

1 Multisatellite observations of the magnetosphere response to changes in the solar wind and inter-
2 planetary magnetic field

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15

16 **Abstract**

17 We employ multipoint observations of the Van Allen Probes, THEMIS, GOES and Cluster to
18 present case and statistical studies of the electromagnetic field, plasma and particle response to
19 interplanetary (IP) shocks observed by Wind. On February 27, 2014 the initial encounter of an
20 IP shock with the magnetopause occurred on the postnoon magnetosphere, consistent with the
21 observed alignment of the shock with the spiral IMF. The dayside equatorial magnetosphere
22 exhibited a dusk-dawn oscillatory electrical field with a period of ~ 330 s and peak to peak
23 amplitudes of ~ 15 mV/m for a period of 30 min. The intensity of electrons in the energy
24 range from 31.5 to 342 KeV responded with periods corresponding to the shock induced ULF
25 electric field waves. We then perform a statistical study of E_y variations of the electric field
26 and associated plasma drift flow velocities for 60 magnetospheric events during the passage of
27 interplanetary shocks. The E_y perturbations are negative (dusk-to-dawn) in the dayside
28 magnetosphere (followed by positive or oscillatory perturbations) and dominately positive
29 (dawn-to-dusk direction) in the nightside magnetosphere, particularly near the Sun-Earth line
30 within an L-shell range from 2.5 to 5. The typical observed amplitudes range from 0.2 to 6
31 mV/m but can reach 12 mV during strong magnetic storms. We show that electric field
32 perturbations increase with solar wind pressure and that the changes are especially marked in
33 the dayside magnetosphere. The direction of the V_x component of plasma flow is in agreement
34 with the direction of the E_y component and is antisunward at all local times except the nightside

35 magnetosphere, where it is sunward near the Sun-Earth line. The flow velocities V_x range from
36 0.2 to 40 km/s and are a factor of 5 to 10 times stronger near noon as they correspond to greater
37 variations of the electric field in this region. We demonstrate that the shock-induced electric
38 field signatures can be classified into four different groups according to the initial E_y electric
39 field response and these signatures are local time dependent. Negative and bipolar pulses
40 predominate on the dayside with positive pulses occur on the nightside. The ULF electric field
41 pulsations of Pc and Pi types produced by IP shocks are observed at all local times and in the
42 range of periods from several tens of seconds to several minutes. We believe that that most
43 electric field pulsations of the Pc5 type in the dayside magnetosphere at $L < 6$ are produced by
44 field line resonances. We show that the direction of the shock normal determines the direction of
45 the propagation of the shock-induced magnetic and plasma disturbances. The observed direc-
46 tions of velocity V_y predominately agree with those expected for the given spiral or orthospiral
47 shock normal orientation.

48

49 **Keywords:** Interplanetary shocks, solar wind – magnetosphere interactions, energetic particles,
50 trapped.

51

52 **1 Introduction**

53 Sudden increases in the solar wind dynamic pressure accompanying interplanetary (IP)
54 shocks cause earthward motion of the bow shock and the magnetopause and launch fast and in-
55 termediate mode waves into the magnetosphere (Tamao, 1964). The fast mode waves propagate
56 both radially inward and azimuthally around the Earth (Araki et al., 1997) whereas the interme-
57 diate mode waves propagate along magnetic field lines to produce transient perturbations in the
58 high-latitude dayside ionosphere (Southwood and Kivelson, 1990; Glaßmeier and Heppner,
59 1992). Using multipoint observations) estimated the propagation speeds to be about 600 km/s in
60 the radial direction from geostationary orbit to the ground and about 910 km/s in the azimuthal
61 direction in the equatorial plane. Nopper et al. (1982) estimated an impulse disturbance speed of
62 about 1500 km/s at geostationary orbit. Schmidt and Pedersen (1988) derived a propagation ve-
63 locity for the radially inward travelling compressive wave of 950 km/s and for the azimuthal
64 wave in the outer magnetosphere of 1100 km/s. Samsonov et al. (2007) used a magneto-hydro-

65 dynamic code to simulate the interaction of a moderately strong interplanetary shock propagat-
66 ing along the Sun-Earth line and obtained the average speed of the primary and reflected fast
67 shocks in the magnetosphere to be about 700 km/s, in agreement with their assumptions con-
68 cerning the mean Alfvén velocity in the outer dayside magnetosphere (1000 km/s) and in the
69 plasmasphere (500 km/s).

