

<sup>1</sup> **Transfer entropy and cumulant based cost as  
2 measures of nonlinear causal relationships in space  
3 plasmas: applications to  $D_{st}$**

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**9 Abstract.** It is well known that the magnetospheric response to the so-  
10 lar wind is nonlinear. Information theoretical tools such as mutual informa-  
11 tion, transfer entropy, and cumulant based analysis are able to characterize  
12 the nonlinearities in the system. Using cumulant based cost, we show that  
13 the nonlinear significance of  $D_{st}$  peaks at 3–12 hours lags that can be attributed  
14 to  $VBs$ , which also exhibit similar behavior. However, the nonlinear signif-  
15 icance that peaks at lags 25, 50, and 90 hours can be attributed to internal  
16 dynamics, which may be related to the relaxation of the ring current. These  
17 peaks are absent in the linear and nonlinear self-significance of  $VBs$ . Our  
18 analysis with mutual information and transfer entropy show that both meth-  
19 ods can establish that there are a strong correlation and transfer of infor-  
20 mation from  $V_{sw}$  to  $D_{st}$  at a time scale that is consistent with that obtained  
21 from the cumulant based analysis. However, mutual information also shows  
22 that there is a strong correlation in the backward direction, from  $D_{st}$  to  $V_{sw}$ ,  
23 which is counterintuitive. In contrast, transfer entropy shows that there is  
24 no or little transfer of information from  $D_{st}$  to  $V_{sw}$ , as expected because it  
25 is the solar wind that drives the magnetosphere, not the other way around.  
26 Our case study demonstrates that these information theoretical tools are quite  
27 useful for space physics studies because these tools can uncover nonlinear  
28 dynamics that cannot be seen with the traditional analyses and models that  
29 assume linear relationships.

## 1. Introduction

30 One of the most practically important concepts in dynamical systems is the notion of  
 31 causality. It is particularly useful to organize observational datasets according to causal  
 32 relationships in order to identify variables that drive the dynamics. Understanding causal  
 33 dependencies can also help to simplify descriptions of highly complex physical processes  
 34 because it constrains the coupling functions between the dynamical variables. Analysis  
 35 of those coupling functions can lead to simplification of the underlying physical processes  
 36 that are most important for driving the system. It is particularly useful from a practi-  
 37 cal standpoint to understand causal dependencies in systems involving natural hazards  
 38 because monitoring of causal variables is closely linked with warning.

A common method to establish causal dependencies in a data stream of two variables, e.g.,  $[a(t)]$  and  $[b(t)]$ , is to apply linear correlation studies such as *Strangeway et al.* [2005], which showed the relationship between downward Poynting flux and ion outflows. Causal relationships are typically identified by considering a time-shifted correlation function

$$\lambda_{ab}(\tau) \triangleq \frac{\langle a(t)b(t + \tau) \rangle - \langle a \rangle \langle b \rangle}{\sqrt{\langle a^2 \rangle - \langle a \rangle^2} \sqrt{\langle b^2 \rangle - \langle b \rangle^2}} \quad (1)$$

39 where  $\langle \dots \rangle$  is an ensemble average obtained by drawing samples at a set of measurement  
 40 times,  $\{t_0, t_1, \dots, t_N\}$ . For example, *[Borovsky et al., 1998]* used such a method to iden-  
 41 tify relationships between solar wind variables and plasma sheet variables. The causal  
 42 dependency that the plasma sheet responds to changes in the solar wind can be identified  
 43 from the time-shift of the peak of the cross correlation indicating a response time. From  
 44 this type of analysis it can be found that the plasma sheet generally responds from the

45 tail to the inner magnetosphere consistent with the notion of earthward convection. Such  
46 analysis has been particularly useful to help understand plasma sheet transport.

47 However, the procedure of detecting causal relationships based on linear cross-  
48 correlation suffers from a number of limitations. First it should be noted that the statisti-  
49 cal accuracy of the correlation function is limited by the resolution and length of the data  
50 stream. Second, the linear time series analysis ignores nonlinear correlations, which may  
51 be important for energy transfer in the magnetospheric system. For example, substorms  
52 are believed to involve storage and release of energy in the magnetotail, which is a highly  
53 nonlinear response. Similarly, magnetosphere-ionosphere coupling may also be highly non-  
54 linear involving the nonlinear development of accelerating potentials along auroral field  
55 lines and nonlinear current-voltage relationships. Third, the cross-correlation may not  
56 be a particularly clear measure when there are multiple peaks or if there is little or no  
57 asymmetry in the forward [i.e.,  $\lambda_{ab}(\tau)$ ] and backward directions [i.e.,  $\lambda_{ba}(\tau) = \lambda_{ab}(-\tau)$ ].  
58 Finally, the cross-correlation does not provide any way to clearly distinguish between two  
59 variables that are passively correlated because of a common driver rather than causally  
60 related.

