



1 **Using the Galilean Relativity Principle to Understand the**  
2 **Physical Basis for Magnetosphere-Ionosphere Coupling**  
3 **Processes**

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9

10 **Abstract** We use the Principle of Galilean Relativity (PGR) to gain insight into the physical basis for  
11 magnetosphere-ionosphere coupling. The PGR states that the laws of physics are the same in all inertial  
12 reference frames, considering relative speeds between such reference frames that are significantly less than  
13 the speed of light. The PGR is a limiting case of the principle of Special Relativity, the latter applicable to  
14 any relative speeds between two inertial reference frames. Although the PGR has been invoked in past works  
15 related to magnetosphere-ionosphere coupling, it has not been fully exploited for the insights it can provide  
16 into such topics as large-scale ionospheric convection and high latitude heating. In addition, the difficulties  
17 of applying the PGR to electrodynamics has not been covered. The PGR can be used to show that in the high  
18 latitude ionosphere there often exists a reference frame where electric fields vanish at lower altitudes where  
19 collisions are important (altitudes near ~100-120 km). In this reference frame, it is problematic to assert that  
20 currents of magnetospheric origin cause horizontal electric fields in the ionosphere, as has been suggested  
21 for the causal origin of Subauroral Polarization Stream electric fields. Electric fields have also been invoked  
22 as the causal origin of large-scale ionospheric convection, which may be a problematic assertion in certain  
23 reference frames. The PGR reinforces the importance of the neutral species and ion-neutral collisions in  
24 magnetosphere-ionosphere coupling, which has been noted by several authors using detailed multi-species  
25 plasma calculations. A straightforward estimate shows that the momentum carried by electron field aligned  
26 currents of magnetospheric origin during disturbed periods is much less than the momentum changes  
27 experienced by the neutral species in an Earth-fixed frame. The primary driver of neutral species momentum  
28 changes during disturbed periods is the momentum imparted by the solar wind to ionospheric ions resulting  
29 from electrodynamic interactions. This is consistent with the idea that electric fields do not lead to large scale  
30 ionospheric convection.

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## 35 **1 Introduction**

36 Electrodynamic as it pertains to magnetosphere-ionosphere coupling is a critical aspect of the ionospheric  
37 response during geomagnetic storms. Large scale convection of the high-latitude ionospheric plasma (auroral  
38 latitudes and higher), and heating of the plasma and neutral species during disturbed conditions is a  
39 consequence of electric and magnetic forces that change dramatically when solar wind conditions lead to  
40 geomagnetic storms. Scientific consensus on fundamental aspects of the physical processes that occur at high  
41 latitude is not yet achieved, including the definition of Joule heating (Vasyliunas and Song, 2005;  
42 Verkhoglyadova et al. 2017). In this paper, we are able to gain insight into these physical processes by using  
43 the simple but powerful Principle of Galilean Relativity (PGR), which states that physical laws are invariant  
44 with respect to inertial reference frame.

45

46 Despite its deceptively simple expression, the PGR has important implications for high latitude  
47 electrodynamic. This arises due to the inertial reference frame-dependent property of the electric field in the  
48 high-latitude ionosphere. The large-scale electric field is directed predominantly in the horizontal direction  
49 and the magnetic field tends to be predominantly in the vertical direction. The PGR requires that electric  
50 fields in the high latitude ionosphere vary substantially according to inertial reference frame. This fact has  
51 been appreciated in the literature but not fully exploited for its physical implications. In particular, the concept  
52 that field-aligned currents cause electric fields that lead to high velocity plasma flow has been invoked for  
53 high latitude phenomena (Anderson et al., 1993). However, if another inertial reference frame is chosen,  
54 these same arguments would seem problematic because the electric field may vanish. The notion of a  
55 “preferred inertial reference frame” for high latitude electrodynamic is often cited (Vasyliūnas and Song,  
56 2005; Leake et al., 2014; Strangeway, 2012) and is useful when considering that high latitude phenomena  
57 occur within the physical media of plasma and neutral gases. However, a reference frame where the electric  
58 field vanishes is also useful to consider in understanding the physical basis of magnetosphere-ionosphere  
59 coupling.

60

61 In this paper, we discuss how the PGR affects high latitude electrodynamic. We discuss the literature on the  
62 topic of Galilean electrodynamic, which is the low-velocity limit of the theory of special relativity applied  
63 to electrodynamic. Galilean electrodynamic is used in the literature of magnetosphere-ionosphere coupling,  
64 but incompletely in the sense that transformation of the source terms – charges and currents – are not  
65 considered along with the field transformation equations. This leads to a contradictory set of equations  
66 whereby magnetic fields do not change between inertial reference frames, but currents do, even though  
67 currents are the source term for magnetic fields. We review how the PGR has been referred to in the literature  
68 and how it can be used to help interpret physical processes. We next discuss how the frame-variant nature of



69 the electric field can be used to interpret the physical basis of Ohm's law and high latitude electromagnetic  
70 energy conversion. Finally, in seeking physical explanations that are robust to choice of inertial reference  
71 frame, we are led to de-emphasize electric fields as a source of ion motion and electron currents at high  
72 latitudes. We discuss how the literature emphasizes the importance of relative velocities between different  
73 species populations within the plasma as the root cause of high latitude changes during disturbed periods.  
74 The question presents itself: what causes the ions to convect and acquire a different velocity than the neutral  
75 species? We address this question.

76

## 77 **2 The Principle of Galilean Relativity in Electrodynamics**

78

79 The Principle of Galilean Relativity (PGR) determines how physical quantities change between inertial  
80 reference frames traveling at speeds significantly less than the speed of light. These quantities include electric  
81 and magnetic fields, and their sources such as charges and currents. The PGR is a limiting case of the principle  
82 of Special Relativity, the latter being applicable for any relative speeds between inertial reference frames.  
83 Either form of relativity is an important symmetry of nature: physical laws do not depend on one's velocity  
84 relative to an "absolute" or preferred reference frame.

85

86 What makes the PGR a useful idea in the context of high latitude electrodynamics is that the electric field  
87 component perpendicular to  $\mathbf{B} - \mathbf{E}$  being largely horizontal at high latitudes – substantially changes with  
88 choice of inertial reference frame, even for relatively low velocities that are characteristic of high latitude  
89 processes (e.g. ~100s of m/s to a few km/s). The PGR can help to interpret physically the equations governing  
90 high latitude electrodynamics by considering these equations in different inertial reference frames.

