



1 Statistical study of ULF waves in the magnetotail by THEMIS 2 observations

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18
19 **Abstract.** Ultra-low frequency (ULF) waves are ubiquitous in the magnetosphere. Previous studies
20 mostly focused on ULF waves in the dayside or near-earth region (with radial distance $R < 12 R_E$). In this
21 study, using the data of Time History of Events and Macroscale Interactions during Substorms (THEMIS)
22 during the period from 2008 to 2015, the Pc5-6 ULF waves in the tail region with $X^*_{GSM} < 0$, $8 R_E < R <$
23 $32 R_E$ (mostly on the stretched magnetic field lines) are studied statistically. A total of 1089 azimuthal
24 oscillating events and 566 radial oscillating events were found. The statistical results show that both the
25 azimuthal and radial oscillating events in the magnetotail region ($12 R_E < R < 32 R_E$) are more frequently
26 observed in the post-midnight region. The frequency decreases with increasing radial distance from Earth
27 for both azimuthal oscillating events ($8 R_E < R < 16 R_E$) and radial oscillating events ($8 R_E < R < 14 R_E$),
28 which is consistent with the field line resonances theory. About 52 % of events (including the azimuthal
29 and radial oscillating events) are standing waves in the region of 8-16 R_E , while only 2 % are standing
30 waves in the region of 16-32 R_E . There is no obvious dawn-dusk asymmetry of ULF wave frequency in
31 $8 R_E < R < 32 R_E$, which contrasts with the obvious dawn-dusk asymmetry found by previous studies in
32 the inner magnetosphere ($4 R_E < R < 9 R_E$). An examination for possible statistical relationships between
33 ULF wave parameters and substorm occurrences is carried out. We find that the wave frequency is higher
34 after the substorm onset than before it, and the frequency differences are more obvious in the midnight
35 region than in the flank region.



36

37 **Keyword.** Magnetospheric physics (Magnetotail; MHD waves and instabilities; Solar wind-
38 magnetosphere interactions)

39

40 **1 Introduction**

41

42 Ultra-low frequency (ULF) waves play a significant role in storing and transferring energy in the Earth's
43 magnetosphere. ULF waves can transport energy from the magnetosphere to the ionosphere, accelerate
44 energetic particles, modulate luminosity of aurorae, mediate reconnection and trigger substorm onset (e.g.,
45 Baumjohann and Glassmeier, 1984; Lessard et al., 1999; Ukhorskiy et al., 2005; Keiling, 2009; Rae et al.,
46 2014; Zong et al., 2017).

47 There are several excitation sources for magnetospheric ULF waves. These sources include the
48 Kelvin-Helmholtz instability (KHI) along the magnetopause (e.g., Walker, 1981; Claudepierre et al.,
49 2008), solar wind dynamic pressure impulse (e.g., Allan et al., 1986; Lee et al., 1989; Shi et al., 2013;
50 Degeling et al., 2014; Shen et al., 2015), periodic solar wind dynamic pressure variations (e.g.,
51 Kepko, 2002; 2003), drift-bounce resonance (e.g., Southwood et al., 1969) and dynamic processes during
52 substorms (e.g., Olson, 1999).

53 Although many previous studies have focused on waves occurring in the dayside magnetosphere
54 (e.g., Samson et al., 1981; Rostoker et al., 1984; Zong et al., 2007; Shen et al., 2017), ULF waves
55 occurring on stretched magnetic field lines in the magnetotail have also been reported in some
56 observational studies (e.g., Zheng et al., 2006; Tian et al., 2012) and simulations (e.g., Rankin et al., 2000;
57 Lui and Cheng, 2001). Pc5 (150-600 s) and Pc6 (600-1800 s) waves are the primarily waves occurring at
58 high latitudes and in the magnetotail. Investigating the source and characteristics of these waves in the
59 magnetotail will help us further understand the solar wind-magnetosphere-ionosphere coupling processes
60 in the night side region.

61 Statistical studies of ULF wave properties in the magnetosphere have been performed using various
62 satellites (e.g., Hudson et al., 2004; Liu et al., 2009; Takahashi et al., 2015). Hudson et al. (2004)
63 performed a statistical study of the occurrence rate of Pc5 magnetic pulsations for both toroidal and
64 poloidal modes at L values from 4 to 9 by using 14 months magnetometer data from Combined Release
65 and Radiation Effects Satellite (CRRES). They found that there is no dawn-dusk asymmetry on the
66 occurrence rate of toroidal mode oscillations inside L=8, however the occurrence rate of poloidal mode
67 oscillations is higher on dusk side. Liu et al. (2009) statistically studied the both occurrence and frequency
68 distributions of Pc5 magnetic pulsations in toroidal and poloidal modes between L= 4 and 9 by using 13
69 months electric and magnetic field measurements from THEMIS-D. They found that the occurrence
70 distribution is similar to the results of Hudson et al. (2004) and the frequency is higher in the dawn side
71 than in the dusk side by a factor 2 and decreases with radial distance. Takahashi et al. (2015) statistically
72 investigated the fundamental toroidal mode oscillations from L = 7 to 12 by using 2008-2013 ion bulk



73 velocity data from THEMIS-D. They found that the occurrence rate and amplitude of toroidal mode
74 oscillations are higher in the dawn side (4-8 MLT) than in the dusk side (16-20 MLT). Moreover, the
75 relationship between ULF wave characteristics and the solar wind conditions/geomagnetic activity level
76 were also studied statistically (e.g. Takahashi and Ukhorskiy, 2007; Susumu Kokubun, 2013; Wang et al.,
77 2015). Takahashi and Ukhorskiy (2007) found that the solar wind dynamic pressure variance has the best
78 correlation with the power of magnetic pulsations at geosynchronous orbit. Susumu Kokubun (2013)
79 statistically studied Pc5 ULF waves (mostly on the 4-8 MLT and 16-20 MLT) using GEOTAIL data during
80 the period of 1995 to 2000. They found that the wave occurrence tends to be larger for higher solar wind
81 velocity (> 400 km/s), smaller IMF Bz, and lower cone angle. Wang et al. (2015) studied the spatial
82 distribution of the Pi2 and Pc4-5 magnetic fluctuation power in the plasma sheet by using THEMIS-
83 A/C/D/E data from 2007 to 2014. They found that the amplitude of Pc-5 fluctuations is larger globally
84 during periods of higher AE index, faster solar wind, and larger solar wind dynamic pressure variations.