70 The IP shock orientation plays an important role in determining the associated geophysical
71 effects (e.g., Oliveira and Raeder, 2015) showed that system evolution times are much longer for
72 shocks with normals oblique to the Sun–Earth line. The pressure pulse model of Sibeck (1990)
73 predicts dawnward moving transient events near local noon when shock normals point per-
74 pendicular to the nominal spiral interplanetary magnetic field (IMF) direction, but duskward
75 moving events occur near local noon for events when shock normals point perpendicular to the
76 orthospiral IMF orientation. The direction of the plasma flow within the magnetosphere is ex-
77 pected to be consistent with the orientation of the shock. That is to say dawnward flow for spi-
78 ral IMF shocks and duskward flow for orthospiral IMF shocks. Here orthospiral refers to IMF
79 longitudes ($0^\circ < \Lambda < 90^\circ$ and $180^\circ < \Lambda < 270^\circ$), spiral refers to IMF longitudes ($90^\circ < \Lambda < 180^\circ$
80 and $270^\circ < \Lambda < 360^\circ$), where longitude $\Lambda=0^\circ$ points sunward, and $\Lambda=90^\circ$ duskward.

81 The magnetic and electric fields are key parameters for understanding of the response of
82 the Earth’s space environment to IP shocks. The propagation and evolution of electric fields in
83 the magnetosphere-ionosphere system in response to IP shocks have been studied for several
84 decades but signatures of the shock related electric field perturbations are still not fully under-
85 stood. Knott et al. (1985) reported that the electric field observed by the GEOS-2 satellite
86 showed a transient signature of about 7 mV/m in the dayside magnetosphere associated with the
87 onset of a Sudden Commencement (SC). These signatures were followed by Pc4-5 oscillations.
88 Schmidt and Pedersen (1988) performed a statistical investigation of the GEOS2 electric field
89 signatures associated with SC that showed a clear tailward flow pattern near local noon. Close to
90 the flanks or in the nightside of the magnetosphere the corresponding flows also exhibited a
91 radially inward component. Shinbory et al (2004) investigated the detailed signatures of the
92 Akebono electric and magnetic fields associated with SCs inside the plasmasphere ($L < 5$). The
93 initial excursion of the electric field associated with SCs was almost directed westward at all lo-
94 cal times. The amplitude did not show a clear dependence on magnetic local time and the inten-

95 sity of the E_y field gradually increased by 0.5-2.0 mV/m about 1-2 minutes after the onset of the
96 initial electric field impulse. The propagation velocity of SCs disturbances derived from the am-
97 plitude ratio of the electric field to magnetic field was about 360 km/s in the equatorial plasmas-
98 sphere. Kim et al. (2009) used an MHD simulation to examine the electric field and suggested
99 that the SC associated electric field seen by Shinbory et al. (2004) was the convection electric
100 field. Takahashi et al. (2017) investigated the spatial and temporal evolution of large-scale elec-
101 tric fields in the magnetosphere and ionosphere associated with SCs using multipoint equatorial
102 magnetospheric and ionospheric satellites together with ground radars and showed that the
103 propagation characteristics of electric fields in the equatorial plane depend on magnetic local
104 time. They showed that the initial variation of the electric field (negative E_y) lasted about one
105 minute and was directed westward throughout the inner magnetosphere. Positive E_y became
106 dominant 2 min after SCs propagated to pre-midnight or post-midnight region with near constant
107 amplitude.

108 Observations and MHD simulations (e.g., Li et al., 1993; Zong et al., 2009; Halford et al.,
109 2014; Schiller et al., 2016) show that the electric fields generated by sudden compressions can
110 resonantly interact with trapped charged particle populations within the Earth magnetosphere,
111 energizing and injecting them deep into the magnetosphere. During the well-known shock
112 event in March 1991, the CRESS satellite observed injected electrons energized to extremely
113 high energies, up to 5 MeV (Blake et al., 1992). Wygant et al. (1994) showed that the shock
114 related electric and magnetic field perturbations observed by the CRRES satellite in the
115 nightside inner magnetosphere exhibited a bipolar waveform with amplitude of about 80 mV/m
116 and 140 nT, respectively, and energized the energetic electrons to energies up to 15 MeV. Foster
117 et al. (2015) found that a shock with an azimuthal electric field impulse of 10 mV/m observed
118 by the Van Allen Probes was responsible for accelerating 1.5-4.5 MeV electrons by 400 KeV
119 in the radial region of $L=3.5-4$.