61 In the remainder of this paper, we will discuss other methods to identify causal rela-  
62 tionships based on entropy based discriminating statistics such as mutual information and  
63 transfer entropy. We will also discuss the cumulant-based method. We will illustrate the  
64 shortcomings and strengths of the various methods for studying causality with examples  
65 from nonlinear dynamics and space physics.

## 2. Linear vs Nonlinear Dependency

66 It is well known that the magnetosphere responds to variation in the solar wind param-  
 67 eters [Clauer *et al.*, 1981; Baker *et al.*, 1983; Crooker and Gringauz, 1993; Papitashvili  
 68 *et al.*, 2000; Wing and Johnson, 2015; Johnson and Wing, 2015; Wing *et al.*, 2016], and  
 69 it has been established that the magnetosphere has a significant linear response to the  
 70 solar wind. However, it is also expected that the magnetosphere has a nonlinear response  
 71 [Tsurutani *et al.*, 1990; Vassiliadis *et al.*, 1990; Klimas *et al.*, 1998; Valdivia *et al.*, 2013;  
 72 Balikhin *et al.*, 2011]. The nonlinear response may be driven by internal dynamics rather  
 73 than driven externally [Wing *et al.*, 2005; Johnson and Wing, 2005]. For example, the  
 74 internal dynamics associated with loading and unloading of magnetic energy associated  
 75 with storms and substorms is nonlinear [e.g., Johnson and Wing, 2014, and references  
 76 therein]. Indeed, the data analysis of Bargatze *et al.* [1985] indicated that the dynamical  
 77 response of the magnetosphere to solar wind input could not be entirely understood using  
 78 linear prediction filters.

Suppose that we consider a set of variables  $\mathbf{a}$  and  $\mathbf{b}$  which could be vectors of variables measured in time and we would like to measure their dependency. Instead of considering the covariance matrix/correlation function, we consider a more general measure of dependency between an input and output is obtained by considering whether

$$P(\mathbf{a}, \mathbf{b}) \stackrel{?}{=} P(\mathbf{a})P(\mathbf{b}). \quad (2)$$

79 where  $P(\mathbf{a}, \mathbf{b})$  is the joint probability of input  $\mathbf{a}$  and output  $\mathbf{b}$  while  $P(\mathbf{a})$  and  $P(\mathbf{b})$  are  
 80 the probability of  $\mathbf{a}$  and  $\mathbf{b}$  respectively. If the relationship holds, then the variables  $\mathbf{a}$   
 81 and  $\mathbf{b}$  are independent. For all other cases, there is some measure of dependency. In the  
 82 case where the system output is completely known given the input,  $P(\mathbf{a}, \mathbf{b}) = P(\mathbf{a})$ . The

83 advantage of considering Equation 2 is that it is possible to detect the presence of higher  
 84 order nonlinear dependencies between the input and output even in the absence of linear  
 85 dependencies [Gershenfeld, 1998].

## 2.1. Mutual Information and Cumulant based cost

86 Mutual information and cumulant-based cost are two useful measures that quantify  
 87 Eq. 2. Mutual information has the advantage that in the limit of Gaussian joint proba-  
 88 bility distributions, it may be simply related to the correlation coefficient  $C_{ab}(\tau)$  defined  
 89 in equation 1 [Li, 1990]. Cumulants have the advantage of good statistics for limited  
 90 datasets and noisy systems [Deco and Schürmann, 2000]. Moreover, for high-dimensional  
 91 systems it is more efficient to compute moments of the data rather than try to construct  
 92 the probability density function.

Correlation studies also only detect linear correlations, so if the feedback involves non-  
 linear processes (highly likely in this case) then their usefulness may be seriously limited.  
 Alternatively, entropy-based measures such as mutual information [Prichard and Theiler,  
 1995; Materassi et al., 2011] and cumulants [Johnson and Wing, 2005] are useful for de-  
 tecting linear as well as nonlinear correlations. The mutual information is constructed  
 from the probability distribution function of the variables and may be computed using  
 an quantization procedure where data is binned such that the samples  $[a(t)]$  are assigned  
 discrete values  $\hat{a} \in \{a_1, a_2, \dots, a_n\}$  of an alphabet  $\aleph_1$  and  $[b(t)]$  is assigned discrete values  
 $\hat{b} \in \{b_1, b_2, \dots, b_m\}$  of an alphabet  $\aleph_2$ . The *ad hoc* time-shifted mutual entropy

$$\mathcal{M}_{ab}(\tau) \triangleq \sum_{\hat{a} \in \aleph_1, \hat{b} \in \aleph_2} p(\hat{a}(t + \tau), \hat{b}(t)) \log \left( \frac{p(\hat{a}(t + \tau), \hat{b}(t))}{p(\hat{a})p(\hat{b})} \right) \quad (3)$$