91

92 In this section, we review the literature of how electric and magnetic fields transform in the low velocity limit  
93 that corresponds to the PGR. Throughout the text, we use the symbol  $\mathbf{v}_r$  to indicate the relative velocity  
94 between two inertial reference frames, for example frames  $\mathcal{A}$  and  $\mathcal{B}$ . We use primed variables to refer to  
95 physical quantities in the inertial reference frame  $\mathcal{B}$  moving with velocity  $\mathbf{v}_r$  relative to reference frame  $\mathcal{A}$ .

96

97 As first deduced by Einstein in 1905, electric and magnetic fields transform between inertial reference frames  
98 according to the following relationships (Rousseaux, 2014; Heras, 2010):

99

$$\mathbf{E}' = \gamma \left( \mathbf{E} - \frac{\gamma - 1}{\gamma} \frac{\mathbf{v}_r (\mathbf{v}_r \cdot \mathbf{E})}{v_r^2} + \mathbf{v}_r \times \mathbf{B} \right) \quad (1)$$

100



$$\mathbf{B}' = \gamma \left( \mathbf{B} - \frac{\gamma - 1}{\gamma} \frac{\mathbf{v}_r (\mathbf{v}_r \cdot \mathbf{B})}{v_r^2} - \frac{1}{c^2} \mathbf{v}_r \times \mathbf{E} \right) \quad (2)$$

101

102 where  $(\mathbf{E}', \mathbf{B}')$  are the electric and magnetic fields in the new inertial reference frame,  $(\mathbf{E}, \mathbf{B})$  are the fields in  
103 the original frame, and  $\gamma = 1/\sqrt{1 - v_r^2/c^2}$  with  $c$  being the speed of light.

104

105 It is clear from these equations that there is not a unique low-velocity limit applicable to  $(\mathbf{E}, \mathbf{B})$  because the  
106 field transformations do not depend on velocity exclusively, but also on the electric and magnetic fields  
107 themselves. Le Bellac and Lévy-Leblond (1973) discuss two low-velocity limiting cases, which have since  
108 spawned a literature on the topic of “Galilean electromagnetism”. These are the electric and magnetic limits,  
109 according to which field magnitude is dominant. The electric limit applies when  $E \gg cB$ , and the magnetic  
110 limit applies when  $cB \gg E$ . High latitude electrodynamics encompasses the magnetic limit, which  
111 corresponds to the fact that Earth’s magnetic field is relatively strong and plasmas are quasi-neutral: the  
112 sources of electric fields (charges) are generally neglected compared to the sources of magnetic fields  
113 (currents). Magnetic fields in a plasma are caused by currents arising in a quasi-neutral medium because  
114 positive and negative charges move in opposing directions.

115

116 In the magnetic limit ( $cB \gg E$ ) and assuming that terms of order  $v_r/c$  are small, the Lorentz transformation  
117 laws (Equations (1) and (2)) become (Le Bellac and Lévy-Leblond, 1973):

118

$$\mathbf{E}' = \mathbf{E} + \mathbf{v}_r \times \mathbf{B} \quad (3)$$

119

$$\mathbf{B}' = \mathbf{B} \quad (4)$$

120

121 which are familiar transformation rules in the context of space physics (e.g. Parks, 2007; Vasyliūnas and  
122 Song, 2005). We refer to these low-velocity limit equations as comprising a Galilean transformation, by  
123 analogy to the more general Lorentz transformation.

124

125 An alternative derivation of the Galilean transformation of electric and magnetic fields is possible by  
126 considering the Lorentz force law (Preti et al., 2009; Heras, 2010). The Lorentz force  $\mathbf{F}$  in inertial reference  
127 frame  $\mathcal{A}$  is given by:

$$\mathbf{F} = q(\mathbf{E} + \mathbf{v} \times \mathbf{B}) \quad (5)$$

128

129 where  $\mathbf{F}$  is the force on a charge  $q$  moving with velocity  $\mathbf{v}$  in frame  $\mathcal{A}$ , and where the electric field is  $\mathbf{E}$  and  
130 the magnetic flux density is  $\mathbf{B}$ . If applied to a charged particle of mass  $m$ , the Lorentz force will result in  
131 acceleration  $\mathbf{F}/m$ . This acceleration is independent of inertial reference frame. Therefore, if  $\mathbf{F}'$  represents the



132 force measured in an inertial reference frame  $\mathcal{B}$  moving with velocity  $\mathbf{v}_r$  with respect to the original frame  
133  $\mathcal{A}$ , we know that  $\mathbf{F}' = \mathbf{F}$ . In reference frame  $\mathcal{B}$  the particle velocity is  $\mathbf{v}' = \mathbf{v} - \mathbf{v}_r$ , and the particle's mass  
134 and charge are invariant with respect to inertial reference frame (Galilean approximation). The equality of  
135 forces between the two inertial reference frames requires that:

136

$$(\mathbf{E} + \mathbf{v} \times \mathbf{B}) = (\mathbf{E}' + (\mathbf{v} - \mathbf{v}_r) \times \mathbf{B}') \quad (6)$$

137

138 which is achieved if Equations (3) and (4) are used.

139

140 Equation 3 shows that only the component of the electric field parallel to  $\mathbf{B}$  is unchanged under a Galilean  
141 transformation (i.e.  $\mathbf{B}$  is Galilean invariant, GI), whereas the perpendicular electric field changes depending  
142 on the relative velocity of frame  $\mathcal{B}$ . The frame-variant nature of the electric field has significant implications  
143 within the context of high-latitude electrodynamics, as we discuss below. Table 1 summarizes how different  
144 physical quantities relevant to high latitude electrodynamics vary under a Galilean transformation in the  
145 magnetic limit.