120 This paper focuses on two major issues. We will inspect multispacecraft electric and mag-
121 netic field and particles and plasma observations to study their response to an IP shock on Feb-
122 ruary 27, 2014. We will time the occurrence of magnetic field disturbances associated with the
123 shock in space and the magnetosphere and will show that it propagated downward consistent
124 with expectations based on the shock orientation. Then we will perform a statistical study of the

125 Van Allen Probes electric field disturbances in the magnetosphere and associated plasma drift
126 V_x and V_y velocities in response to IP shocks. We will show that there are four categories of
127 electric field perturbations that occur in response to shock-induced compressions and that these
128 signatures have a clear dependence on magnetic local time. We will show that the direction of
129 the shock normal has an important effect on the propagation of the shock induced magnetic and
130 plasma disturbances and that our statistical results are consistent with MHD simulation predic-
131 tion.

132

133 **2 Data sets**

134 The extensive Van Allen Probes, THEMIS, Cluster and GOES multi-instrument data sets
135 provide numerous opportunities to observe the magnetospheric response to the changes in the
136 solar wind and interplanetary magnetic field monitored by Wind. The five THEMIS spacecraft
137 were launched in 2007 and carry identical instruments and operated in highly elliptical, near-
138 equatorial, orbits that precess about the Earth with apogees of 12, 20, and 30 R_E and orbital peri-
139 ods of 1, 2, and 4 days. With the outermost two spacecraft ARTEMIS now at the Moon, three
140 THEMIS spacecraft remain on the innermost orbits. We use magnetic field data with 3 s time
141 resolution from the THEMIS FGM triaxial fluxgate magnetometers (Auster et al., 2008). The
142 ESA electrostatic analyzer on the THEMIS spacecraft measures the distribution functions of
143 0.005 to 25 keV ions and 0.005 to 30 keV electrons over 4π -str and provides accurate 3 s time
144 resolution plasma moments, pitch angle and gyrophase particle distributions (McFadden et al.,
145 2008).

146 The two Van Allen Probes were launched in August 2012 into nearly identical equatorial
147 and low inclination ($\sim 10^\circ$) orbits with perigee altitudes of 605 and 625 km and apogee altitudes
148 of 30410 and 30540 km (Mauk et al., 2012). Both satellites carry identical sets of instruments to
149 measure charged particle populations, fields, and waves in the inner magnetosphere. In this pa-
150 per, we employ observations from the Energetic Particle, Composition, and Thermal Plasma
151 Suite (ECT: MagEIS, 20-4000 keV for electrons) (Spence et al., 2013; Blake et al., 2013). Elec-
152 tric and Magnetic Field Instrument Suite and Integrated Science (EMFISIS) (Kletzing et al.,
153 2013), and the Electric Field and Waves Suite (EFW) (Wygant et al., 2013). In particular, we
154 inspect electric and magnetic field observations with 11 and 4 s time resolution, respectively,

155 and differential particle flux measurements with ~ 11 s (spin period) time resolution. The electric
156 field data were obtained from sites <http://www.space.umn.edu/rbspew-data> and CDAWEB
157 where they are presented in an MGSE (modified GSE) coordinate system. They provide two
158 components Y and Z of the electric field. Both components are in the spin plane of the spacecraft
159 and are measured with the 50 m long booms. The spin axis X is oriented within 37 degrees of
160 the Earth-Sun line. The spin axis component of the electric field can be obtained from the $\mathbf{E} \cdot \mathbf{B} = 0$
161 assumption. For this to succeed the magnetic field should be at least 15 degrees out of the
162 spin plane. To calculate Van Allen Probes plasma flow velocities we converted the electric field
163 data from modified MGSE coordinates into GSE coordinates. Additionally we used magnetic
164 field data from GOES 13 and 15 with 0.5 s time resolution (Singer et al., 1966) and Cluster with
165 4 s time resolution (Balogh et al., 1997). We use Wind solar wind magnetic field and SWE
166 plasma data with 3 s (Lepping et al., 1995) and 1 min, respectively (Ogilvie et al., 1995).