93 has been used as an indicator of causality, but suffers from the same problems as time-  
 94 shifted cross correlation when it has multiple peaks and long range correlations.

Similarly, examination of time-shifted cumulants could be used as an indicator of causal-  
 ity in a nonlinear system. In this case, we can define a discriminating statistic

$$D^C = \sum_{q=1}^{\infty} \sum_{i_1, \dots, i_q \in \Pi_q} K_{1i_2 \dots i_q}^2 \quad (4)$$

where

$$\begin{aligned} K_i &= C_i = \langle z_i \rangle & (4) \\ K_{ij} &= C_{ij} - C_i C_j = \langle z_i z_j \rangle - \langle z_i \rangle \langle z_j \rangle \\ K_{ijk} &= C_{ijk} - C_{ij} C_k - C_{jk} C_i - C_{ik} C_j + 2C_i C_j C_k \\ K_{ijkl} &= C_{ijkl} - C_{ijk} C_l - C_{ijl} C_k - C_{ilk} C_j - C_{ljk} C_i \\ &\quad - C_{ij} C_{kl} - C_{il} C_{kj} - C_{ik} C_{jl} + 2(C_{ij} C_k C_l \\ &\quad + C_{ik} C_j C_l + C_{il} C_j C_k + C_{jk} C_i C_l + C_{jl} C_i C_k \\ &\quad + C_{kl} C_i C_j) - 6C_i C_j C_k C_l \end{aligned}$$

are the cumulants

$$C_{i \dots j} = \int d\mathbf{z} P(\mathbf{z}) z_i \dots z_j \equiv \langle z_i \dots z_j \rangle \quad (5)$$

95 of the joint probability distribution for variables  $z_1, \dots, z_j$ .

With only two variables,  $a$  and  $b$ , defined above, we can consider the cost function

$$D_{a,b}^C(\tau) = D_C(a(t), b(t + \tau)) \quad (6)$$

96 The presence of nonlinear dependence has been identified by comparing the cumulant cost  
 97 for a time series with the cumulant based cost of surrogate time series, which are con-  
 98 structed to have the same linear correlations as in [Johnson and Wing, 2005]). Significance  
 99 measures the difference in the discriminating statistic from the mean of the discriminating  
 100 statistic of the surrogates in terms of the spread of the surrogates,  $\sigma$ .

101 In Section 3, we will show an application of cumulant based analysis to the distur-  
 102 bance storm-time index ( $D_{st}$ ). In principle, the cross-correlation, mutual information,  
 103 and cumulant-based cost should be independent of the selection of measurement points

<sup>104</sup> if the system is stationary; therefore, time stationarity can be examined by comparing  
<sup>105</sup> these discriminating statistics for groups of measurements drawn from different windows  
<sup>106</sup> of time as in [*Johnson and Wing, 2005; Wing et al., 2016*].

## 2.2. Transfer entropy

Another method for determining causality is the one-sided transfer entropy [*Schreiber, 2000; De Michelis et al., 2011; Materassi et al., 2014; Wing et al., 2016, 2018*], which is based upon the conditional mutual information

$$\mathcal{M}_C(x, y|z) \triangleq \sum_{x \in \mathbb{N}_1} \sum_{y \in \mathbb{N}_2} \sum_{z \in \mathbb{N}_3} p(x, y, z) \log \left( \frac{p(x, y, z)p(z)}{p(x, z)p(y, z)} \right) \quad (7)$$

<sup>107</sup> The conditional mutual information measures the dependence of two variables,  $x$  and  $y$ ,  
<sup>108</sup> given a conditioner variable,  $z$ . If either  $x$  or  $y$  are dependent on  $z$  the mutual information  
<sup>109</sup> between  $x$  and  $y$  is reduced, and this reduction of information provides a method to  
<sup>110</sup> eliminate coincidental dependence, or conversely to identify causal dependence.