85 Although statistical studies of ULF waves have been performed, most have focused on the dayside
86 or near-earth region. The distributions and excitation mechanisms of ULF waves on stretched magnetic
87 field lines are still unclear. Our work focuses on ULF waves on stretched magnetic field lines ($X^*_{\text{GSM}} < 0$
88 and $8 R_E < R < 32 R_E$).

89 This paper will be organized as follows. In section 2, the data set and the selection criteria of the ULF
90 wave event are presented. In section 3, we show the statistical results. In section 4, we discuss the
91 occurrence and frequency distributions of ULF waves on the stretched field lines and the influence factors
92 of solar wind parameters and geomagnetic activity level. The main conclusions of this study are given in
93 section 5.

94

95 **2 Data and statistical methods**

96

97 In this study, we use 3 s resolution magnetic field data from Flux Gate Magnetometer (FGM) (Auster et
98 al., 2008) and 3 s resolution plasma data from Electrostatic Analyzer (ESA) (Mcfadden et al., 2008) of
99 THEMIS mission from 2008 to 2015. The THEMIS mission consists of five satellites (THEMIS
100 A/B/C/D/E), each with an orbital inclination of about 10° (Angelopoulos, 2008). In the first two years,
101 the apogees were about $12 R_E$ for THEMIS A/D/E, $20 R_E$ for THEMIS C and $30 R_E$ for THEMIS B. After
102 2010, THEMIS B/C were transferred to a lunar orbit which is about $60 R_E$ from earth. Because THEMIS
103 A/D/E have similar orbits, in this study we only use data from THEMIS A, B and C. In addition, we use
104 1 minute resolution interplanetary magnetic field (IMF) and solar wind plasma data from the OMNI
105 database (<https://spdf.sci.gsfc.nasa.gov/>), which is calculated by time shifting satellite data taken in the
106 solar wind to the Earth's bow shock sub-solar point. Figure 1 shows the binned spatial distribution of the
107 total observation time over the 2008-2015 interval for THEMIS A/B/C in the magnetosphere.

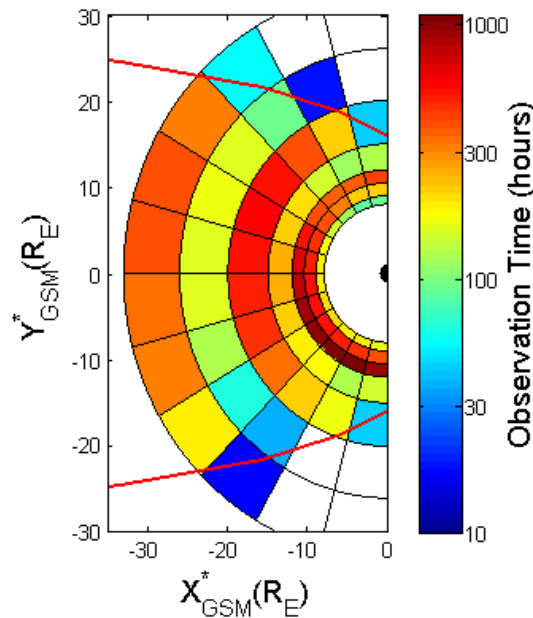
108 We use the aberration coordinate GSM* whose X axis rotated 4° from the X axis of GSM
109 coordinates for spacecraft position to remove the effect of Earth's revolution. Field-aligned coordinates



110 (FAC) are used to analyze waves and separate the azimuthal and radial oscillating wave components. The
 111 FAC system is defined in Eq. (1).

112
$$\mathbf{z} = \frac{\mathbf{B}_0}{|\mathbf{B}_0|}; \mathbf{a} = \frac{\mathbf{z} \times \mathbf{R}}{|\mathbf{z} \times \mathbf{R}|}; \mathbf{r} = \mathbf{a} \times \mathbf{z} \quad (1)$$

113 In this equation, \mathbf{B}_0 is the background magnetic field vector, derived by taking a 30 minutes sliding
 114 average of the magnetic field data, \mathbf{R} is the vector from Earth's center to the satellite, \mathbf{z} is the parallel unit
 115 vector, \mathbf{a} is the unit vector pointing east and \mathbf{r} completes the right-hand rule. It should be noted that the
 116 direction of \mathbf{r} is approximately radial due to the equatorial orbits of THEMIS.



117 **Figure 1.** The distribution of total observation time of THEMIS A/B/C in the GSM* X-Y plane between
 118 2008 and 2015. The red line is the average magnetopause, calculated by Shue et al.'s (1998) model with
 119 $D_p=1.66$ nPa and $B_z=0.16$ nT. The blank bins indicate regions where the residence time of THEMIS is
 120 less than 10 hours.
 121

122
 123 In this study, we mainly use ion velocity data to identify ULF waves, following the technique of
 124 Takahashi (2015). They suggested that using velocity is better than using magnetic field data, because
 125 fundamental mode magnetic field fluctuations (considered most likely in the Pc5 range) give rise to a
 126 node near the equatorial plane, making their measurement problematic along the low-inclination THEMIS
 127 orbital path. On the other hand, the fundamental mode has an antinode for the electric field and plasma
 128 velocity fluctuations under ideal MHD conditions. The electric field data is therefore estimated by $\mathbf{E} =$
 129 $-\delta\mathbf{V} \times \mathbf{B}$, where $\delta\mathbf{V}$ indicates the variation of velocity, which is obtained by subtracting the 30 minutes
 130 sliding average values.