167 **3 Observations**

168 Figure 1 presents Wind magnetic and plasma data from 15:30 to 16:10 UT on February 27,
169 2014. The arrival of the shock at Wind at 15:50 UT (X, Y, Z GSM = (220.9, 93.9, 30.7 Re)) is
170 revealed by an enhancement in the interplanetary magnetic field strength from 6 to 16 nT and
171 total plasma velocity from 350 to 420 km/s. The IMF had positive B_x and negative B_y compo-
172 nents during the whole interval that both increased the shock arrived. The solar wind density
173 increased from 18 to 45 cm⁻³, and the dynamic pressure increased from 3 to 13 nPa. This fast
174 forward (FF) shock was oblique. Its normal was calculated using magnetic field coplanarity and
175 pointed in the GSM $[n_x, n_y, n_z] = [-0.8, -0.4, -0.3]$ direction, i.e., antisunward, dawnward, and
176 southward. Consequently the shock should first strike the northern dusk bow shock and magne-
177 topause. i.e., it has a spiral IMF orientation. We will use the direction of the shock normal to in-
178 terpret the timing results for the IP shock arrival observed by THEMIS, GOES, Cluster and the
179 Van Allen Probes spacecraft for this event.

180 Figure 2 shows the GSM locations of The THEMIS, Cluster, Van Allen Probes and GOES
181 spacecraft at $\sim 16:50$ UT (Their coordinates are given in Table 1). All the spacecraft located in
182 the solar wind observed the enhanced magnetic field strength, densities, velocities and tempera-
183 tures associated with the IP shock. The shock induced disturbances were seen just upstream

184 from the bow shock by Cluster 1 and 3, located at high southern postnoon latitudes at 16:48:46
185 UT and 16:48:57 UT, respectively .

186 Figures 3 (a, b) show the THEMIS D and A observations of the magnetic field, plasma and
187 energy spectra of ion fluxes from 16:40 to 17:20 UT. The spacecraft were initially located in the
188 magnetosheath. At 16:49:04 UT the IP shock hit THEMIS D as indicated by enhanced densities,
189 magnetic field strength and velocities. Particles from low to high energies showed the increase
190 of energy and enhanced fluxes. The shock produced compression caused the bow shock to move
191 inward at 16:49:36 UT, past the spacecraft as indicated by the decrease in the magnetic field
192 strength and, decrease in density and temperature and spectra expected for its entry into the solar
193 wind. THEMIS A observed the IP shock at 16:49:12 UT and in about 1 min and 34 s later its
194 magnetic field, density and temperature traces indicate that the bow shock moved inward past
195 THEMIS A.

196 Figures 4 (a, b) show GOES 13 and 15 observations of the magnetic field from 16:40 to
197 17:20 UT. Following the arrival of the transmitted IP shock at GOES 13 near local noon at
198 16:50:07 UT there was a sharp increase of magnetic field variations with amplitudes of ~ 70 nT
199 in the H component. The shock induced compression was so strong that at 17:02 UT GOES 13
200 briefly entered the sheath. The shock front was then detected at GOES 15 in the morning local
201 hours 33 sec later at 16:50:40 UT, where it caused a gradual increase of the magnetic field am-
202 plitudes by ~ 20 nT followed by compressional pulsations that fall in the category of Pc5 pulsa-
203 tions.

204 The upper and middle panels of Fig. 5 (a, b) present the Van Allen Probes A and B mag-
205 netic field and electric observations from 16:40 to 17:20 UT. The arrival of the shock character-
206 ized by a strong (~ 50 nT) increase in the total magnetic field strength and bipolar variations in
207 all three components of the electric field at $\sim 16:50:26$ UT at Probe B and 7 sec later at Probe
208 A. The initial electric field perturbations in the E_y component observed by Van Allen Probes A
209 and B were directed downward with amplitudes of -9.4 and -8.2 mV/m, respectively, but ~ 4
210 minutes later the sense changed direction towards dusk (with amplitudes of 5.3 and 5.8 mV/m).
211 We interpret these variations as due to a compression of the magnetosphere followed by a reflec-
212 tion (Samsonov et al., 2007). The E_z and E_x components show variations with amplitudes that
213 are a factor of 1.5-2 smaller than those of the E_y component. The bipolar electric field wave-

214 forms are followed by geomagnetic pulsations with periods of ~ 330 s that damp within ~ 30
215 min.