Transfer entropy considers the conditional mutual information between two variables using the past history of one of the variables as the conditioner.

$$\mathcal{T}_{a \rightarrow b}(\tau) = \sum_{\hat{a} \in \mathbb{N}_1} \sum_{\hat{a}^{(k)} \in \mathbb{N}_1^{(k)}} \sum_{\hat{b} \in \mathbb{N}_2} p(\hat{a}(t + \tau), \hat{a}^{(k)}(t), \hat{b}(t)) \log \left( \frac{p(\hat{a}(t + \tau) | \hat{a}^{(k)}(t), \hat{b}(t))}{p(\hat{a}(t + \tau) | \hat{a}^{(k)}(t))} \right) \quad (8)$$

<sup>111</sup> where  $\hat{a}^{(k)}(t) = [\hat{a}(t), \hat{a}(t - \Delta), \dots, \hat{a}(t - (k - 1)\Delta)]$ . The standard definition of transfer  
<sup>112</sup> entropy takes  $k = 1$  (no lag), but keeping a higher embedding dimension could in prin-  
<sup>113</sup> ciple provide a more precise measure (for example, if  $a$  has periodicity a dimension of 2  
<sup>114</sup> may provide better prediction of future values of  $a$  from its past time series and therefore  
<sup>115</sup> lower the transfer entropy. Transfer entropy as a discriminating statistic has the following  
<sup>116</sup> advantages. First in the absence of information flow from  $a$  to  $b$  (i.e.,  $a(t + \tau)$  has no  
<sup>117</sup> additional dependence from  $b(t)$  beyond what is known from the past history of  $a^{(k)}(t)$ )

118  $p(\hat{a}(t+\tau)|\hat{a}^{(k)}(t), \hat{b}(t)) = p(\hat{a}(t+\tau|\hat{a}_{(k)}(t))$  and the transfer entropy vanishes. The transfer  
 119 entropy is also highly directional so that  $\mathcal{T}_{a \rightarrow b} \neq \mathcal{T}_{b \rightarrow a}$ . The advantage can be clearly  
 120 seen for dynamical systems where variables are forward differenced and the transfer en-  
 121 tropy is clearly one-sided while mutual information and correlation functions can even be  
 122 symmetric [Schreiber, 2000]. This measure also accounts for static internal correlations,  
 123 which can be used to determine whether two variables are driven by a common driver or  
 124 whether the variable  $b$  is causally driving the variable  $a$ .

125 Both mutual information and transfer entropy require binning of data. As mentioned  
 126 in Wing *et al.* [2016], the number of bins ( $n_b$ ) needs to be chosen properly and there are  
 127 some guidelines that can be followed. In general, we would like to maximize the amount  
 128 of information. Having too few bins would lump too many points into the same bin,  
 129 leading to loss of information. Conversely, having too many bins would leave many bins  
 130 with 0 or a few number of points, which also would lead to loss of information. Sturges  
 131 [1926] proposed that for a normal distribution, optimal  $n_b = \log_2(n) + 1$  and bin width  
 132  $w = \text{range}/n_b$ , where  $n$  = number of points in the dataset,  $\text{range}$  = maximum value –  
 133 minimum value of the points. In practice, there is usually a range of  $n_b$  that would work.

### 3. Application to space weather: $D_{st}$ analysis

134  $D_{st}$  (disturbance storm time index) is an hourly index that gives a measure of the  
 135 strength of the symmetric ring current that, in turn, provides a measure of the dynamics  
 136 of geomagnetic storms [Dessler and Parker, 1959]. Because of its global nature,  $D_{st}$  is  
 137 often used as one of the several indices that represent the state of the magnetosphere.  
 138 For example, Balasis *et al.* [2011] used the cumulative square amplitude of  $D_{st}$  time series  
 139 as a proxy for energy dissipation rate in the magnetosphere and found that it fits well

<sup>140</sup> a power law with log-periodic oscillations, which was interpreted as evidence for discrete  
<sup>141</sup> scale invariance in the  $D_{st}$  dynamics.

<sup>142</sup> When plasma sheet ions are injected into the Earth inner magnetosphere, they drift  
<sup>143</sup> westward around the Earth, forming the ring current. Studies have shown that the  
<sup>144</sup> substorm occurrence rate increases with solar wind velocity (high speed streams) [e.g.,  
<sup>145</sup> *Kisslinger et al.*, 2011; *Newell et al.*, 2016]. An increase in the solar wind electric field,  
<sup>146</sup>  $VB_z$ , can increase the dawn-dusk electric field in the magnetotail, which in turn deter-  
<sup>147</sup> mines the amount of plasma sheet particles that move to the inner magnetosphere [e.g.,  
<sup>148</sup> *Friedel et al.*, 2001]. Studies have shown that the electric field,  $VBs$  ( $V_{sw} \times$  southward  
<sup>149</sup> IMF  $B_z$ ) or  $VB_z$ , has a strong effect on the ring current dynamics [*Burton et al.*, 1975;  
<sup>150</sup> *O'Brien and McPherron*, 2000; *McPherron and O'Brien*, 2001; *Weygand and McPherron*,  
<sup>151</sup> 2006].