131 As shown in Fig. 1, the region concerned in this work is $X_{\text{GSM}}^* < 0 R_E$ and $8 R_E < R < 32 R_E$. In order
132 to remove the likelihood of identification of ULF wave events when THEMIS enters the magnetosheath
133 or solar wind regions, only events for which density values less than 1 cm^{-3} if $|Y_{\text{GSM}}^*| > 10 R_E$ are included
134 in the database.

135 The following criteria are used to select ULF waves in the magnetotail: (i) the wave frequency is
136 below 7 mHz; (ii) the wave is quasi monochromatic, and includes at least three cycles; (iii) the maximum
137 of peak to trough value of fluctuations is more than 50 km/s; (iv) mirror-like structures, indicated by anti-
138 phase variations of magnetic field and density are excluded; (v) magnetotail flapping events, characterized
139 by sign changes in Bx are excluded. A quantitative standard is used to determine the beginning and ending
140 time of each event, namely that the beginning and ending time is at the points where the amplitude is 20
141 km/s. Additionally, if the interval time between two events is less than 20 minutes and they have similar
142 frequency (within 0.5 mHz), we consider them as a single event.

143 The process of selecting wave events and distinguishing the wave mode in this study is as follows.
144 Firstly, we conduct wavelet analysis to THEMIS ion velocity and magnetic field data in GSM coordinates
145 and choose the wave events which roughly satisfy criteria mentioned above. Then, we transform from
146 GSM to FAC coordinates for magnetic field and ion velocity data, and calculate the electric field. To
147 quantitatively distinguish the azimuthal or radial oscillating waves, Fast Fourier Translation (FFT)
148 analysis is applied to all three components of ion velocity (Fig. 2).

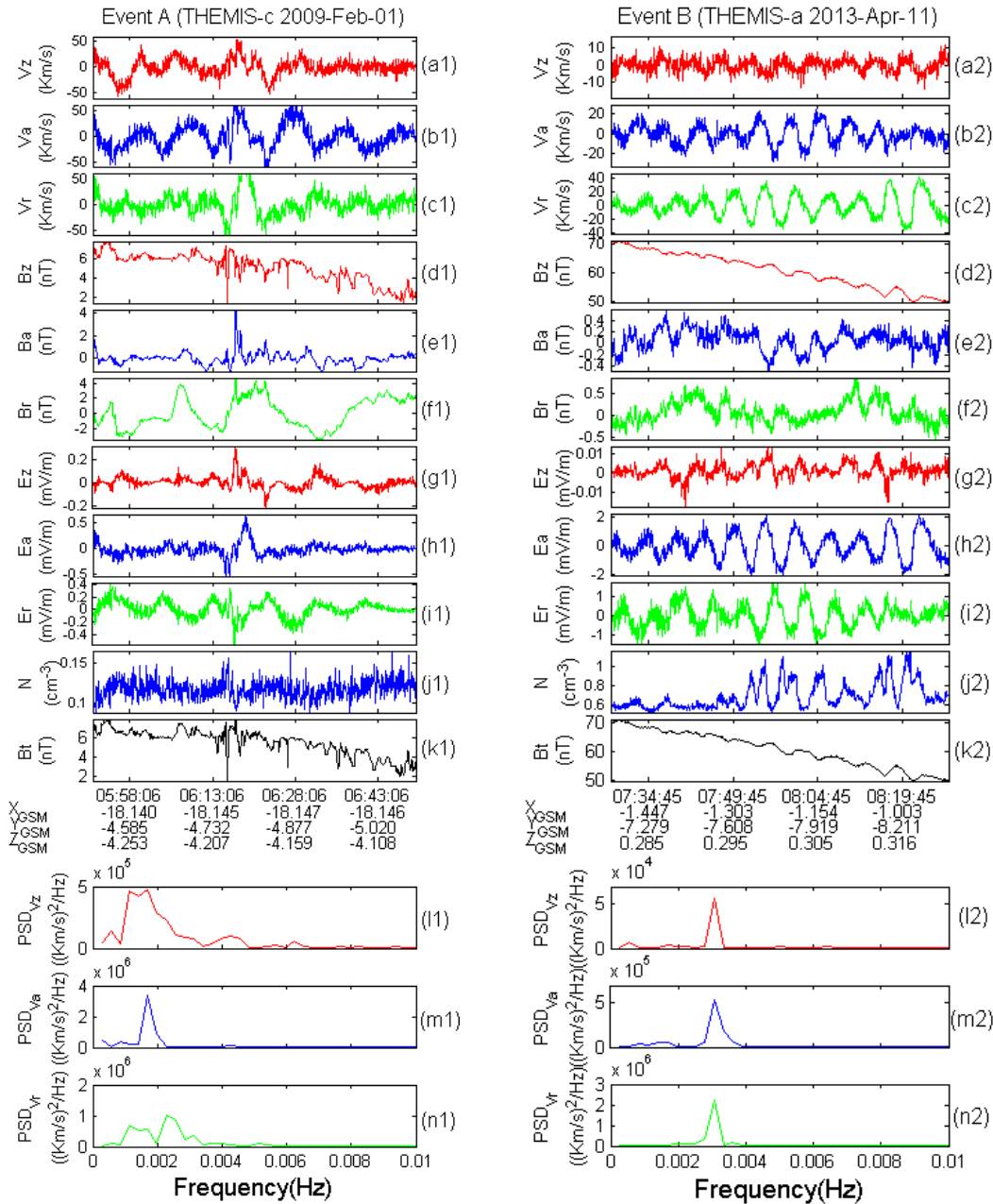
149 Figure 2 shows two typical events (labelled “A” and “B”) with Event A occurring near $R \approx 19 R_E$,
150 from 0550 to 0650 UT on 01 February 2009 and showing azimuthal oscillations, and Event B occurring
151 near $R \approx 8 R_E$ from 0728 to 0828 on 11 April 2013 and showing radial oscillations. Figure 2 shows three
152 components of the ion velocity (a-c), magnetic field (d-f), and the calculated electric field (g-i), in addition
153 to the total ion density (j) and total magnetic field (k) which are used for excluding mirror-like structures.
154 Figure 2(l-n) show the Power Spectral Density (PSD) of three components of the ion velocity derived by
155 FFT. In events A and B, the peak in PSD of the dominant wave component exceeds its counterpart by a
156 factor of 4, enabling their unambiguous designation as an azimuthal and radial oscillation event
157 respectively. Events for which the peak in PSD in V_a and V_r have similar magnitudes are simply regarded
158 as both an azimuthal oscillating event and a radial oscillating event.