146

147 The source terms of the fields, charge density  $\rho$  and current density  $\mathbf{J}$ , also transform according to the principle  
148 of special relativity. As shown by Rousseaux (2014), they transform as a four vector according to:

149

$$\mathbf{J}' = \mathbf{J} - \gamma \mathbf{v}_r \rho + (\gamma - 1) \frac{\mathbf{v}_r (\mathbf{v}_r \cdot \mathbf{J})}{v_r^2} \quad (7)$$

150

$$\rho' = \gamma \left( \rho - \frac{(\mathbf{v}_r \cdot \mathbf{J})}{c^2} \right) \quad (8)$$

151

152

153 In the magnetic limit, the transformation equations become (Le Bellac and Lévy-Leblond, 1973):

154

$$\rho' = \rho - \frac{(\mathbf{v}_r \cdot \mathbf{J})}{c^2} \quad (9)$$

155

$$\mathbf{J}' = \mathbf{J} \quad (10)$$

156

157 The magnetic limit is associated with the condition  $c\rho \ll J$ . The Galilean invariance of current expressed by  
158 Equation (10) is familiar in the context of ionospheric electrodynamics (e.g. Thayer and Semeter, 2004) and  
159 is intuitive when charge density is zero. (Currents in the presence of no charge density arise from oppositely  
160 charged particles moving in opposite directions). The invariance of currents is consistent with the Galilean



161 invariance of magnetic fields. We note that zero charge density in the original reference frame leads to a  
 162 small charge density in the moving frame according to Equation (9). The literature of ionospheric  
 163 electrodynamics admits of non-zero charge densities or “charge accumulation” (Figure 3 in Vasylunas,  
 164 2012) leading to “polarization electric fields” (Richmond and Thayer, 2000) that is inconsistent with the  
 165 transformation law Equation (4). In a moving frame the net charge becomes a current that must lead to  
 166 modification of the magnetic field. Thus, electric fields that arise due to charge accumulation are inconsistent  
 167 with the usual field transformation equations adopted in the space physics literature.

168

169 In Figure 1 we depict an idealized situation where electric and magnetic fields are at perpendicular angles to  
 170 each other, to represent approximately the high latitude ionosphere where the Earth’s magnetic field is close  
 171 to vertical and large-scale convection electric fields are predominantly horizontal. We will refer to this  
 172 geometry elsewhere in the text. For electric fields of magnitude  $\sim 50$  mV/m, which can occur at high latitude  
 173 in a reference frame rotating with the Earth, and high latitude magnetic field magnitudes of  $\sim 50,000$  nT (near  
 174 120 km altitude), the electric field will be nearly zero in a reference frame moving with a speed of  $\sim 1.6$  km/s  
 175 relative to the Earth. The necessary direction of such a moving frame is shown in Figure 1. For situations  
 176 relevant to high latitude electrodynamic, an electric field that is non-negligible observed from an Earth fixed  
 177 frame can be zero viewed from an inertial reference frame moving at speeds consistent with the PGR and the  
 178 transformation rules of Table 1. Although an Earth fixed frame is not inertial, the small acceleration in this  
 179 frame is typically ignored for the purposes of these estimates.

180

181 Table 1: Transformation of physical quantities under change of inertial reference frame in the Galilean  
 182 magnetic limit. The primed quantities are in the new reference frame moving at velocity  $\mathbf{v}_r$  relative to the  
 183 original frame. All quantities are assumed to be quasi-static or slowly varying. Galilean invariance refers to  
 184 quantities that are the same in all inertial reference frames.

Physical Quantity	Transformation	Comment
Electric field $\mathbf{E}$	$\mathbf{E}' = \mathbf{E} + \mathbf{v}_r \times \mathbf{B}$	The electric field in a direction parallel to $\mathbf{B}$ is invariant, as is the electric field in the absence of a magnetic field
Magnetic field $\mathbf{B}$	$\mathbf{B}' = \mathbf{B}$	
Current density $\mathbf{J}$	$\mathbf{J}' = \mathbf{J}$	Assumes no net charge. Not invariant if charges are present.
Charge density $\rho$	$\rho' = \rho - \frac{1}{c^2} \mathbf{v}_r \cdot \mathbf{J}$	Charge per unit volume.
Heat energy $\mathbf{Q}$	$\mathbf{Q}' = \mathbf{Q}$	Heat energy and temperature are Galilean invariant (GI)
Velocity $\mathbf{v}$	$\mathbf{v}' = \mathbf{v} - \mathbf{v}_r$	



Relative velocity between two species	Invariant	
Collisional forces	Invariant	Depends on relative velocities, which are invariants

185

186

187 There are several implications of the electric and magnetic field transformation rules in Table 1. First,  
188 Maxwell's equations are not invariant under these transformations, unless the displacement current term is  
189 neglected in Ampere's law (Le Bellac and Lévy-Leblond, 1973; Preti et al., 2009; Heras, 2010). Second,  
190 currents in a moving reference frame that arise from accumulated charges in the original frame do not  
191 generate magnetic fields (Le Bellac and Lévy-Leblond, 1973). Fortunately, the quasi-neutral plasmas  
192 characteristic of geospace do not lead to large errors because of the small magnitudes of such charges.

193

194 The PGR is mentioned in textbooks on electrodynamics such as Jackson (1975) and Pollock and Stump  
195 (2001), primarily to show how the PGR fails in the context of electromagnetism. A point is made, however,  
196 that the physical principle of relativity demonstrates that "**E** and **B** have no independent existence" (Jackson,  
197 1975), which is true for Galilean as well as Special relativity.

198

199 Galilean relativity has been invoked within plasma physics in the context of wave-particle interaction and  
200 Landau damping (Chen, 2016; Dawson, 1961). Electrons gaining or losing energy from a plasma wave  
201 depends on the relative velocity of the electrons and the wave's phase velocity, hence the possibility of the  
202 plasma waves to alter the velocity distribution function of the electrons. Although relative velocity is a  
203 Galilean invariant, and hence so is the damping, the usual mathematical description of Landau damping does  
204 not explicitly contain relative velocities, so it appears superficially that Landau damping might depend on  
205 reference frame. (This is similar to how it superficially appears that the Lorentz force depends on reference  
206 frame). The authors resolve this apparent paradox by showing how a reference frame change affects both the  
207 electron distribution function and the wave energy, thus preventing Landau damping from violating the PGR.