<sup>152</sup> For the present study, we examine the relationships between solar wind velocity ( $V_{sw}$ )  
<sup>153</sup> and  $VBs$  with  $D_{st}$ . We use  $D_{st}$  records in the period 1974 – 2001 obtained from  
<sup>154</sup> Kyoto University World Data Center for Geomagnetism (<http://swdcwww.kugi.kyoto-u.ac.jp/index.html>). The corresponding solar wind data are obtained from IMP-8, ACE,  
<sup>155</sup> WIND, ISEE1, and ISEE3 observations. The ACE SWEPAM and MAG data; and  
<sup>156</sup> the WIND MAG data are obtained from CDAWeb (<http://cdaweb.gsfc.nasa.gov/>). The  
<sup>157</sup> WIND 3DP data are obtained from the 3DP team directly. The ISEE1 and ISEE3  
<sup>158</sup> data are obtained from UCLA (these datasets are also available at NASA NSSDC  
<sup>159</sup> [<http://nssdc.gsfc.nasa.gov/space/>]). The IMP8 data come directly from the IMP teams.  
<sup>160</sup> The solar wind is propagated with minimum variance technique [*Weimer et al.*, 2003] to

<sub>162</sub> GSM (X, Y, Z) = (17, 0, 0)  $R_E$  to produce 1-min files, from which hourly averaged solar  
<sub>163</sub> wind parameters are constructed.

### 3.1. Cumulant based analysis

Section 2.1 presents the method of cumulant based cost. Here, we show an application of cumulant based cost to detect nonlinear dynamics in  $D_{st}$ . We consider the forward coupling between a solar wind variable such as  $VB_s$  and  $D_{st}$ , which characterizes the ring current response to the solar wind driver. We therefore consider the nonlinear cross-correlations of the vector

$$\mathbf{c}(t, \tau) = \{VB_s(t), D_{st}(t + \tau)\} = \{z_1, z_2\} \quad (9)$$

<sub>164</sub> The generalization of cost is based on realizations of  $\{z_1, z_2\}$ . In this case, each variable  
<sub>165</sub> is Gaussianized with unit variance to eliminate static nonlinearities (i.e. higher order  
<sub>166</sub> self-correlations in  $VB_s$  and  $D_{st}$  are eliminated so that the cost measures only cross-  
<sub>167</sub> dependence between  $VB_s$  and  $D_{st}$ ). This procedure is explained in the next paragraph.

<sub>168</sub> The distribution of  $D_{st}$  and  $VB_s$  are generally non-Gaussian. As such, the raw dis-  
<sub>169</sub> tributions (e.g., distribution of values of  $D_{st}$ ) may have nonzero higher-order cumulants  
<sub>170</sub> (e.g., they can have a skew and kurtosis). This property makes it more difficult to in-  
<sub>171</sub> terpret whether the higher order cumulants in the time evolution arise from the overall  
<sub>172</sub> shape of the distribution of data points or from the time-ordering of the data. To elim-  
<sub>173</sub> inate the inherent nonzero cumulants in the overall distribution of data, we construct a  
<sub>174</sub> rank-ordered map from the original dataset to a proxy dataset of the same length drawn  
<sub>175</sub> from a Gaussian distribution [Kennel and Isabelle, 1992; Schreiber and Schmitz, 1996;  
<sub>176</sub> Deco and Schürmann, 2000]. The distribution of the proxy dataset ensures that all cu-

mulants of the distribution beyond second order should in principle vanish. However, the time-ordering of the data can still lead to nonzero cumulants, because the joint probability distribution of  $D_{st}(t+\tau)$  and  $D_{st}(t)$  may be non-Gaussian even if the distribution of  $D_{st}$  is Gaussian. Moreover, it is simple to construct surrogate data from the Gaussianized data that shares the same autocorrelation by using the same power spectrum, but randomly shifting the phases of the Fourier coefficients. The surrogate data therefore has the same autocorrelation as the original data. Any deviation from the linear statistic is apparent from comparison with the surrogate data, and we interpret these deviations as evidence of nonlinear dependence because we have falsified the hypothesis that the data can be adequately described by linear statistics. This method has been successfully employed in *Johnson and Wing [2005]* where  $K_p$  record was analyzed with mutual information and cumulants.