159 In total, we find 1089 azimuthal oscillating wave events and 566 radial oscillating wave events, with
160 an average event-time duration of ~ 54 minutes.

162 3 Statistical Analysis

163 3.1 Occurrence rate

164 Figure 3 shows the occurrence rates of azimuthal oscillating wave events (left panel) and radial oscillating
165 wave events (right panel) in the $\text{GSM}^* X\text{-}Y$ plane. The color in each bin indicates the occurrence rate

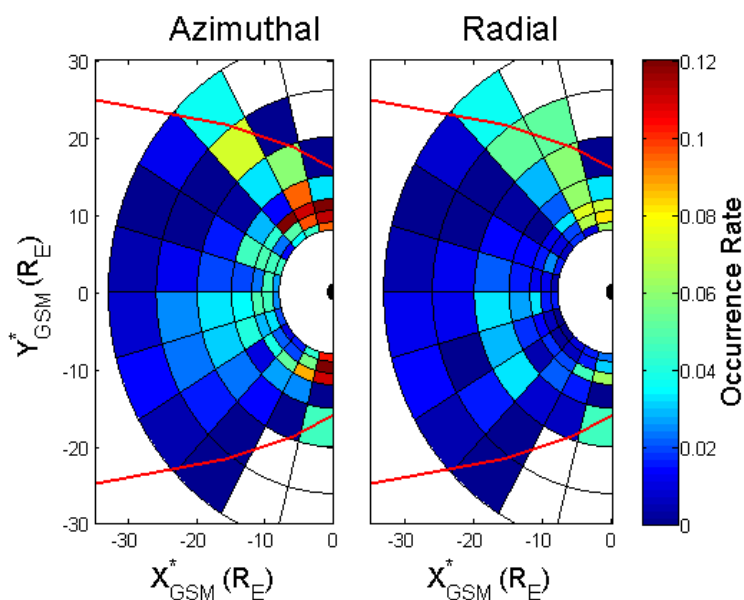


168

169 **Figure2.** Examples of an azimuthal oscillating event (Event A) from 0550 to 0650 UT on 01 February
 170 2009, and radial oscillating event (Event B) from 0728 to 0828 on 11 April 2013: (a-c) velocity
 171 components, (d-f) magnetic components, (g-i) electric field components, (j) total ion density, (k) total
 172 magnetic field, (l-n) FFT analysis of ion velocity.
 173



174 calculated by dividing the total event times by the total observation times shown in Fig. 1. In the near-
175 earth region ($8 R_E < R < 12 R_E$), we can see that the occurrence rates of both azimuthal and radial
176 oscillating events in the dusk and dawn flanks (18-21 MLT and 3-6 MLT) are higher than the midnight
177 regions (21-03 MLT). For azimuthal oscillating events, there is no clear dawn-dusk asymmetry in the
178 occurrence rates, while for the radial oscillating events, the occurrence rates of waves are higher on the
179 dusk side than dawn side. In the magnetotail region ($12 R_E < R < 32 R_E$), the occurrence rates of both
180 modes of waves are slightly higher in the post-midnight region. Note that, although no wave events are
181 found in the dawn side flank region ($20 R_E < R < 32 R_E$, 3-6 MLT), the total observation time is also very
182 short (< 38 hours) in this region. So, we cannot conclusively say that the occurrence rates on the dusk
183 side flank region is higher than that of the dawn side.



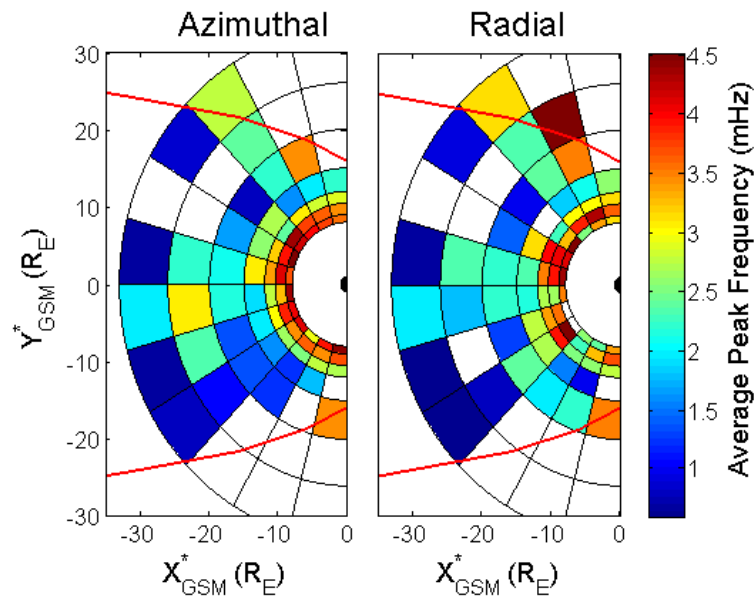
184
185 **Figure 3.** The occurrence rates of azimuthal oscillating wave events (left panel) and radial oscillating
186 wave events (right panel) in the GSM* X-Y plane.