208

## 209 **2.1 The PGR in Geospace**

210 Discussions of the PGR in the context of geospace are highly diverse. On the one hand, widely-used textbooks  
211 that describe ionospheric electrodynamics (Kivelson and Russell, 1995; Kelley, 2009; Brekke, 2013) do not  
212 refer to the PGR. In the journal-based literature, several authors discuss inertial reference frames in the  
213 context of high latitude electrodynamics (examples are provided in the text below). The reference frame-  
214 dependent property of the electric field is mentioned on occasion, but not emphasized or exploited in many  
215 cases. To our knowledge, inconsistency between the full set of Maxwell's equations and the Galilean



216 transformation laws for electrodynamics has not been emphasized in the context of geospace. The  
217 transformation of source terms is generally not discussed.

218

219 It should be noted that a “preferred reference frame” is a useful construct in plasma physics because of the  
220 importance of material media that obey the laws of classical mechanics, such as the plasma and the neutral  
221 atmosphere (Vasyliūnas and Song, 2005; Leake et al., 2014). Song et al. (2001) derive different versions of  
222 Ohm’s law appropriate to different reference frames. The expressions for the conductivities depend on  
223 whether one is in the inertial reference frame of the plasma or of the neutral species. This is related to the  
224 discussion in Jackson (1975, Section 11.1) on why the wave equation for sound waves depends on velocity  
225 relative to the medium carrying the waves. For electromagnetic waves, of course, there is no such medium  
226 and no preferred reference frame.

227

228 An important construct to examine from the perspective of the PGR, and widely seen in the literature, is  
229 based on the following quantity:  $\mathbf{E} + \mathbf{V} \times \mathbf{B}$ , where  $\mathbf{E}$  is an electric field,  $\mathbf{V}$  is the velocity of a constituent of  
230 the material medium (ions, electrons, neutrals, etc.), and  $\mathbf{B}$  is the magnetic field. In several publications the  
231 quantity  $\mathbf{E} + \mathbf{V} \times \mathbf{B}$  is referred to as “the electric field in the reference frame of species X” where  $\mathbf{V}$  is the  
232 bulk velocity of that species (Vasyliūnas, 2012; Vasyliūnas and Song, 2005; Thayer and Semeter, 2004;  
233 Richmond, 1995; Leake et al., 2014; Hidekatsu, 2009; Strangeway, 2012). The symbol  $\mathbf{E}'$  is often used to  
234 denote this quantity, but in this work, we will use the convention  $\mathbf{E}^* = \mathbf{E} + \mathbf{V} \times \mathbf{B}$  (following Vasyliūnas  
235 and Song, 2005) to maintain use of the prime symbol to refer to transformations between inertial reference  
236 frames.

237

238 Using the PGR, it is immediately clear that  $\mathbf{E}^*$  is not an electric field because of its transformation properties.  
239  $\mathbf{E}^*$  is a Galilean invariant, whereas an electric field depends on inertial reference frame. For similar reasons,  
240 the term  $\mathbf{V}_n \times \mathbf{B}$  is not an electric field, although it has been referred to as a dynamo electric field, where  $\mathbf{V}_n$   
241 is the velocity of the neutral species (Richmond, 1995).

242

243 The GI property of  $\mathbf{E}^*$  holds no matter what the velocity  $\mathbf{V}$  refers to, whether that of the neutral species, or  
244 the ions, etc. The use of  $\mathbf{E}^*$  has encouraged the exploration of reference frames tied to material media (e.g.  
245 Leake et al., 2014), whereas for electrodynamic quantities there is no need for a preferred inertial reference  
246 frame. In many situations applicable to the high latitude ionosphere, *there exists a reference frame for which*  
247 *the electric field is zero* (Figure 1). This frame is not tied to any particular medium, but is instructive to  
248 consider. As we show in the discussion of Ohm’s law, currents can arise in the absence of electric fields for  
249 precisely the reason that currents depend on  $\mathbf{E}^*$  rather than  $\mathbf{E}$ . Considering this special inertial reference frame  
250 – where the electric field vanishes – provides insight into the physical basis for momentum and energy  
251 changes at high latitudes.

252





253

## 254 2.2 The PGR and Ohm's Law

255 Ohm's law is derived from the plasma force balance equations taking collisions into account (Song et al.,  
256 2001; Richmond, 1995). These equations can be written, for ions and electrons, as:

$$qN_e(\mathbf{E} + \mathbf{u}_i \times \mathbf{B}) = N_e m_i v_{in}(\mathbf{u}_i - \mathbf{u}_n) + N_e m_i v_{ie}(\mathbf{u}_i - \mathbf{u}_e) \quad (11)$$

257

$$-qN_e(\mathbf{E} + \mathbf{u}_e \times \mathbf{B}) = N_e m_e v_{en}(\mathbf{u}_e - \mathbf{u}_n) - N_e m_e v_{ei}(\mathbf{u}_i - \mathbf{u}_e) \quad (12)$$

258

259 Charge neutrality is assumed such that  $N_e$  is the charge density of either electrons or ions.  $q$  is the elementary  
260 charge,  $\mathbf{u}_e$ ,  $\mathbf{u}_n$  and  $\mathbf{u}_i$  are the electron, neutral and ion velocities, respectively, and  $m_i$  and  $m_e$  are the ion and  
261 electron masses, respectively. We assume a single ion species for simplicity.  $v_{in}$ ,  $v_{ie}$ ,  $v_{en}$ , and  $v_{ei}$  are the ion-  
262 neutral, ion-electron, electron-neutral and electron-ion momentum transfer rates, respectively. These rates  
263 are sometimes referred to as "collision rates". The following reciprocity relation applies to these rates:  
264  $m_k v_{kl} = m_l v_{lk}$ , where  $k, l$  represent one of the species: electron, ion, or neutral (see Gombosi, 2004), and  
265  $k \neq l$ . We are ignoring pressure forces and other forces such as gravity, centrifugal, etc. Since these are force  
266 balance equations, the acceleration of each species is zero. It is clear these equations are consistent with PGR  
267 and are valid in any inertial reference frame, as can be readily seen by applying the transformations from  
268 Table 1. A simplification that collision frequencies do not depend on relative velocities is assumed; see  
269 Richmond (1995) for a statement regarding this limitation.