In Figure 1 we plot the significance obtained from the year 1999 as a function of time delay,  $\tau$ . Significance extracted from  $\{VBs(t), D_{st}(t + \tau)\}$  and  $\{VBs(t), VBs(t + \tau)\}$  for 1999 are plotted in panels (a) and (b), respectively. It should be noted that there is a strong linear response at around 3 hour time delay. As shown in Figure 1a, there is a clear nonlinear response with peaking around 3–10, 25, 50 and 90 hours lasting for approximately 1 week. In contrast, in Figure 1b, the nonlinearity only has one broad peak around 3 – 12 hours in the self-significance for  $VBs$ , suggesting that the nonlinear and linear peaks at  $\tau = 3–12$  hours in in Figure 1a i may be associated with  $VBs$ . We will revisit the solar wind causal relationship with  $D_{st}$  using transfer entropy in Section 3.2.

The absence of the nonlinear peaks at  $\tau = 25, 50$ , and 90 hours in the self-significance for  $VBs$  (Figure 1b) suggest that these nonlinearities in  $\{VBs(t), D_{st}(t+\tau)\}$  are related to

internal magnetospheric dynamics. As the  $D_{st}$  index is thought to reflect storm activity, it is reasonable that nonlinear significance would decay on the order of 1 week as storms commonly last around that time. The strong nonlinear responses at  $\tau = 25, 50$ , and  $90$  hours are likely related to multiple modes of relaxation of the ring current following the commencement of storms. It should also be noted that other nonlinearities detected by even higher order cumulants may also be present; however, the calculation demonstrates the nonlinear nature of the underlying dynamics.

A common scenario for storm-ring current interaction is the following. A storm compresses the magnetosphere, intensifies the magnetic field in the magnetosphere, and injects energetic particles into the ring current region. The ring current intensifies during the main phase of the storm, which can last  $\sim 6$  hours [Weygand and McPherron, 2006]. Once the injection stops, the ring current begins to decay and the storm enters the recovery phase. Conservation of magnetic moment implies that anisotropies develop in the ring current and plasma sheet. Anisotropy drives the ring current plasma unstable to ion cyclotron waves. The ion cyclotron waves scatter energetic ions into the loss cone so that they are lost from the ring current. Nonlinear interaction between waves and particles keeps the plasma near marginal instability with a steady loss of energetic particles due to wave-particle scattering. Other loss mechanisms include charge exchange, coulomb scattering, and convective of ions to the front of the magnetopause. The ring current decay can have two stages [Kozyra et al., 2002]. In the first stage, the ring current decays rapidly and the loss mechanisms can be attributed to convective out flow, pitch-angle scattering in the ring current, and  $O^+$  charge exchange [e.g., Weygand and McPherron, 2006; Hamilton et al., 1988]. The second stage may typically begin about one day from

223 the commencement of the storm (see, for example, Figure 7 of *Kozyra et al.* [2002]). In  
224 the second stage, the decay rate is slower and is attributed mainly to  $H^+$  charge exchange  
225 [*Hamilton et al.*, 1988] and can take several days to deplete the ring current to the baseline  
226 level [*Smith et al.*, 1976]. We can speculate that the multiple nonlinear response lag times  
227 that are detected with the cumulant-based approach are likely the relaxation of the ring  
228 current due to complex interplay of multiple loss processes.

### 3.2. Transfer entropy

229 As mentioned in Section 2.2, transfer entropy gives a measure of how much information  
230 is transferred from one variable to another. We have applied transfer entropy and mutual  
231 information to the relationship between the  $V_{sw}$  and  $D_{st}$  for the period 1974 – 2001. The  
232 result is shown in Figure 2. Note that the mutual information measure suggests strong  
233 correlations between prior values of  $D_{st}$  and  $V_{sw}$ . This finding suggests that  $D_{st}$  could be  
234 a driver of  $V_{sw}$ , which is counterintuitive. On the other hand, the transfer entropy clearly  
235 shows that this information transfer in the backward direction ( $D_{st} \rightarrow V_{sw}$ ) does not rise  
236 above the noise level (the horizontal blue lines indicate mean and standard deviation of  
237 100 surrogate data sets where the data was randomly reordered.) This result is expected  
238 because it is the solar wind that drives the magnetosphere, not the other way around.  
239 The transfer of information from  $V_{sw}$  to  $D_{st}$  peaks at  $\tau = 8 – 11$  hours. The cumulant  
240 based analysis in Section 3.1 shows that the response of  $D_{st}$  to  $VBs$  has similar time scale.  
241 This time scale is consistent with the 4 to 15 hours transport time for the solar wind to  
242 reach the midnight and noon regions of the geosynchronous orbit, respectively, from the  
243 dayside magnetopause [*Borovsky et al.*, 1998]. The analysis presented here illustrates the  
244 power of the transfer entropy for accessing causality.

