside in the magnetosphere during intense substorms.

# 459 Data availability

460 All data used in this study are publicly accessible. MMS data are available at the MMS Science Data Center (https://lasp.colorado.edu/mms/sdc/public/). The OMNI data can be downloaded from the 461 462 NASA Goddard Flight Coordinated Space Center Data Analysis Web 463 (CDAWeb:http://cdaweb.gsfc.nasa.gov/).

#### 464 Competing interests

465 The authors declare that they have no conflict of interest.

## 466 Author contribution

467 C. Z. conducted the majority of the data processing, analysis and writing for this study. S.P.D, C.W,
468 L.D and S.F participated in the interpretation of the data and modified this paper. J.B, R.T, B.G and C.R
469 produced the data and controlled the data quality. All the authors discussed the results and commented
470 on the paper.

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#### 477 References

- Angelopoulos, V., Cruce, P., Drozdov, A. et al. (2019), The Space Physics Environment Data Analysis
   System (SPEDAS), Space Sci Rev., 215(9), doi:10.1007/s11214-018-0576-4.
- Birn, J., Thomsen, M. F., and Borovsky, J. E. et al.: Substorm ion injections: Geosynchronous
  observations and test particle orbits in three-dimensional dynamic MHD fields, J. Geophys.
  Res., 102, 2325–2341, 1997.
- Bouhram, M., Klecker, B., Paschmann, G., Haaland, S., Hasegawa, H., Blagau, A., et al. (2005).
  Survey of energetic O<sup>+</sup> ions near the dayside mid-latitude magnetopause with Cluster. Annales
  Geophysicae, 23(4), 1281–1294. doi:10.5194/angeo-23-1281-2005.

486	Burch, J. L., Moore, T. E., Torbert, R. B., and Giles, B. L. (2016). Magnetospheric Multiscale
487	Overview and Science Objectives. Space Science Reviews, 199(1–4), 5–21.
488	doi:10.1007/s11214-015-0164-9.
489	Daglis, I. A., N. P. Paschalidis, E. T. Sarris, W. I. Axford, G. Kremser, B. Wilken, and G. Gloeckler
490	(1991), Statistical features of the substorm expansion phase as observed by AMPTE/CCE
491	spacecraft, in Magnetospheric Substorms, Geophys. Monogr. Ser, vol. 64, edited by J. R. Kan
492	et al., pp. 323–332, AGU, Washington, D. C.
493 494 495	<ul> <li>Daglis, I. A., Livi, S., Sarris, E. T., and Wilken, B. (1994), Energy density of ionospheric and solar wind origin ions in the near - Earth magnetotail during substorms, J. Geophys.</li> <li>Res., 99(A4), 5691–5703, doi:10.1029/93JA02772.</li> </ul>
496	Daglis, I. A., <u>&amp; and Axford, W. I. (1996)</u> . Fast ionospheric response to enhanced activity in geospace:
497	Ion feeding of the inner magnetotail. Journal of Geophysical Research, 101(A3), 5047–5065,
498	doi:10.1029/95JA02592.
499 500	Daglis, I. A., R. M. Thorne, W. Baumjohann, and S. Orsini (1999), The terrestrial ring current: Origin, formation, and decay, Rev. Geophys., 37, 407, doi:10.1029/1999RG900009.
501	Dai, L., J. R. Wygant, C. A. Cattell, S. Thaller, K. Kersten, A. Breneman, X. Tang, R. H. Friedel, S. G.
502	Claudepierre, and X. Tao (2014), Evidence for injection of relativistic electrons into the
503	Earth's outer radiation belt via intense substorm electric fields, Geophys. Res. Lett., 41, 1133–
504	1141,doi:10.1002/2014GL059228.
505	Dai, L., Wang, C., Duan, S., He, Z., Wygant, J. R., Cattell, C. A.,Tang, X. (2015). Near-Earth
506	injection of MeV electrons associated with intense dipolarization electric fields: Van Allen
507	Probes observations. Geophysical Research Letters, 42, 6170–6179. https://doi.org/10.1002/
508	2015GL064955.
509	Duan, S. P., Liu, Z. X., Liang, J., Zhang, Y. C., & and Chen, T. (2011). Multiple magnetic
510	dipolarizations observed by THEMIS during a substorm. Annales de Geophysique, 29(2),
511	331–339, doi:10.5194/angeo- 29- 331- 2011.
512	Duan, S. P., L. Dai, C. Wang, J. Liang, A. T. Y. Lui, L. J. Chen, Z. H. He, Y. C. Zhang, and V.
513	Angelopoulos (2016), Evidence of kinetic Alfvén eigenmode in the near-Earth magnetotail
514	during substorm expansion phase, J. Geophys. Res. Space Physics, 121, 4316–4330,
515	doi:10.1002/2016JA022431.
516	Duan, S., Dai, L., Wang, C., He, Z., Cai, C., Zhang, Y. C., Khotyaintsev, Y. V. (2017). Oxygen ions
517	O <sup>+</sup> energized by kinetic Alfvén eigenmode during dipolarizations of intense substorms.
518	Journal of Geophysical Research: Space Physics, 122. https://doi.org/10.1002/ 2017JA024418
519 520 521	Ebihara, Y., M C. Fok, T. J. Immel, and P. C. Brandt (2011), Rapid decay of storm time ring current due to pitch angle scattering in curved field line, J. Geophys. Res., 116, A03218, doi:10.1029/2010JA016000.
522 523 524	Elliott, H. A., Comfort, R. H., Craven, P. D., Chandler, M. O., and Moore, T. E. (2001), Solar wind influence on the oxygen content of ion outflow in the high - altitude polar cap during solar minimum conditions, J. Geophys. Res., 106(A4), 6067–6084, doi:10.1029/2000JA003022.