270

271 A form of Ohm's law that is derived from the force balance equations is as follows (see derivation provided  
272 by Song et al., 2001):

$$\mathbf{J} = \sigma_{\parallel} \mathbf{E}_{\parallel} + \sigma_p (\mathbf{E}_{\perp} + \mathbf{V} \times \mathbf{B}) + \sigma_H \mathbf{B} \times (\mathbf{E} + \mathbf{V} \times \mathbf{B}) / B \quad (13)$$

273

274 where the current density  $\mathbf{J}$  is given by  $qN_e(\mathbf{u}_i - \mathbf{u}_e)$  (charge neutrality is assumed).  $\sigma_{\parallel}$ ,  $\sigma_p$  and  $\sigma_H$  are the  
275 parallel, Pedersen and Hall conductivities, respectively.  $\mathbf{V}$  might be the bulk plasma velocity or the neutral  
276 wind velocity (see Song et al., 2001, Equation (17)). In our Equation (13), we have corrected typographical  
277 errors of Song et al.'s (2001) Equation 17). The conductivities typically involve terms that are GI, such as  
278 collision frequencies and gyrofrequencies that depend on the invariant magnetic field. Since conductivities  
279 are GI, this form of Ohm's law is explicitly GI, since both the left-hand side and right-hand side do not  
280 change under the Galilean transformation. In fact, any velocity  $\mathbf{V}$  ensures that the right hand side of Eq. (13)  
281 is GI.

282

283 Ohm's law Equation (13) states that that currents can be generated by 1) electric fields alone, 2) neutral winds  
284 alone, or 3) a combination of the two. The choice of inertial reference frame influences these three different  
285 possibilities. The causal connection between electric fields and currents is reference frame dependent in the



286 collisional ionospheric plasma where Ohm's law applies. If an inertial reference frame is chosen such that  
287 the perpendicular electric field is zero, then in that reference frame the currents are not generated by electric  
288 fields. Strangeway (2012) also questions whether electric fields cause the currents, instead suggesting that  
289 both the electric field and the currents are a consequence of the plasma flow. From the perspective of the  
290 PGR, it must be concluded that it is inconsistent to assert that electric fields cause plasma flow, at least in  
291 certain reference frames. It is consistent with the PGR to suggest that flow differences between plasma and  
292 neutrals are responsible for currents.

293  
294 Further clarification is found in Vasyliunas (2012) work on the "physical basis for ionospheric  
295 electrodynamics". Vasyliunas (2012) suggests that the ionospheric current is primarily a stress-balance  
296 current ultimately due to the relative motion between plasma and neutrals. That paper thus makes the direct  
297 link between two GI quantities, currents and relative velocities, without the problematic intermediary of the  
298 reference frame-dependent electric field. In the discussion section, we remark on the implications of requiring  
299 a direct relationship between currents and electric fields, in the context of Subauroral Polarization Stream  
300 (SAPS) electric fields.

301  
302 The PGR is also relevant to Figures (3) and (4) in Vasyliunas (2012) which diagram two differing physical  
303 interpretations of the neutral wind dynamo. We note that the interpretation on the left, the "conventional  
304 approach", requires the creation of polarization charges, which is inconsistent with the Galilean invariance  
305 of the magnetic field (Equation (4)). The charges generate currents in a moving reference frame, thus  
306 modifying the magnetic field in the moving frame.

### 307 **2.3 The PGR and Poynting's Theorem**

308 Poynting's theorem (PT) and Poynting flux are often used in the context of high latitude electrodynamic to  
309 understand energy deposition (Kelley, 1989; Thayer and Semeter, 2004). In this section, we discuss how  
310 Poynting's theorem is incompatible with Galilean electromagnetism in the magnetic limit.

311  
312 The *conservation of total energy* must hold in all inertial reference frames, that is, total energy does not  
313 change with time within a given reference frame. However, the different contributors to total energy, and the  
314 value of the total energy, can vary between inertial reference frames. For example, the kinetic energy of a  
315 particle will depend on its velocity, which varies with inertial reference frame. In contrast, heat energy is  
316 invariant with respect to inertial reference frame because it depends on relative velocities between ions,  
317 electrons and neutrals in motion.

318

319 Poynting's theorem is:

$$\frac{\partial W}{\partial t} = -\nabla \cdot \mathbf{S} - \mathbf{J} \cdot \mathbf{E} \quad (14)$$

320



321 where  $W$  is the energy density of the electromagnetic field:

$$W = \frac{1}{2} \left( \epsilon_0 \mathbf{E} \cdot \mathbf{E} + \frac{1}{\mu_0} \mathbf{B} \cdot \mathbf{B} \right) \quad (15)$$

322

323 and  $\mathbf{S}$  is the Poynting vector:

$$\mathbf{S} = \mathbf{E} \times \frac{1}{\mu_0} \mathbf{B} \quad (16)$$

324

325 Physically, Equation (14) represents energy conservation. It states that the rate of change of  
326 electromagnetic energy density  $W$  within a volume equals the energy leaving that volume, via  
327 divergence of Poynting flux ( $\nabla \cdot \mathbf{S}$ ), plus the rate of work done by the electromagnetic field within the  
328 volume ( $\mathbf{J} \cdot \mathbf{E}$ ). In the steady state approximation Poynting's theorem has been used to analyze energy  
329 partitioning at high latitudes (e.g. Equation 5 in Thayer and Semeter, 2004). In this case,  $\frac{\partial W}{\partial t} = 0$  and  
330 Poynting's theorem becomes:

331

$$\nabla \cdot \mathbf{S} = -\mathbf{J} \cdot \mathbf{E} \quad (17)$$

332

333 Referring to Table 1, it is clear that none of the individual terms in this equation are GI. Of interest is  
334 whether the equality holds in all inertial reference under the assumptions of Galilean electrodynamics  
335 in the magnetic limit.