187 188 3.2 Frequency distribution

189
190 Figure 4 shows the spatial distribution of average frequency for azimuthal (left panel) and radial (right
191 panel) oscillating wave events in the equatorial plane. The color in each bin is the average of all event
192 frequencies (obtained by FFT analysis as described earlier) in that bin. A blank bin inside the
193 magnetopause indicates that there are no events. It can be seen roughly that the frequency decreases with
194 increasing radial distance both for azimuthal and radial oscillating wave events, for regions where



195 $R < 15R_E$. Note that the crimson bin in the upper right corner (19-20 MLT and $20 < R < 26 R_E$) of the right
196 panel is caused by short residence time (~ 19 hours) and only one wave event with frequency of 5.71 mHz.

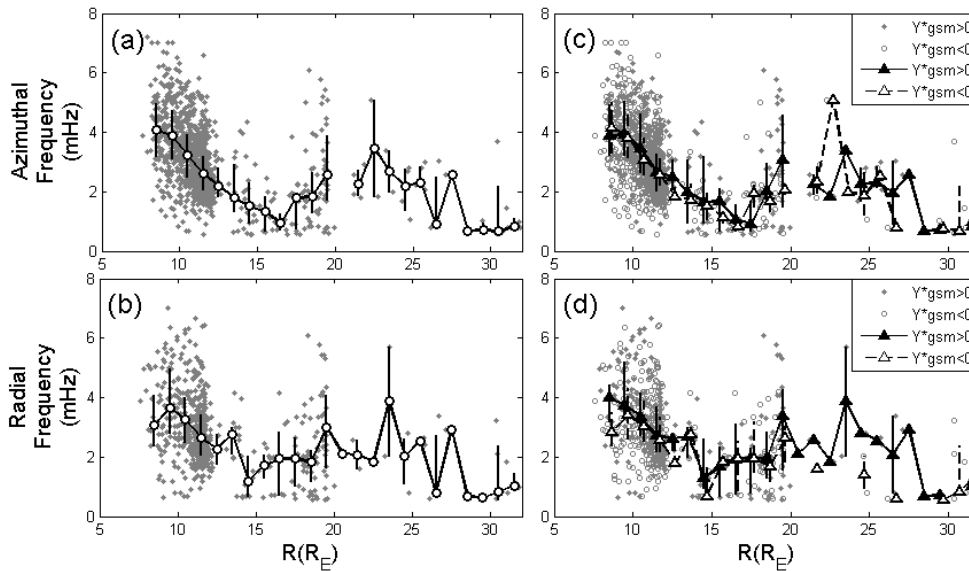


197

198 **Figure 4.** The average frequencies of azimuthal oscillating wave events (left panel) and radial oscillating
199 wave events (right panel) in the GSM* X-Y plane.

200

201 We further plot the relationship between the peak frequency and the distance from earth in Fig. 5. It
202 shows that the frequency can be as low as 0.55 mHz. As shown in Fig. 5a and 5b, the median frequency
203 of azimuthal oscillating events decreases with increasing radial distance from the Earth in the region with
204 $8 R_E < R < 16 R_E$, and the same trend is found for the radial oscillating events with $8 R_E < R < 14 R_E$.
205 Figure 5c and 5d show frequency distribution of events in the dawn side ($Y_{gsm}^* < 0$) and dusk side
206 ($Y_{gsm}^* > 0$) regions, respectively. The frequency for both azimuthal and radial oscillating events show
207 no obvious dawn dusk asymmetry. This is verified by the Wilcoxon rank sum test applied to the dawn and
208 dusk datasets. The Wilcoxon rank sum test is a non-parametric statistical hypothesis test that can be used
209 to assess whether two samples have the same distribution or not (Gibbons and Chakraborti, 2011).
210 Specifically, in the Wilcoxon rank sum test, a “P-value” result greater than 0.01 means there is no
211 significant statistical difference between two datasets. The P-value for the dawn and dusk side data sets
212 is 0.4535 (for all azimuthal and radial oscillating events). This confirms that the dawn and dusk side
213 frequency data sets belong to the same distribution.



214

215 **Figure 5.** The wave frequency versus radial distance for azimuthal (a and c) and radial (b and d) oscillating
 216 event. In panels a and b, the grey dots are individual events, the open circles are the median values of
 217 frequencies in each $1 R_E$ bin. The vertical bars connect the lower and upper quartiles. In panel c and d, the
 218 grey dots and circles indicate the dusk and dawn events, respectively. The solid and open triangles are the
 219 same as the open circles in Fig. 5a, but for dusk and dawn events respectively.