216 Figures 6 (a, b) present Van Allen Probes A and B observations of the E_z component of the
217 electric field and pitch angle distributions for electron energies of 31.5, 53.8, 108.3, 183.4, 231.8,
218 and 342 KeV measured by the MagEIS instrument. The electrons exhibit enhanced intensities at
219 all energies but the most intense occur at pitch angles near 90° , immediately after the arrival of
220 the IP shock. Kanekal et al. (2016) suggested that the shock-injection mechanism can be effec-
221 tive for energizing particles over a substantial range of pitch angles. The initial flux enhancement
222 is more pronounced by comparison with the following pulses. One of the interesting feature in
223 Figures 6 (a, b) is that the intensity of electrons in the energy range of 31.5-342 KeV exhibits
224 a regular periodicity with periods corresponding to the ULF electric field waves. The oscilla-
225 tions in electron fluxes are in quadrature with the E_y component. This component is of special
226 interest because some charged particles that drift azimuthally as a consequence of the gradient
227 and curvature drifts in the Earth magnetic field can traverse this electric field acquiring a signifi-
228 cant amount of energy. Figures 7 (a, b, c, d, e, f) present the response of the energetic electrons
229 to the IP shock in the energy range from 31 to 183 keV. Panels d and b show that after the shock
230 arrival the electron population increased, especially for the lower energies. In the electrical field
231 of 15 mV/m electron fluxes increased by factors of 21 and 14 at Van Allen Probes B and A, re-
232 spectively, in less than a drift period (panels e and f). The energetic electron fluxes do not dis-
233 play obvious phase differences across the energies. We interpret these observations as evidence
234 for prompt energization of electrons due to shock induced ULF electric fields with an additional
235 contribution for the initial acceleration from the compressional effect of the shock. It should be
236 noted that electrons can be accelerated most significantly via drift resonance (Southwood and
237 Kivelson, 1981) when resonant particles drift with the same velocity as the wave front. Claudep-
238 ierre et al. (2013) showed Van Allen Probes observations of the energy dependence of the ampli-
239 tude and phase of the electron flux modulations which were consequences of drift resonance be-
240 tween ~ 60 keV electrons and fundamental poloidal Pc5 waves. Hao et al. (2014) presented Van
241 Allen Probes observations of electron injections caused by the IP shock and showed that the in-
242 jected electrons with energies between 150 KeV and 230 KeV were in drift resonance with the
243 excited poloidal ULF waves. Considering the process for energizing drift resonant electrons, the

244 value for the $E \times B$ drift velocities of the particles in the wave fields provides important infor-
245 mation. We calculated the V_x and V_y drift velocities at Van Allen Probes A and B for the inter-
246 val from 16:40 to 17:20 UT and present them in the two bottom panels of Fig. 5 (a, b). The V_x
247 and V_y components associated with the minimum peak of the E_y electric field are about -40
248 km/s and -15 km/s for Van Allen Probe B and -35 km/s and -6 km/s for Van Allen Probe A, i. e.,
249 the initial direction of the plasma flow is tailward and dawnward consistent with expectation for
250 the spiral orientation of the IP shock.

251 Interaction with the initial fast mode pulse and subsequent ULF electrical field oscillations
252 can have an important effect on particle acceleration. In considering the energization of elec-
253 trons on February 27, 2014, an encounter with the observed electric field for a period of 240 s
254 will transport the electrons earthward by $\delta R_e = 1.3$ to $1.6 R_e$ from their original position at $L =$
255 6.4 for Van Allen Probe A and at $L = 7.1$ for Van Allen Probe B. Conservation of the first adia-
256 batic invariant implies that such particles will be energized by a factor of about $1.9 - 2.3$ in only
257 one cycle of the electric field pulsations. The studies of Wygant et al. (1994) using CRRES data
258 and Foster et al. (2015) using Van Allen Probes data, and others have demonstrated that the tail-
259 ward propagation of the strong shock-induced electric field impulse and subsequent ULF pro-
260 cesses can result in the extremely fast acceleration of relativistic electron populations inside the
261 plasmasphere.