#### 4. Summary

245 We recently used mutual information, transfer entropy, and conditional mutual infor-  
 246 mation to discover the solar wind drivers of the outer radiation belt electrons [*Wing et al.*,  
 247 2016]. Because  $V_{sw}$  anticorrelates with solar wind density ( $n_{sw}$ ), it is hard to isolate the  
 248 effects of  $V_{sw}$  on radiation belt electrons, given  $n_{sw}$  and vice versa. However, using condi-  
 249 tional mutual information, we were able to determine the information transfer from  $n_{sw}$   
 250 or any other solar wind parameters to radiation belt electrons, given  $V_{sw}$  (or any other  
 251 solar wind parameters). We also showed that the triangle distribution in the radiation  
 252 belt electron vs. solar wind velocity plot [*Reeves et al.*, 2011] can be understood better  
 253 when we consider that  $V_{sw}$  and  $n_{sw}$  transfer information to radiation belt electrons with  
 254 2 days and 0 day (< 24 hr) lags, respectively. Also recently, we used transfer entropy to  
 255 better understand the causal parameters in the solar cycle and their response lag times  
 256 [*Wing et al.*, 2018].

257 As a follow up to *Wing et al.* [2016, 2018], the present study demonstrates further how  
 258 information theoretical tools can be useful for space physics and space weather studies.  
 259 Cumulant based analysis can be used to distinguish internal vs. external driving of the  
 260 system. Both mutual information and transfer entropy give a measure of shared infor-  
 261 mation between two variables (or vectors). However, unlike mutual information, transfer  
 262 entropy is highly directional. To illustrate, we apply mutual information, transfer entropy,  
 263 and cumulant based analysis to investigate the dynamics of  $D_{st}$  index.

264 Our analysis with mutual information and transfer entropy indicates that there are  
 265 strong linear and nonlinear correlations and transfer of information, respectively, in the  
 266 forward direction between  $V_{sw}$  and  $D_{st}$  ( $V_{sw} \rightarrow D_{st}$ ). However, mutual information indi-

267 cates that there is also a strong correlation in the backward direction ( $D_{st} \rightarrow V_{sw}$ ), which  
268 is puzzling and counterintuitive. In contrast, the transfer entropy indicates that there is  
269 no information transfer in the backward direction ( $D_{st} \rightarrow V_{sw}$ ), as expected because it is  
270 the solar wind that drives the magnetosphere, not the other way around. The transfer of  
271 information from  $V_{sw}$  to  $D_{st}$  peaks at  $\tau = 8 - 11$  hours.

272 Using the cumulant-based significance, we have established that the underlying dynam-  
273 ics of  $D_{st}$  is in general nonlinear exhibiting a quasiperiodicity which is detectable only if  
274 nonlinear correlations are taken into account. The strong nonlinear responses of  $D_{st}$  to  
275  $VBs$  at  $\tau = 25, 50$ , and  $90$  hours are likely related to multiple modes of relaxation of the  
276 ring current from multiple loss mechanisms following the commencement of storms. It is,  
277 of course, possible that these nonlinearities are caused by solar wind drivers other than  
278  $VBs$ . However, the timing of these nonlinearities would put them well in the recovery  
279 phase of a storm and previous studies suggested that the ring current decays in the recov-  
280 ery phase are strongly influenced by  $VBs$  [Burton *et al.*, 1975; O'Brien and McPherron,  
281 2000; McPherron and O'Brien, 2001]. The nonlinearities at  $\tau = 3 - 12$  hours are not  
282 caused by internal dynamics but rather by the solar wind driver, which is similar with  
283 the time scale for the solar wind transport time from the dayside magnetopause to the  
284 inner magnetosphere. This time scale is consistent with the time scale for the information  
285 transfer from the solar wind to  $D_{st}$  obtained from transfer entropy analysis.

286 Although linear models are useful, our results indicate that these models have to be  
287 used with cautions because solar wind – magnetosphere system is inherently nonlinear.  
288 Hence, nonlinearities generally need to be taken into account in order to describe the  
289 system accurately. Local-linear models (which include slow evolution of parameters) may

290 be able to handle some nonlinearities, but it is expected that these local-linear models  
 291 would have difficulties if the dynamics suddenly and rapidly change.