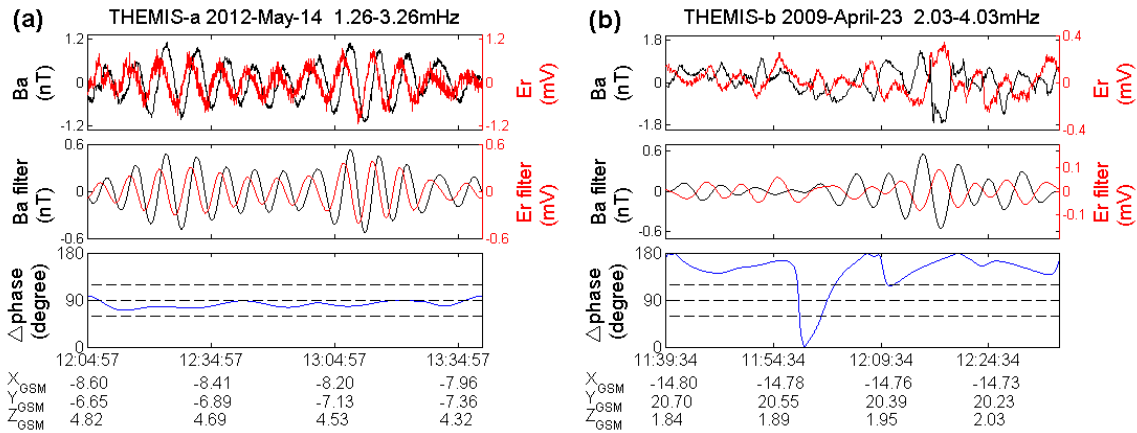
220

221 3.3 Standing wave

222

223 According to Singer et al. (1982), Alfvénic standing wave oscillations are characterized by a phase
 224 difference of 90° between the electric field and magnetic field components. Figure 6 shows the standing
 225 wave analysis of two azimuthal oscillating events. The first row shows the magnetic field component B_a
 226 and electric field component E_r . The second row shows the 1.26-3.26 mHz (Fig. 6a) and 2.03-4.03 mHz
 227 (Fig. 6b) band-pass filtered B_a and E_r components. The phase differences between the band-pass filtered
 228 B_a and E_r are shown in the bottom panels, in which three dotted lines indicate the 60° , 90° , 120° phase
 229 differences respectively. We can see that the first event (Fig. 6a) shows characteristics of standing wave
 230 as indicated by the $\sim 90^\circ$ phase difference between B_a and E_r , while the second event (Fig. 6b) does not
 231 have this characteristic. We quantify the criteria of standing azimuthal (radial) oscillating waves as that
 232 the phase differences between the filtered B_a and E_r (B_r and E_a) that falls within the range 60° - 120° and
 233 lasts for at least three cycles.

234



235

236 **Figure 6.** Examples of: (a) a standing azimuthal oscillating event and (b) a non-standing azimuthal
 237 oscillating event.

238

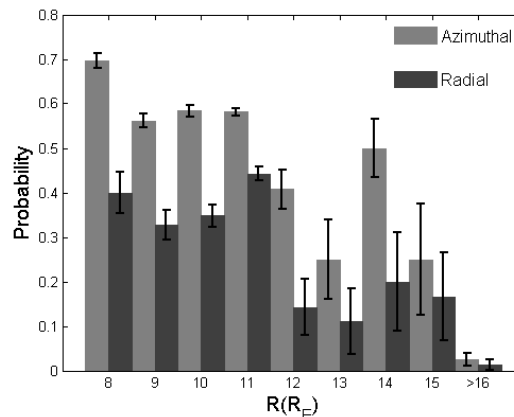
239 Figure 7 shows the distribution of the probability of standing waves with the distance R. The light
 240 and deep histogram represents the probability for azimuthal and radial oscillating event respectively. The

241 errorbars shown are calculated by $\varepsilon = \frac{n}{N} * \left(\frac{\sqrt{n}}{n} + \frac{\sqrt{N}}{N} \right)$, where n is the number of standing wave events and

242 N is the total number of waves events in each bin for each polarization. It is obvious that standing waves
 243 occupy a larger proportion in the region of 8-16 R_E , while almost no standing waves are identifiable in

244 the region of 16-32 R_E . We find that about 52% events (including the azimuthal and radial oscillating
 245 events) are standing waves in the region of 8-16 R_E , while only 2 % are standing waves in the region of

246 16-32 R_E . This figure also shows that the probability of standing waves is higher for the azimuthal
 247 oscillating events than for the radial oscillating events.



248

249 **Figure 7.** The radial distribution of the probability of identifying standing waves, for azimuthal and radial
 250 oscillating events (light and dark histograms respectively).



251

252 **4 Discussion**

253

254 Using THEMIS data during the period from 2008 to 2015, we find 1314 Pc5-6 ULF wave events in the
255 region of $X_{\text{GSM}}^* < 0$ and $8 R_E < R < 32 R_E$. The elevation angle of the magnetic field of each event was
256 calculated by the formula $\tan^{-1} \left(\frac{B_z}{\sqrt{B_x^2 + B_y^2}} \right) * \frac{180}{\pi}$, where B_x , B_y , B_z are the three magnetic field
257 components in GSM* coordinates. We find that 58.25% of the events have an elevation angle larger than
258 45° . This suggests that most of our events are observed near the magnetic equatorial plane. It is reasonable
259 to consider that most of our standing wave events belong to the fundamental eigenmode. Furthermore,
260 Lui and Cheng. (2001) indicated that the magnetic field line in the nightside is very stretched in the region
261 of $R > 8 R_E$, especially during intervals of high Kp index. We therefore consider it likely that most of our
262 events should be observed on stretched magnetic field lines.

263

264 **4.1 Occurrence rate**

265

266 As shown in Fig. 3, in the region of $8 R_E < R < 12 R_E$, there is no obvious dawn-dusk asymmetry in the
267 occurrence rates for azimuthal oscillating waves, while the occurrence rates are higher on the dusk side
268 than dawn side for radial oscillating waves. This is consistent with the wave mode distributions in the
269 inner magnetosphere ($4 R_E < R < 9 R_E$) presented in previous works (Hudson et al., 2004, Liu et al., 2009).
270 One possible reason is that westward drifting ions injection associated with substorm may excite more
271 radial oscillating wave events in the dusk side via the ion drift bounce resonance (Southwood et al., 1969;
272 Chen and Hasegawa, 1988). In contrast to that of the inner magnetosphere ($4 R_E < R < 9 R_E$), the
273 occurrence rates for both azimuthal and radial oscillating events in the region of $12 R_E < R < 32 R_E$ are
274 slightly higher on the post-midnight region than the pre-midnight region. It is possible that the K-H
275 instability may play an important role on the generation of ULF waves on the stretched magnetotail, given
276 that the K-H instabilities are inclined to occur in the dawn side than in the dusk side (Nykyri et al., 2013)
277 and even can happen in the down tail flanks up to the lunar orbit ($\sim 60 R_E$) (Wang et al., 2017). In view of
278 the limited observation times in the dawn side magnetopause, more events are needed in further study on
279 the definite reasons of the dawn-dusk asymmetry of occurrence rate in the outer side region ($12 R_E < R <$
280 $32 R_E$).