525	Ergun, R. E., Tucker, S., Westfall, J., Goodrich, K. A., Malaspina, D. M., Summers, D., et al. (2016).
526	The Axial Double Probe and Fields Signal Processing for the MMS Mission. Space Science
527	Reviews, 199(1-4), 167-188. doi:10.1007/s11214-014-0115-x.
528	Fok, M., Moore, T. E., Brandt, P. C., Delcourt, D. C., Slinker, S. P., & and Fedder, J. A. (2006).
529	Impulsive enhancements of oxygen ions during substorms. Journal of Geophysical Research,
530	111, A10222, doi:10.1029/2006JA011839.
531	Frank, L. A., W. R. Paterson, and J. B. Sigwarth (2002), Observations of plasma injection into the ring-
532	current during substorm expansive phase, J. Geophys. Res., 107(A11), 1343, -
533	doi:10.1029/2001JA000169
534	Fuselier, S. A., Klumpar, D. M., Peterson, W. K., and Shelley, E. G. (1989), Direct injection of
535	ionospheric O <sup>+</sup> into the dayside low latitude boundary layer. Geophysical Research Letters,
536	<u>16(10)</u> , 1121–1124. doi.:10.1029/GL016i010p01121.
37 38 39	Fuselier, S. A., D. M. Klumpar, and E. G. Shelley (1991), Ion reflection and transmission during reconnection at the Earth's subsolar magnetopause, Geophys. Res. Lett., 18, 139–142, doi:10.1029/90GL02676.
40 41	Fuselier, S. A., et al. (2016 <u>a</u> ), Magnetospheric ion influence on magnetic reconnection at the duskside magnetopause, Geophys. Res. Lett., 43,1435–1442, doi:10.1002/2015GL067358.
42	Fuselier, S. A., Lewis, W. S., Schiff, C., Ergun, R., Burch, J. L., Petrinec, S. M., and Trattner, K. J.
43	(2016b). Magnetospheric Multiscale Science Mission Profile and Operations. Space Science
44	<u>Reviews</u> , 199(1–4), 77–103. doi:10.1007/s11214-014-0087-x.
15	Fuselier, S. A., Trattner, K. J., Petrinec, S. M., Denton, M. H., Toledo - Redondo, S., André, M., et al.
16	(2019).Mass loading the Earth's dayside magnetopause boundary layer and its effect on
17	magnetic reconnection. Geophysical Research Letters, 46, doi:10.1029/2019GL082384.
8 9 0	Ganushkina, N. Y., Pulkkinen, T. I., <u>&amp; and</u> Fritz, T. (2005). Role of substorm-associated impulsive electric fields in the ring current development during storms. Annales Geophysicae, 23(2), 579–591, doi:10.5194/angeo-23-579-2005.
51	Hsu, T S., and McPherron, R. L. (2003), Occurrence frequencies of IMF triggered and nontriggered-
52	substorms, J. Geophys. Res., 108, 1307, doi:10.1029/2002JA009442, A7.
53	Kim, K. C., Lee, D Y., Lee, E. S., Choi, C. R., Kim, K. H., Moon, Y. J., Cho, K. S., Park, Y. D.,
54	and Han, W. Y. (-2005), A new perspective on the role of the solar wind dynamic pressure in
55	the ring current particle loss through the magnetopause, J. Geophys. Res., 110, A09223,
56	doi:10.1029/2005JA011097.
57	<ul> <li>Kronberg, E. A., Haaland, S. E., Daly, P. W., Grigorenko, E. E., Kistler, L. M., Fränz, M., and</li></ul>
58	Dandouras, I. (2012), Oxygen and hydrogen ion abundance in the near-Earth magnetosphere:
59	Statistical results on the response to the geomagnetic and solar wind activity conditions, J.
60	Geophys. Res., 117, A12208, doi:10.1029/2012JA018071.
561	Keika, K., P. C. Brandt, M. Nosé, and D. G. Mitchell (2011), Evolution of ring current ion energy-
562	spectra during the storm recovery phase: Implication for dominant ion loss processes, J
563	Geophys. Res., 116, A00J20, doi:10.1029/2010JA015628