336

337 It is straightforward to derive how the terms in Equation (17) vary with inertial reference frame using the  
338 transformation equations in Table 1. The Poynting vector  $\mathbf{S}$  transforms as:

339

$$\mathbf{S}' = \mathbf{S} + \frac{1}{\mu_0} (\mathbf{v}_r \times \mathbf{B}) \times \mathbf{B} \quad (18)$$

340

341 Using the chain rule applied to vector fields, we find the divergence of  $\mathbf{S}$  transforms as:

342

$$\nabla \cdot \mathbf{S}' = \nabla \cdot \mathbf{S} + \frac{1}{\mu_0} (-\mathbf{v}_r \times \mathbf{B}) \cdot \nabla \times \mathbf{B} - \mathbf{B} \cdot (\mathbf{v}_r \cdot \nabla) \mathbf{B} \quad (19)$$

343

344 where we have used the algebraic rules for the divergence of a cross-product and the curl of a cross  
345 product, the fact that the divergence of  $\mathbf{B}$  is zero, and  $\mathbf{v}_r$  has zero spatial derivatives. Using Ampere's  
346 law ( $\nabla \times \mathbf{B} = \mu_0 \mathbf{J}$ ) in the steady-state approximation (neglecting displacement currents) we have:

347



$$\nabla \cdot \mathbf{S}' = \nabla \cdot \mathbf{S} + \mathbf{v}_r \cdot (\mathbf{J} \times \mathbf{B}) - \frac{1}{\mu_0} \mathbf{B} \cdot (\mathbf{v}_r \cdot \nabla) \mathbf{B} \quad (20)$$

348

349 The rate at which the electromagnetic field does work ( $\mathbf{J} \cdot \mathbf{E}$ ) transforms as:

350

$$(\mathbf{J} \cdot \mathbf{E})' = \mathbf{J} \cdot \mathbf{E} - \mathbf{v}_r \cdot (\mathbf{J} \times \mathbf{B}) \quad (21)$$

351

352 Combining Equations (17), (20) and (21) we have:

$$\nabla \cdot \mathbf{S}' = -(\mathbf{J} \cdot \mathbf{E})' - \frac{1}{\mu_0} \mathbf{B} \cdot (\mathbf{v}_r \cdot \nabla) \mathbf{B} \quad (22)$$

353

354 Equation (22) shows that Poynting's theorem does not hold in the moving inertial reference frame. The  
 355 practical consequences of using the widely-used Galilean transformation rules will depend on the specific  
 356 situations where the Poynting's theorem is applied.

357

#### 358 2.4 The PGR and Joule Heating

359

360 The PGR can provide physical insight into Joule heating in the context of high latitude electrodynamics. Heat  
 361 is a Galilean invariant, since it is due to the random motion of the species independent of their uniform  
 362 translational motion. Several authors have expressed the Joule heating term in the energy conservation  
 363 equation as (e.g. Cole, 1962; Thayer and Vickrey, 1992; Matsuo and Richmond, 2008):

$$\mathbf{JH} = \mathbf{J} \cdot (\mathbf{E} + \mathbf{V}_n \times \mathbf{B}) = \mathbf{J} \cdot \mathbf{E}^*, \quad (23)$$

364

365 where  $\mathbf{V}_n$  is the velocity of the neutral wind. (More precisely,  $\mathbf{V}_n$  is the mass-weighted velocity of the plasma,  
 366 which is dominated by the neutral species).

367

368 The discussion in the Appendix of Thayer and Semeter (2004) reaches the same conclusion as in Vasyliūnas  
 369 and Song (2005) that heating at high latitudes is due to friction between plasma and neutrals (see also Brekke  
 370 and Rino, 1978). Such frictional heating depends on the relative velocity between two species and is thus GI,  
 371 consistent with its generating heat, another GI quantity. Thayer and Semeter (2004) note that  $\mathbf{JH}$  has been  
 372 equated to the following quantities:  $\mathbf{j} \cdot \mathbf{E}^*$ ,  $\sigma_p E^{*2}$  and  $j^2 / \sigma_c$  where  $\sigma_p$  is the Pedersen conductivity and  $\sigma_c$  is  
 373 the Cowling conductivity. These expressions are independent of inertial reference frame because they depend  
 374 on conductivities that depend on collision frequencies or relative velocities (Leake et al., 2014; Song et al.,  
 375 2001). The current density  $j$  is assumed to exist in the absence of a net charge density, so  $j$  is also reference  
 376 frame independent.  $\mathbf{E}^*$  refers to the electric field in a particular inertial reference frame and thus is reference



377 frame independent. However, it should be noted that  $\mathbf{E}^*$  is not an electric field and so does not obey Maxwell's  
378 equations.

379

380 Strangeway (2012) concludes that Joule dissipation in the reference frame of the neutrals results in heating.  
381 Since heating is GI, this would tend to suggest that Joule dissipation must also be GI. However, Strangeway  
382 (2012) initially defines Joule dissipation as  $\mathbf{J} \cdot \mathbf{E}$ , which is not GI. We know that  $\mathbf{J} \cdot \mathbf{E}$  is not necessarily  
383 dissipative since it is the work done by the electromagnetic field, and can consist of dissipative heating but  
384 also mechanical work done (Matsuo and Richmond, 2008) and the latter is not considered a form of  
385 dissipation.

386

387 According to Equation (23), JH can occur in the absence of an electric field (i.e. in an inertial reference frame  
388 where the electric field is zero; see Figure 1). This is consistent with the physical interpretation of Vasyliūnas  
389 (2012) that the “ionospheric current is thus primarily a stress-balance current”, and “not an Ohmic current in  
390 the usual sense.” The reason the neutral wind velocity appears in Eq. (23) is based on momentum  
391 considerations, as revealed when derived using the full multi-species plasma equations (Brekke and Rino,  
392 1978; Vasyliūnas and Song, 2005; Thayer and Vickrey, 1992). When JH occurs in the absence of an electric  
393 field, the currents associated with the heating are generated by ion-neutral velocity differences and not by an  
394 electric field (Mannucci et al., 2018; Vasyliūnas 2012; Thayer and Semeter 2004).

## 395 Discussion

396 The PGR is relevant to the physical basis of high latitude electrodynamics. A standard interpretation of  
397 magnetosphere-ionosphere coupling is that magnetospheric currents lead to electric fields that in turn lead to  
398 plasma motion, ion-neutral velocity differences and finally heating (Milan et al., 2017; Cowley, 2000; Kan,  
399 1997). The electric field is not an invariant, so explanations linking currents to electric fields are problematic  
400 from the perspective of the PGR. In addition, Vasyliūnas (2001) has shown that electric fields do not cause  
401 bulk plasma motion.