281

282 **4.2 Frequency distribution**

283

284 As shown in Fig. 5a and 5b, the frequency decreases with increasing radial distance from the Earth for
285 both azimuthal oscillating events ($8 R_E < R < 16 R_E$) and radial oscillating events ($8 R_E < R < 14 R_E$). This
286 is consistent with the Alfvén continuum of field line resonance (FLR) theory (e.g., Allan and



287 Poulter, 1992; Waters et al., 2000). However, this trend does not continue for $R > 16 R_E$. Previous
288 observation and simulation studies have shown that standing waves can exist on the stretched magnetic
289 field lines (Lui and Cheng, 2001; Zheng et al., 2006; Tian et al., 2012). Our statistical result shows that
290 52 % of all event types are standing waves in the region of 8-16 R_E , while only 2 % can be confirmed as
291 standing waves in the region of 16-32 R_E as shown in Fig. 7. Given the likelihood that most of our wave
292 events belong to the fundamental mode, the uncertainty in the phase measurement of the weak magnetic
293 field signal near the equatorial plane will affect the identification of standing waves. Moreover, the
294 complicated phase relationship between the electric field and the magnetic field caused by magnetic field
295 disturbances in the farther deeper magnetotail will also affect the identification of standing waves. These
296 suggest that our data may underestimate the proportion of standing wave events. Even so, the finding that
297 only 2 % of events in the down-tail region ($R > 16 R_E$) can be identified as standing waves suggests that
298 the standing waves are far less common on the highly stretched field lines.

299 As shown in Fig. 5c and 5d, there is no obvious dawn-dusk asymmetry in the ULF wave frequency
300 for $8 R_E < R < 32 R_E$. This is different from previous studies in the near-earth region (Liu et al., 2009;
301 Takahashi et al., 1982; 2015). Takahashi et al. (1982) found that the frequencies of Pc3-4 ULF waves
302 were higher on the dawn side than dusk side at geosynchronous orbit. They suggested that the quasi-
303 parallel shock and the associated turbulent magnetosheath flow is more likely to occur on the dawn side,
304 which leads to higher harmonic waves to be excited in the dawn side. Takahashi et al. (2015) found that
305 the frequencies of Pc5 toroidal waves in the region with L values between 7 and 12 R_E is lower in the
306 dusk side (16-20 MLT) than dawn side (04-08 MLT). They suggest that this is due to the higher mass
307 density in the dusk side near-earth region, supplied by the particles from ionosphere. However, the wave
308 frequency distributions shown in this paper ($X^*_{GSM} < 0$, $8 R_E < R < 32 R_E$) show a different distribution
309 feature from that of the events in the inner or dayside magnetosphere. This suggests that neither of the
310 above mechanisms for producing asymmetry are important within the region of interest in our study. This
311 may be expected for the turbulent magnetosheath flow mechanism more applicable to higher frequencies.
312 The influence of particle injection from the ionosphere may be weakened by higher ExB drift speeds and
313 longer field line lengths in the nightside magnetotail region, compared to the near-earth region.

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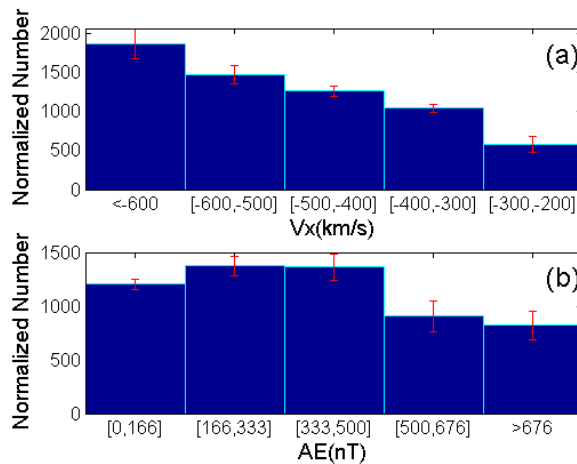
315 **4.3 The influence of Solar wind parameters and geomagnetic activity level**

316

317 Figure 8 shows the relationship between the occurrence rate of wave events and solar wind velocity V_x
318 (panel a) and the AE index (panel b). The Y-axis indicates the normalized event number, which is
319 calculated by dividing the number of waves with the proportion of background solar wind velocity V_x
320 (panel a) and the background AE index (panel b) in each bin. The background values are calculated with
321 OMNI data from 2008 to 2015. We can see that the ULF waves occurrences increase with increasing solar
322 wind speed $|V_x|$. This implies that the K-H instability could be a sources of ULF waves in the magnetotail
323 region ($8-32 R_E$), since the higher shear velocity is an important factor for exciting K-H instabilities



324 (Miura, 1992). Figure 8b shows that the waves occurrences are higher when the AE values are less than
325 500 nT. Note that about 74.8 % of the waves occurred when the AE values are less than 250 nT. This
326 suggests that most of the wave events in the magnetotail are observed during quiet times or weak substorm
327 times. The relationship between the occurrences of ULF waves and the relative variation of solar wind
328 dynamic pressure (P_d) and the IMF B_z values were also examined (not shown). We find that the
329 occurrence rates are higher for larger solar wind P_d variance and during periods of northward IMF B_z .

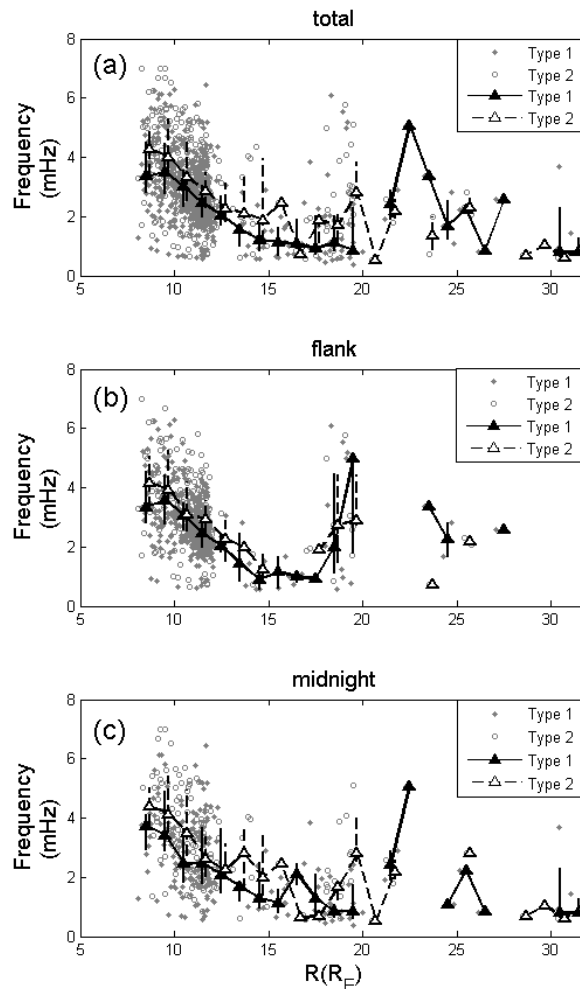


330
331 **Figure 8.** The normalized occurrence number of waves versus (a) Solar wind velocity V_x , and (b) AE
332 index.