564	Keika, K., L. M. Kistler, and P. C. Brandt (2013), Energization of O+ ions in the Earth's inner-
565	magnetosphere and the effects on ring current buildup: A review of previous observations and-
566	possible mechanisms, J. Geophys. Res. Space Physics, 118, 4441–4464,-
567	doi:10.1002/igra.50371
568 569 570 571	Kronberg, E. A., Ashour-Abdalla, M., Dandouras, I., Delcourt, D. C., Grigorenko, E. E., Kistler, L. M.,Zelenyi, L. M. (2014). Circulation of heavy ions and their dynamical effects in the magnetosphere: Recent observations and models. Space Science Reviews, 184(1-4), 173–235, doi:10.1007/s11214-014-0104-0.
572	Kronberg, E. A., E. E. Grigorenko, S. E. Haaland, P. W. Daly, D. C. Delcourt, H. Luo, L. M. Kistler,
573	and I. Dandouras (2015), Distribution of energetic oxygen and hydrogen in the near-Earth
574	plasma sheet, J. Geophys. Res. Space Physics, 120, 3415–3431, 681
575	doi:10.1002/2014JA020882.
576 577	Lennartsson, W., <u>&amp;-and</u> Shelley, E. G. (1986). Survey of 0.1- to 16-keV/ e plasma sheet ion composition. Journal of Geophysical Research, 91(A3), 3061.
578 579	Li, X., M. Hudson, A. Chan, and I. Roth (1993), Loss of ring current O <sup>+</sup> ions due to interaction with Pc 5 waves, J. Geophys. Res., 98, 215–231, doi:10.1029/92JA01540.
580	Liao, J., Kistler, L. M., Mouikis, C. G., Klecker, B., Dandouras, I., and Zhang, J
581	<u>C. (2010), Statistical study of O<sup>+</sup> transport from the cusp to the lobes with Cluster CODIF</u>
582	<u>data, J. Geophys. Res., 115, A00J15, doi:10.1029/2010JA015613.</u>
583	Liu, Y. H., C. G. Mouikis, L. M. Kistler, S. Wang, V. Roytershteyn, and H. Karimabadi (2015), The
584	heavy ion diffusion region in magnetic reconnection in the Earth's magnetotail, J.Geophys.
585	Res., 120, 3535-3551, doi:10.1002/2015JA020982.
586	Lindqvist, PA., Olsson, G., Torbert, R. B., King, B., Granoff, M., Rau, D., et al. (2016). The Spin-
587	Plane Double Probe Electric Field Instrument for MMS. Space Science Reviews, 199(1–4).
588	137–165. doi:10.1007/s11214-014-0116-9.
589 590 591 592 593 594	<ul> <li>Lui, A. T. Y., Liou, K., Nosé, M., Ohtani, S., Williams, D. J., Mukai, T., et al. (1999). Near-Earth dipolarization: Evidence for a non-MHD process. Geophysical Research Letters, 26(19), 2905–2908. Doi:10.1029/1999GL003620.</li> <li>Luo, H., E. A. Kronberg, K. Nykyri, K. J. Trattner, P. W. Daly, G. X. Chen, A. M. Du, and Y. S. Ge (2017), IMF dependence of energetic oxygen and hydrogen ion distributions in the near-Earth magnetosphere, J. Geophys. Res. Space Physics, 122, 5168–5180,</li> </ul>
595	doi:10.1002/2016JA023471.
596	Moore, T. E., Arnoldy, R. L., Feynman, J., and Hardy, D. A.: Propagating substorm injection fronts, J
597	Geophys. Res., 86, 6713–6726, 1981, doi:10.1029/JA091iA03p03061.
598	Nosé, M., Lui, A. T. Y., Ohtani, S., Mauk, B. H., McEntire, R. W., Williams, D. J.,Yumoto, K.
599	(2000). Acceleration of oxygen ions of ionospheric origin in the near-Earth magnetotail during
600	substorms. Journal of Geophysical Research, 105(A4), 7669–7677,
601	doi:10.1029/1999JA000318.