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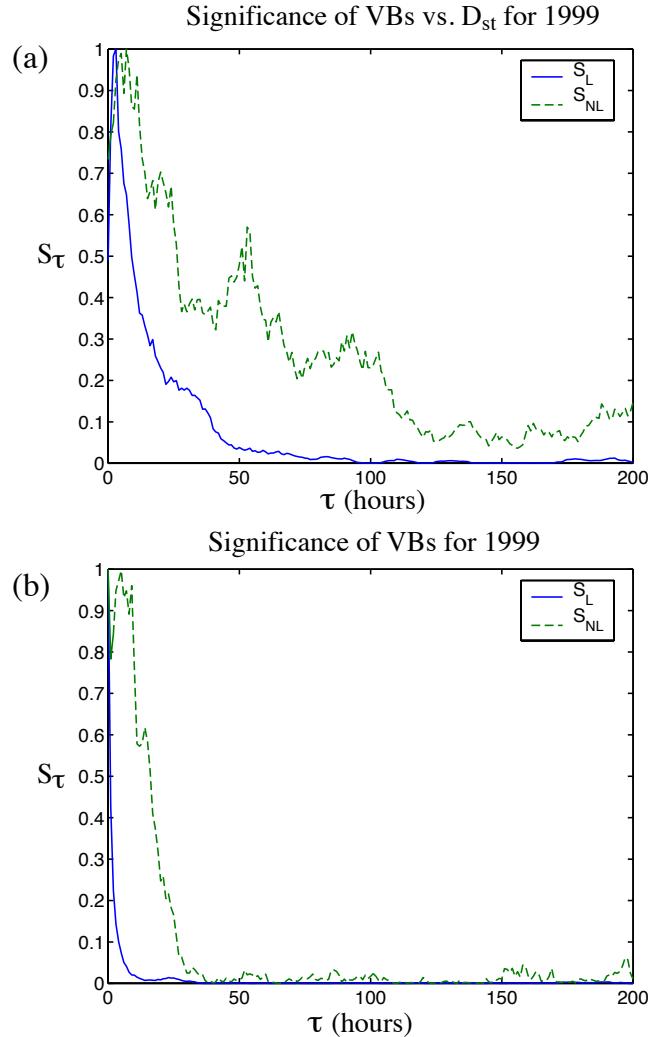
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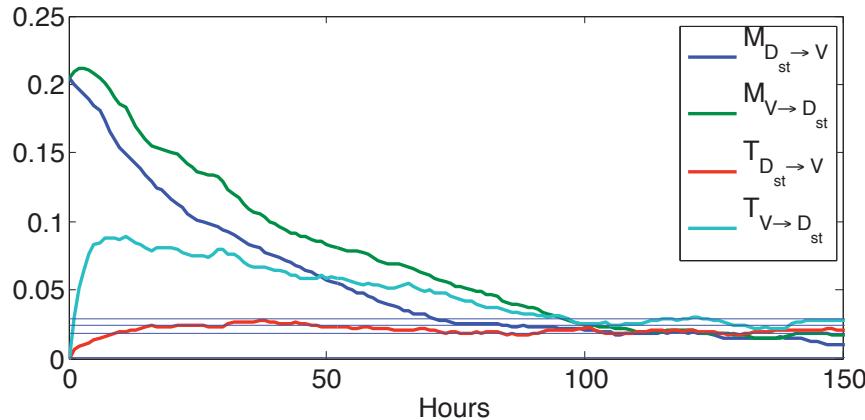
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**Figure 1.** Significance extracted from (a)  $\{VBs(t), D_{st}(t-\tau)\}$  and (b)  $\{VBs(t), VBs(t-\tau)\}$  for 1999. It should be noted that there is a strong linear response at around 3 hour time delay. There is a clear nonlinear response with a strong peak around 50 hours lasting for approximately 1 week. The longterm nonlinear response is absent in the solar wind data indicating that the longterm nonlinear correlations between  $VBs$  and  $D_{st}$  are the result of internal magnetospheric dynamics.



**Figure 2.** Comparison of mutual information and transfer entropy measures to determine causal driving of the magnetosphere as characterized by  $D_{st}$ . Note that causal driving appears to peak somewhat later (11 hours) than indicated by mutual information (2 hours) indicating that internal dynamics likely are very important initially. The backward transfer entropy is below the noise level for all values indicating that  $D_{st}$  in no way influences the upstream solar wind velocity. Such a conclusion could not be inferred from the mutual information measure.