274 Knowing the distances between the satellites and the lag times for the propagation of shock
275 induced disturbances we calculated the shock propagation velocities. Table 1 summarizes the
276 onset times of the shock driven encounters at different spacecraft. In the solar wind Cluster 1
277 observed the shock earlier than Cluster 3, respectively, that is the shock moved dawnward. The
278 shock perturbations occurred almost simultaneously in the magnetosheath at Themis A and D
279 ($\Delta t < 10$ s) suggesting the front strikes a broad region of the magnetopause at once. The
280 shock induced impulse propagated antisunward, southward and both dawnward and, presuma-
281 bly, duskward (thick arrows in Fig. 2) from the point of origin on the magnetopause (depicted as
282 a red oval in Fig. 2) that is consistent with the orientation of the IP shock. In the outer magne-
283 tosphere the propagation velocity for the disturbance was about 1348 km/s between Goes 13
284 and 15 but only about 390 km/s between Van Allen Probes B and A. We believe that the
285 shock induced pulse propagated with the velocity of fast mode waves. The local fast mode speed

286 can be evaluated from Van Allen Probe measurements of the magnetic field and density. At the
287 time of the shock encounter Van Allen Probes A and B were in the high-density plasmasphere at
288 $L = 5.5$ and $L = 5.1$, respectively. For a measured local magnetic field of 255 nT for Probe A
289 and 220 nT for Probe B and density of $\sim 200 \text{ cm}^{-3}$ derived from the potential of both space-
290 craft, the fast-mode speeds will be $\sim 395 \text{ km/s}$ and 337 km/s , respectively, which are con-
291 sistent with our estimates of the propagation velocity derived from the time difference of shock
292 arrivals at the spacecraft. The decrease of the fast mode wave speed in the plasmasphere relative
293 to that in the outer magnetosphere agrees well with earlier studies (e.g., Wilken, 1982; Foster et
294 al., 2016).

295

296 **4 Statistical study of shock-initiated signatures of the electric field**

297 The list of IP shocks used in this study was obtained from Heliospheric Shock Database
298 maintained and generated by the University of Helsinki [<http://ipshocks.fi>]. They identify shocks
299 by visual inspection and an automated shock detection algorithm. To be included in the data-
300 base a shock should satisfy the following upstream to downstream jump conditions: $B_{\text{down}}/B_{\text{up}}$
301 > 1.2 , $N_{\text{down}}/N_{\text{up}} > 1.2$, $T_{\text{down}}/T_{\text{up}} > 1/1.2$, for FF $V_{\text{up}} - V_{\text{down}} > 20 \text{ km/s}$. The normal vec-
302 tor of the shock (\mathbf{n}) was calculated from the magnetic field data and velocities using the mixed
303 mode method (Abraham-Shrauner and Yun, 1976). When there is data gap in the velocity com-
304 ponents the normal was calculated using magnetic field coplanarity (Colburn and Sonett, 1966).

305 In view of the importance of the electric field in energizing particles we performed a statis-
306 tical study of E_y variations of the electric field and associated plasma drift V_x and V_y velocities
307 during the passage of interplanetary shocks. We identified more than 60 events observed by Van
308 Allen Probes A and B associated with FF IP shocks for the period from 2013 to 2015. The
309 shocks arrived from Wind with lag times in the time range from 26 min to 58 min and pro-
310 duced magnetic field perturbations in the magnetosphere from several to 130 nT. Discontinui-
311 ties in the solar wind plasma such as shocks have often been considered as possible triggers for
312 the release of energy stored within the magnetotail in the form of magnetospheric substorms.
313 Most previous studies of shocks leading to substorms have relied on ground magnetometer ob-
314 servations. Recently it has been shown that the use of global auroral images to identify substorm
315 onsets has some advantages over many other alternative substorm onset signatures, such as low-

316 latitude Pi2 pulsations, auroral kilometric radiation (AKR), and dispersionless particle injections
317 and magnetic field dipolarization at geosynchronous orbits (e.g., Liou et al., 2000). To identify
318 substorms triggered by shocks in our study we considered negative magnetic bays by examining
319 the westward auroral electrojet AL index at the times when SSC were determined from low-
320 latitude ground magnetograms.. As a quantitative definition for the substorm bay does not exist
321 we used the criteria of Liou et el. (2003) that there should be a sharp decrease in AL of at least
322 100 nT occurring within a 20 min window starting at the SSC. We found that shocks triggered
323 a substorm in the magnetosphere in 17 of the 30 examined events. Further study whether
324 these negative magnetic bays are unique identifiers of substorms is beyond the scope of the pa-
325 per. Other effects in the magnetosphere initiated by IP shocks are perturbations in the electric
326 field (Wygant et al., 1994) and the radiation belt (Blake et al., 2015). Understanding and pre-
327 dicting such responses is important for reducing the risks associated with space exploration. We
328 found that 55 events showed an electron enhancement at energies of 32-54 keV measured by
329 MagEIS at all local time and three of them were accompanied by intensity decreases at high-
330 er energies. Five events showed a decrease of the 32-54 keV energy electrons observed in
331 the nightside magnetosphere.