602 Ohtani, S., M. Nosé, S. P. Christon, and A. T. Y. Lui (2011), Energetic O<sup>+</sup> and H<sup>+</sup> ions in the plasma 603 sheet: Implications for the transport of ionospheric ions, J. Geophys. Res., 116, A10211, 604 doi:10.1029/2011JA016532. 605 Ono, Y., M. Nosée, S. P. Christon, and A. T. Y. Lui (2009), The role of magnetic field fluctuations in nonadiabatic acceleration of ions during dipolarization, J. Geophys. Res., 114, A05209, 606 607 doi:10.1029/2008JA013918. 608 Phan, T. D., Dunlop, M. W., Paschmann, G. et al.: Cluster observations of continuous reconnection at 609 the magnetopause under steady interplanetary magnetic field conditions, Ann. Geophys., 22, 610 2355-2367, 2004, 611 Pollock, C., Moore, T., Jacques, A., Burch, J., Gliese, U., Saito, Y., et al. (2016). Fast Plasma 612 Investigation for Magnetospheric Multiscale. Space Science Reviews, 199(1–4), 331–406. 613 doi:10.1007/s11214-016-0245-4. 614 Russell, C. T., Anderson, B. J., Baumjohann, W., Bromund, K. R., Dearborn, D., Fischer, D., et al. 615 (2016). The Magnetospheric Multiscale Magnetometers. Space Science Reviews, 199(1-4), 616 189-256. doi:10.1007/s11214-014-0057-3. 617 Shue, J. - H., et al. (1998), Magnetopause location under extreme solar wind conditions, J. Geophys. 618 Res., 103(A8), 17691-17700, doi:10.1029/98JA01103. 619 Sonnerrup, B. U. Ö., G. Paschmann, I. Papamastorakis, N. Sckopke, G. Haerendel, S. J. Bame, J. R. 620 Asbridge, J. T. Gosling, and C. T. Russell (1981), Evidence for magnetic field reconnection at 621 the Earth's magnetopause, J. Geophys. Res., 86, 10,049–10,067, 622 doi:10.1029/JA086iA12p10049. 623 Slapak, R., H. Nilsson, L. G. Westerberg, and A. Eriksson (2012), Observations of oxygen ions in the 624 dayside magnetosheath associated with southward IMF, J. Geophys. Res., 117, A07218, 625 doi:10.1029/2012JA017754. 626 Slapak, R., Nilsson, H., Westerberg, L. G., and Larsson, R (2015), O<sup>+</sup> transport in the dayside 627 magnetosheath and its dependence on the IMF direction, Ann. Geophys., 33, 301-307, 628 doi:10.5194/angeo-33-301-2015. 629 Tang, B., and Wang, C. (2018). Large scale current systems developed from substorm onset: Global 630 MHD results. Science China Technological Sciences, 61(3), 389-396, doi:10.1007/s11431-631 017-9132-y. Wang, S., L. M. Kistler, C. G.Mouikis, Y. Liu, and K. J. Genestreti (2014), Hot magnetospheric O<sup>+</sup> and 632 633 cold ion behavior in magnetopause reconnection: Cluster observations, J. Geophys.Res., 119, 9601-9623, doi:10.1002/2014JA020402. 634 635 Welling, D. T., Jordanova, V. K., Zaharia, S. G., Glocer, A., and Toth, G. (2011), The effects of 636 dynamic ionospheric outflow on the ring current, J. Geophys. Res., 116, A00J19, doi:10.1029/2010JA015642. 637 638 Winglee, R. M., Lewis, W., and Lu, G. (2005), Mapping of the heavy ion outflows as seen by IMAGE and multifluid global modeling for the 17 April 2002 storm, J. Geophys. Res., 110, A12S24, 639 640 doi:10.1029/2004JA010909.