402

403 Consistency with the PGR is achieved if ion-neutral velocity differences are viewed as the primary causative  
404 factor of heating (as per Thayer and Semeter (2004) or Vasyliūnas and Song (2005)) and currents (Vasyliūnas  
405 (2012)). In the introduction section of Vasyliūnas and Song (2005) is it stated: “by virtue of the plasma  
406 momentum equation” that the ion-neutral velocity difference is proportional to the current density  $\mathbf{J}$ . From a  
407 cause and effect perspective, ion-neutral velocity difference *causes* the existence of  $\mathbf{J}$ . As we show below  
408 from simple momentum considerations, it would not be possible for field-aligned magnetospheric currents  
409 to lead to electric fields that then cause the ions to flow.

410

411 The PGR is relevant to explanations for the high velocity ion flows known as Subauroral ion drifts (SAID),  
412 which are often considered as a consequence of Sub-auroral Polarization Stream (SAPS) electric fields. SAPS



413 are postulated to be the result of magnetospheric currents closing in a low-conductivity region of the  
414 ionosphere, thus leading to large electric fields (Anderson et al., 1993; Clausen et al., 2012). Currents as the  
415 driver for high velocity plasma flows in the lower ionosphere is problematic from the perspective of the PGR,  
416 because the flow velocity depends on inertial reference frame, but the currents do not. The Rice Convection  
417 Model (Wolf et al., 2007; Vasyliunas, 1970) and its variants could be interpreted as implying that currents of  
418 magnetospheric origin closing in the conducting ionosphere are a driver of high latitude convection via  
419 electric fields. From the perspective of the PGR, it is more satisfactory to suggest that large-scale convection  
420 is driven by the velocity differences between the solar wind and magnetospheric plasmas without requiring  
421 an electric field as intermediary in the causal chain.

422

423 One might well ask how the velocity of the neutral species (Equation (23)) is relevant when calculating  
424 heating that results from work done by the electromagnetic field, since the neutral species do not experience  
425 the electromagnetic force. A partial answer is that the neutral species carry nearly all of the momentum in  
426 the high latitude ionosphere, and the plasma equations are based in part on considerations of momentum  
427 conservation. The causal chain of forcings relevant to high latitude phenomena would then appear to be:

428

- 429 • Coupling between the solar wind and magnetosphere imparts momentum to the  
430 magnetospheric plasma at altitudes of a few Earth radii.
- 431 • Through flux conservation, this momentum is imparted to the magnetospheric plasma at  
432 progressively lower altitudes, eventually down into the ionosphere.
- 433 • The momentum imparted to the ionospheric plasma creates a velocity difference between  
434 the plasma and neutral species, leading to transfer of momentum to the neutral species via  
435 collisional processes. The ion-neutral velocity differences lead to the presence of currents  
436 and heating in the ionosphere (Vasyliūnas, 2012; Mannucci et al., 2018).

437

438 In this picture, partly alluded to in Vasyliunas and Song (2005), the currents are by-products of momentum  
439 transfer (as the plasma equations suggest) and are not fundamental to the causal chain that couples the  
440 ionosphere to the magnetosphere (see also Vasyliunas, 2012). Currents in the absence of plasma-neutral  
441 collisions do not significantly heat the ionosphere in this picture. The detailed multispecies plasma  
442 calculations in Vasyliunas and Song (2005) show why *the neutral wind velocity is associated with Joule*  
443 *heating in Eq. (23), despite the fact that neutral species are not coupled to electromagnetic forces* (see also  
444 the Appendix of Thayer and Semeter, 2004).

445

446 Momentum transfer between ions and neutral species appears central to understanding cause and effect  
447 relationships in magnetosphere-ionosphere coupling (e.g. using the multispecies plasma equations).  
448 Momentum transfer is independent of inertial reference frame, and so can be considered a primary focus for  
449 understanding high latitude processes. Conversely, in the high latitude ionosphere, the electric field depends



450 on the inertial reference frame, so electric fields as the primary cause of high latitude electrodynamics can be  
451 problematic. In particular, electric fields caused by large-scale field-aligned currents (FACs) closing in the  
452 ionosphere should be reconsidered as a primary driver of high latitude dynamics.

453

454 We present an argument, based on momentum conservation, of why electric fields resulting from FACs  
455 closing in the ionosphere are unlikely to *cause* large-scale momentum changes to the ions and neutrals. We  
456 estimate the momentum carried by large-scale high latitude FACs that close horizontally in the ionosphere.  
457 We use an Earth-fixed frame to estimate relative momentum contributions involved in MI-coupling (relative  
458 momentum contributions are independent of reference frame). We focus on the large-scale FACs that carry  
459 most of the current coupling to the magnetosphere, and acknowledge that FACs can occur at multiple scales.  
460 Joule heating that appears simultaneously with large-scale FACs appears to be the dominant factor in energy  
461 transfer from the magnetosphere to the ionosphere (Verkhoglyadova et al., 2017).

462

463 FACs originating in the magnetosphere are often carried by electrons due to their high mobility (Carlson et  
464 al., 1998; Sugino et al., 2002). For typical large-scale disturbance FACs of approximately  $\sim 2 \mu\text{A}/\text{m}^2$  (e.g.  
465 Iijima and Potemra, 1976), and assuming an Earth-fixed frame, electron velocities are in the range  $\sim 156 \text{ m/s}$ ,  
466 assuming typical densities of  $8 \times 10^{10} \text{ el}/\text{m}^3$  in the lower ionosphere (near 110 km). These velocities are  
467 low enough that relativistic effects are not needed, i.e. the PGR is valid. Larger charged particle densities,  
468 e.g. created by intense electron precipitation, imply lower electron velocities for the same current.