333
334 The possibility that substorm activity may affect the frequency of ULF waves, and thereby influence
335 the distribution of ULF frequencies in our database, is examined using the following method, based on
336 the substorm event list of (Forsyth et al., 2015). The ULF wave events were divided into two categories
337 based on their start time relative to the onset time of individual substorm events. The first category (“type
338 one”) consists of events that occurred more than two hours after the most recent substorm onset, and more
339 than one hour before the next substorm onset. These events are considered to be independent of substorm
340 activity. The second category (“type two”) consists of events that occurred between zero and two hours
341 after the most recent substorm onset. In principle a third category consisting of events that occur less than
342 one hour before the next substorm onset could be defined, however this category contains very few events,
343 so their frequency characteristics will not be discussed here. The radial dependence of median frequency
344 for type one and two events is shown in Fig. 9a. This plot clearly shows that the median frequencies for
345 type two events are higher than type one events. A plausible explanation for this difference could be that
346 field line depolarization following substorm onset results in an increase in local magnetic field strength
347 compared to more stretched magnetotail field lines during quiet times. The resulting higher Alfvén speed
348 profile raises the fundamental mode eigenfrequency for the type two events, compared to the type one
349 events.



350 Figure 9b and 9c show the radial dependence of median frequency for type one and two events
351 occurring in the dawn/dusk flank (3-6 MLT and 18-21 MLT) and midnight sectors (21-03 MLT),
352 respectively. According to these plots, the frequency differences between type one and type two wave
353 events are more obvious in the midnight region than in the flank region. This is understandable, given
354 that the configuration of field lines will be changed much more in the midnight region than in the flank
355 regions during substorm times. It should be noted that, only the possible influence of field lines
356 configuration or plasma environment associated with weak substorms on the ULF wave frequencies are
357 discussed here. The question of whether substorms could trigger or be triggered by ULF waves still cannot
358 be answered by the present analysis.



359 **Figure 9.** The wave frequency versus the distance from the earth for: (a) the type one and type two wave
360



361 events, (b) the wave events in the flank region, and (c) the wave events in the midnight region, respectively.
362 The grey dots and circles indicate type one and type two wave individual events respectively. The solid
363 and open triangles are the median values of frequencies in each $1R_E$ bin for the type one and type two
364 wave event respectively. The vertical bars connect the lower and upper quartiles for each category.

365 366 **5 Summary**

367
368 We have statistically studied the distributions of the occurrence rate and frequency of the Pc5-6 ULF
369 waves in the region of $X_{GSM}^* < 0$ and $8 R_E < R < 32 R_E$ (occurring mostly on stretched magnetic field
370 lines) using 8 years of THEMIS data. We also examined the influence of Solar wind parameters and
371 geomagnetic activity level on the features of these ULF waves. Some new results that differ from those
372 of ULF waves observed in the inner magnetosphere are obtained. The main results are summarized as
373 follows:

374 (1) In the far magnetotail region ($12 R_E < R < 32 R_E$), the occurrence rates of both azimuthal and
375 radial oscillating events are higher in the post-midnight region than in the pre-midnight region. In the
376 near-earth magnetotail ($8 R_E < R < 12 R_E$), the occurrence rates of azimuthal oscillating events are
377 comparable on the dawn and dusk side, while the occurrence rates of radial oscillating events are higher
378 on the dusk side, which is similar to the distributions in the inner magnetosphere ($4 R_E < R < 9 R_E$).

379 (2) Statistically, the peak frequency decreases as the increase of radial distance from Earth for both
380 azimuthal oscillating events ($8 R_E < R < 16 R_E$) and radial oscillating events ($8 R_E < R < 14 R_E$). A possible
381 explanation for this distribution is that at least 52 % events (including both azimuthal and poloidal
382 oscillating events) are standing waves in the region of 8-16 R_E , while only 2 % are unambiguous standing
383 waves in the region of 16-32 R_E . Moreover, the frequencies for all the events in this paper do not show
384 obvious dawn-dusk asymmetry as found in previous studies for waves in the inner magnetosphere ($4 R_E$
385 $< R < 9 R_E$), where the wave frequencies are higher in the dawn side than in the dusk side.

386 (3) The ULF wave occurrence rates are higher for larger solar wind velocity and solar wind Pd
387 variations. Therefore, we suggest that the solar wind maybe the main energy source of the ULF waves in
388 the region of $8 R_E < R < 32 R_E$. About 74.8 % of the ULF waves occurred when the AE values are less
389 than 250nT, which indicates that the ULF waves most likely to occur during in the quiet times or weak
390 substorm times. We have further studied the frequency change between the quiet time and the weak
391 substorm time events. We found that the wave frequency is higher during the substorm time (0-2 hours
392 after substorm onset). The frequency differences are clearer in the midnight region than in the flank region.
393 We suggest that the field lines configuration or plasma environment variation during weak substorm times
394 could increase the eigen frequencies of ULF waves in the magnetotail, leading to the observed change in
395 the frequency distribution.

396
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398 <http://themis.ssl.berkeley.edu/>, and SPDF web service for OMNI data at <https://spdf.sci.gsfc.nasa.gov/>.
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