- 641 Winglee, R. M., and Harnett, E. (2011), Influence of heavy ionospheric ions on substorm onset, J.
  642 Geophys. Res., 116, A11212, doi:10.1029/2011JA016447.
- Yau, A. W., and M. André (1997), Sources of ion outflow in the high latitude ionosphere, Space Sci.
  Rev., 80, 1.
- Yau, A. W., A. Howarth, W. K. Peterson, and T. Abe (2012), Transport of thermal-energy ionospheric
  oxygen (O<sup>+</sup>) ions between the ionosphere and the plasma sheet and ring current at quiet times
  preceding magnetic storms, J. Geophys. Res., 117, A07215, doi:10.1029/2012JA017803.
- Yu, Y., and A. J. Ridley (2013), Exploring the influence of ionospheric O<sup>+</sup> outflow on magnetospheric dynamics: dependence on the source location, J. Geophys. Res. Space Physics, 118, 1711–
  1722, doi:10.1029/2012JA018411.
- Young, D. T., Burch, J. L., Gomez, R. G., De Los Santos, A., Miller, G. P., Wilson, P., et al. (2016).
   Hot Plasma Composition Analyzer for the Magnetospheric Multiscale Mission. Space Science
   Reviews, 199(1–4), 407–470. doi:10.1007/s11214-014-0119-6.
- Zeng, C., Duan, S., Wang, C., Dai, L., Fuselier, S., et al. (2019). Statistical study of oxygen ions
   abundance and spatial distribution in the dayside magnetopause boundary layer: MMS
   observations, J. Geophys. Res. Space Physics, under review.
- Zong, Q. G., Wilken, B., Fu, S. Y., Fritz, T. A., Korth, A., Hasebe, N., Williams, D. J., and Pu, Z. Y. (2001),Ring current oxygen ions escaping into the magnetosheath, J. Geophys.
  Res., 106(A11), 25541–25556, doi:10.1029/2000JA000127.

# 660 Figures and Captions

Figure 1. The three components of the IMF B<sub>x</sub>, B<sub>y</sub>, B<sub>z</sub> in the Geocentric Solar Magnetospheric
 coordinates, solar wind dynamic pressure, as well as AU, AL and AE index from CDAweb OMNI data.
 The two blue dashed lines indicate the interval of the magnetopause boundary layer crossing.