469

470 Since the electron current  $\mathbf{j}_e$  is given by  $\mathbf{j}_e = qN_e\mathbf{u}_e$  and electron momentum density  $\mathbf{p}_e$  per unit volume is  
471 given by  $\mathbf{p}_e = m_eN_e\mathbf{u}_e$ , we have that  $\mathbf{p}_e = m_e\mathbf{j}_e/q$ , where  $m_e$  is the electron mass. Thus, momentum density  
472 of the FACs is approximately given by:  $\mathbf{p}_e \sim 1.1 \times 10^{-17} \text{ kg}\cdot\text{m}\cdot\text{s}^{-1}\cdot\text{m}^{-3}$ . We compare this momentum density  
473 residing with electron currents to changes in momentum density of the neutral species  $\mathbf{p}_n$  that occur during  
474 disturbed periods. We assume that: 1) neutral densities are 100 times or more larger than ion densities as is  
475 typical in the lower ionosphere, 2) the neutral species are dominated by atomic oxygen and 3) velocity  
476 changes of the neutrals under disturbed conditions can be comparable to those of the electrons ( $\sim 156 \text{ m s}^{-1}$ )  
477 (e.g. Zhang and Shepherd, 2002). These assumptions apply to the lower ionosphere where collisions between  
478 ions and neutrals are important. Under these assumptions, we find that the change in neutral momentum due  
479 to disturbed conditions is  $\Delta\mathbf{p}_n \sim 3.3 \times 10^{-11} \text{ kg}\cdot\text{m}\cdot\text{s}^{-1}\cdot\text{m}^{-3}$ , which is several orders of magnitude larger than the  
480 momentum density carried by electronic field-aligned currents.

481

482 It appears that during disturbed conditions, the momentum change of the neutral species dominates by orders  
483 of magnitude the momentum carried by the disturbance-related currents. This would remain true even when  
484 the currents are carried by ions, since from the magnetosphere these currents would be carried by light ions  
485 (hydrogen) versus the dominant heavier neutrals (e.g. oxygen) at altitudes where collisions occur. Velocity  
486 changes for the neutral species can easily reach 100 m/s or larger at high latitudes during disturbed conditions



487 (Kosch et al., 2010). We conclude that momentum changes associated with neutrals dominate over  
488 momentum carried by FACs entering from the magnetosphere that “close” as horizontal currents in the  
489 ionosphere. The changes in neutral momentum are thus due to collisions with ions, not due to the field aligned  
490 currents.

#### 491 **Conclusion**

492 We have reviewed electromagnetism in the context of high latitude electrodynamics to show that commonly-  
493 used relationships that transform electric and magnetic fields between inertial frames correspond to a limiting  
494 case known as the “magnetic limit” of Galilean electromagnetism. This limit is used in the literature related  
495 to magnetosphere-ionosphere coupling (Equations (3) and (4)). The magnetic limit of special relativity  
496 applies when electric and magnetic field magnitudes are related by  $cB \gg E$  which is valid in the collisional  
497 region of the high latitude ionosphere (altitudes near 110 km).

498

499 The Principle of Galilean Relativity is used to gain insight into the physical basis for magnetosphere-  
500 ionosphere coupling. We have considered Ohm’s law from the perspective of the PGR and noted that Ohm’s  
501 law need not be a relationship between currents and electric fields. We suggest that physical insight is gained  
502 by considering an inertial reference frame that is not tied to a particular species of the material medium: the  
503 reference frame in which the large-scale electric field is zero. We have also considered the steady-state  
504 formulation of Poynting’s Theorem from the perspective of the PGR and shown that, in the magnetic limit  
505 of Galilean electromagnetism, Poynting’s Theorem does not necessarily hold in all inertial reference frames.  
506 Previous authors have noted that Galilean electromagnetism is not consistent with the full set of Maxwell’s  
507 equations: the displacement current term must be removed to achieve consistency.

508

509 We have used the PGR to gain insight into why Joule heating at high latitudes is due to friction between  
510 plasma and neutrals, as shown by detailed multi-species plasma calculations (Brekke and Rino, 1978,  
511 Vasyliunas and Song, 2005 and Thayer and Semeter, 2004). Despite the term “Joule heating” that is  
512 associated with experiments where currents generated by electric fields cause heating, high latitude heating  
513 cannot depend on electric fields, the latter being reference frame-dependent. Plasma motion relative to the  
514 neutrals is of course a consequence of electromagnetic forces, originating in relative motion between  
515 planetary plasma and solar wind plasma in the collisionless regime of the interplanetary medium.

516

517 The Galilean transformation rules for currents and electric fields are very different from each other. This  
518 suggests it is problematic to assert that currents of magnetospheric origin *cause* horizontal electric fields in  
519 the ionosphere. Such a causal relation would seem to be reference frame dependent. At E-layer altitudes in  
520 the ionosphere, currents are independent of inertial reference frame, whereas perpendicular electric fields are  
521 not. Thus it is problematic to assert that the cause of high-velocity plasma flows known as SAID are  
522 horizontal currents closing in the ionosphere. During disturbed conditions generally, the momentum carried





523 by large-scale currents of magnetospheric origin is much less than the momentum changes of the neutral  
524 species occurring during disturbed conditions.

525

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533

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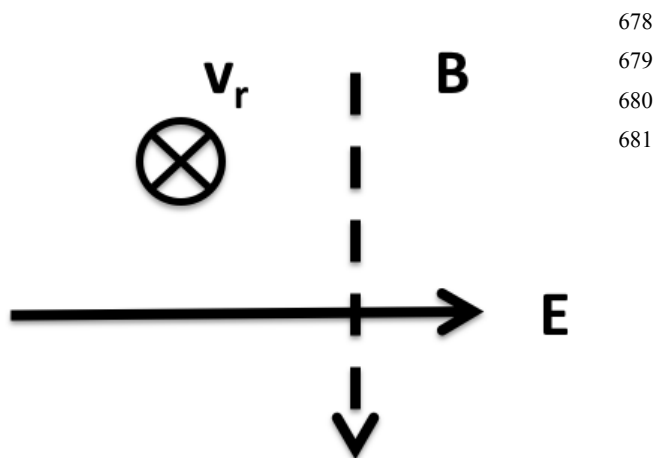
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677 **Figures**



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**Figure 1.** Schematic representation of electric and magnetic fields at northern high latitudes. Magnetic fields are directed vertically ( $B$ ). Electric fields are directed horizontally (e.g. in an Earth fixed frame, resulting from magnetospheric convection). In a reference frame moving with velocity  $v_r$  in the direction shown (into the page), the electric field is zero if the reference frame moves with speed  $\|E\|/\|B\|$ . For “typical” disturbed conditions, non-relativistic speeds of  $\sim 1.6$  km/s are sufficient.