332 The passage of a shock causes electric field perturbations and their amplitudes to increase
333 in proportion to the intensity of the IP shocks. The E field vectors prior to each compression dif-
334 fer greatly from those during the compressional activity. We classified the shock-induced elec-
335 tric field signatures into four different groups according to the examples presented in the upper
336 panels of Fig. 8. Group A presents a negative pulse in the Ey component, B group presents a
337 negative-positive pulse, C group presents a positive pulse and D group presents pulsations.
338 Figure 9 presents occurrence patterns for events with the four different signatures of the
339 electric field initiated by IP shocks. It provides evidence that they are local time dependent.
340 Negative and bipolar pulses predominate on the dayside with positive pulses occur on the
341 nightside. The ULF electric field pulsations of Pc and Pi types produced by IP shocks are ob-
342 served at all local times and in the range of periods from several tens of seconds to several
343 minutes. We believe that the magnetic field as well the electric field pulsations initiated by IP
344 shocks are generated by a wide variety of mechanisms including plasma instabilities, transient
345 reconnection, pressure pulses, and often correspond to field line resonances. Their characteris-

346 tic features are determined to large extent by local time. In the dayside magnetosphere typi-
347 cal pulsations are of the Pc5 type. Sometimes they last for more than twenty wave cycles
348 without noticeable damping which could be explained by a continuous input of the solar wind
349 energy into the magnetosphere. In the nightside magnetosphere during substorms, the generation
350 of Pi2 pulsations is more common. They exhibit an irregular form, last 3-5 wave cycles, and
351 often exhibit damping. Figure 10 presents periods of the pulsations (measured for the first wave
352 cycle of oscillations) as a function of radius and shows that periods increase with increasing
353 radius. A simple explanation for this behavior of pulsation frequencies with radial distance can
354 be given in terms of standing Alfvén waves along resonant field lines (Sugiura and Wilson,
355 1964). The length of the field lines, the magnetic field strength, and the plasma density distribu-
356 tion determine the Alfvén velocity, and the periods of the pulsations. This plot indicates that
357 most electric field pulsations of the Pc5 type in the dayside magnetosphere at $L < 6$ are pro-
358 duced by field line resonances.

359 Figure 11 presents the amplitudes and direction of the initial E_y response to IP shocks in
360 the X-Y GSM plane. The perturbations are negative (dusk-to-dawn) in the dayside
361 magnetosphere (followed by the positive or oscillatory perturbations) and dominately positive
362 (dawn-to-dusk direction) in the nightside magnetosphere, particularly mostly near the Sun-
363 Earth line within an L-shell range from 2.5 to 5. The typical observed amplitudes range from
364 0.2 to 6 mV/m but can reach 12 mV during strong magnetic storms. In the nightside
365 magnetosphere the response of E_y is rather weak and its amplitudes do not exceed 3 mV/m.
366 To demonstrate the impact of IP shocks Fig. 12 shows amplitudes of the initial electric field
367 variations (blue and red crosses) as a function of dynamic pressure observed at Wind. The
368 electric field perturbations increase with the solar wind pressure and that the changes are
369 especially marked in the dayside magnetosphere (red points) as this region is more fully
370 exposed to compression than the nightside sector that is shielded from the frontside
371 compression.

372 To determine the V_x direction of the plasma after the impact of IP shocks we used the
373 formula $V = E \times B/B^2$ for the 60 events under the study. Figure 13 presents the amplitudes
374 and direction of the plasma drift velocities V_x that occur in response to IP shocks (in red -
375 sunward and in blue – tailward directions). The direction of the V_x component of plasma flow

376 is in agreement with the direction of the E_y component (except three peculiar events) and is
377 antisunward at all local times except the nightside magnetosphere, where it is sunward near the
378 Sun-Earth line. The tailward velocities are associated with tailward magnetic field line motion
379 in the dayside magnetosphere. Numbers show that the magnitudes of the flow velocities V_x
380 range from 0.2 to 40 km/s and are a factor of 5 to 10 times stronger near noon as they
381 correspond to greater variations of the electric field in this region.