664 Figure 2. The energetic O<sup>+</sup> is observed at the magnetopause during an intense substorm on 03 October 665 2015 by MMS 4. From top to bottom are (a) the magnetic field three components, Bx (blue line), By 666 (gree line), Bz (red line) and the total magnitude Bt (black line), (b) the electric field three components, 667 Ex (blue), Ey (gree) and Ez (red), (c) ion parallel (red) and perpendicular (black) temperatures, as well 668 as electron parallel (blue) and perpendicular (green) temperatures, (d) The density of ion (green) and 669 electron (blue), (e) three components of the ion velocity, (f) the H<sup>+</sup> (over the full HPCA energy range 670 from 1 eV to 40 keV) and O<sup>+</sup> (at energies from 1 keV to 40 keV) densities, (g-h) electron and ion omnidirectional differential energy fluxes (keV/(cm<sup>2</sup> s sr KeV)<sup>-1</sup>), (i) to (l) present differential particle 671 672 fluxes (cm<sup>2</sup> s sr eV)<sup>-1</sup> of H<sup>+</sup>, O<sup>+</sup>, He<sup>+</sup>, He<sup>++</sup>, respectively. The Geocentric Solar Magnetospheric coordinate 673 system is adopted. The thick bars at the top of the panel present different regions encountered on this 674 magnetopause crossing event. The orange and blue bars represent the magnetosheath and the 675 magnetosphere, respectively. The green bar represents the magnetopause boundary layer. The black 676 <u>horizontal line in figure 2j is at 1 keV and the O<sup>+</sup> contamination from high H<sup>+</sup> fluxes is indicated by the</u>
677 red box. The FPI data in Figure 2c-e and g-h are from FPI L2 data products and in the fast mode.

**Figure 3.** Maps of 31 events of energetic  $O^+$  at the <u>dusk flankduskside</u> magnetopause during intense substorms with AE index larger than 500 nT in XY<sub>GSM</sub> plane. The O+ <u>number</u> density and the <u>number</u> density ratios of O<sup>+</sup>/H<sup>+</sup> are shown by the color signatures at the corresponding magnetopause location in Figures 3a and 3b, respectively. The blue curve line represents the nominal magnetopause. The diamond and circle represent the event at the magnetopause during the intense substorm expansion phase and recovery phase, respectively.

**Figure 4.** The relationship between the energetic O<sup>+</sup> at the <u>dusk flankduskside</u> magnetopause and AE index during intense substorms. From top to bottom, panels show\_<u>that</u> the O<sup>+</sup> and H<sup>+</sup> number density<u>ies</u>, (-Figure 4a), the <u>O<sup>+</sup>/H<sup>+</sup></u> number density ratio of O<sup>+</sup>/H<sup>+</sup>, (Figure 4b), and <u>O<sup>+</sup>/H<sup>+</sup></u> particle flux ration of O<sup>+</sup>/H<sup>+</sup>, (Figure 4c), respectively. Error bars indicate <u>9590</u>% confidence intervals.

Figure 5. The relationship between the energetic O<sup>+</sup> at the dusk flankduskside magnetopause and IMF
B<sub>y</sub> during intense substorms. The format is the same as that of Figure 4. From top to bottom, panels show
that the O+ and H+ number density, Figure 7a, the number density ratio of O+/H+, Figure 7b, and particle
flux ration of O+/H+ Figure 7c, respectively. Error barsindicate 95% confidence intervals.

Figure 6. The relationship between the energetic O+ at the dusk flankduskside magnetopause and IMF
 B<sub>z</sub> during intense substorms. The format is the same as that of Figure 4. From top to bottom, panels show
 that the O+ and H+ number density, Figure 7a, the number density ratio of O+/H+, Figure 7b, and particle
 flux ration of O+/H+ Figure 7c, respectively. Error barsindicate 95% confidence intervals.

Figure 7. The relationship between the energetic O+ at the dusk flankduskside magnetopause and solar
wind dynamic pressure during intense substorms. The format is the same as that of Figure 4. From top to
bottom, panels show that the O+ and H+ number density, Figure 7a, the number density ratio of O+/H+,
Figure 7b, and particle flux ration of O+/H+ Figure 7c, respectively. Error bars indicate 95% confidence
intervals.