382 Our results are consistent with the results of global 3D MHD code simulation for the geo-
383 synchronous magnetic field response in the nightside magnetosphere to IP shocks by Wang et
384 al. (2010) presented in Fig. 14. The figure shows contours of δB_z and velocity vectors in
385 the equatorial plane (blue regions - B_z negative, red regions - B_z positive). Their model re-
386 vealed that when a IP shock sweeps over the magnetosphere there are mainly two regions in the
387 nightside magnetosphere, a positive response region in B_z caused by the compressive effect of
388 the shock and a negative response region (blue) which is associated with the temporary en-
389 hancement of earthward convection. They believe that the displacement of the nightside magne-
390 topause caused by the IP shock launches a flow in the magnetosphere near the magnetopause
391 that has a significant y-component, and converges toward the X axis. In the vicinity of the Sun-
392 Earth line at $\sim -5, -6 R_E$ the flow diverges, producing both an earthward flow (consistent with
393 the sunward direction of plasma flow in the nightside magnetosphere presented in Fig. 13) and
394 a tailward flows.

395 As the direction of the shock normal should determine the direction of propagation of
396 transient perturbations and expected flow direction in the magnetosphere initiated by an IP shock
397 we calculated plasma drift velocities V_y for 30 events for which the E_x component could be
398 obtained from $E \cdot B = 0$. We categorized them into two groups for spiral and orthospiral ori-
399 entation of the shock normal. Figure 15 presents the amplitudes and direction of the plasma drift
400 velocities V_y observed by Van Allen Probes A and B in response to IP shocks (red- sunward
401 V_x and blue – tailward V_x directions) for spiral and orthospiral orientations of IP shocks. We
402 excluded several events from the list of shocks that lacked well defined shock normal. As an-
403 ticipated, the shock orientation controls the sense of dawn/dusk flows in magnetosphere. The
404 observed directions of velocity V_y predominately agree with those expected for the given shock

405 normal orientation: dawnward for shocks that sweeps dawnward across the magnetosphere,
406 duskward for shocks that sweep duskward.

407 **5 Conclusions**

408 We presented multipoint observations concerning the response of the electric and magnetic
409 fields, plasma and particles in the magnetosphere to an IP shock on February 27, 2014. We used
410 a multi-spacecraft timing method to determine the propagation speed and direction of the wave
411 front induced by the IP shock. The propagation velocity of the disturbances was about 1348
412 km/s between Goes 13 and 15 in the outer magnetosphere, but it was only about 390 km/s be-
413 tween Van Allen Probes B and A in the inner magnetosphere consistent with expectations for a
414 plasmasphere with limited radial extent. We deduced that the initial encounter of the IP shock
415 with the magnetopause occurred on the post-noon magnetosphere and the shock induced im-
416 pulse propagated as a fast mode wave both dawnward and, presumably, duskward from the
417 point of origin consistent with the spiral orientation of the IP shock. The multipoint measure-
418 ments provide evidence for a dusk-dawn oscillatory electrical field in the dayside equatorial
419 magnetosphere with a peak-to-peak amplitude of ~ 15 mV/m for a period of 30 min. Both
420 spacecraft observed enhanced fluxes of energetic electrons in the range of energies from 31.5 -
421 342 KeV and their intensity shows a regular periodicity with periods corresponding to the elec-
422 tric field pulsations. We interpret these observations as evidence for prompt energization of elec-
423 trons due to shock induced ULF electric fields with an additional contribution for the initial ac-
424 celeration from the compressional effect of the shock. An encounter with the observed electric
425 field for a period of 240 s will transport the electrons earthward by $\delta Re = 1.3$ to $1.6 Re$ from their
426 original positions at $L = 6.4$ for Van Allen Probe A and at $L = 7.1$ at Van Allen B. Conservation
427 of the first adiabatic invariant implies that such a particle will be energized by a factor of about
428 1.9 - 2.3 in only one cycle of the electric field pulsations. The initial plasma flow velocity in the
429 magnetosphere was directed tailward and dawnward, consistent with expectation for the spiral
430 orientation of the IP shock.

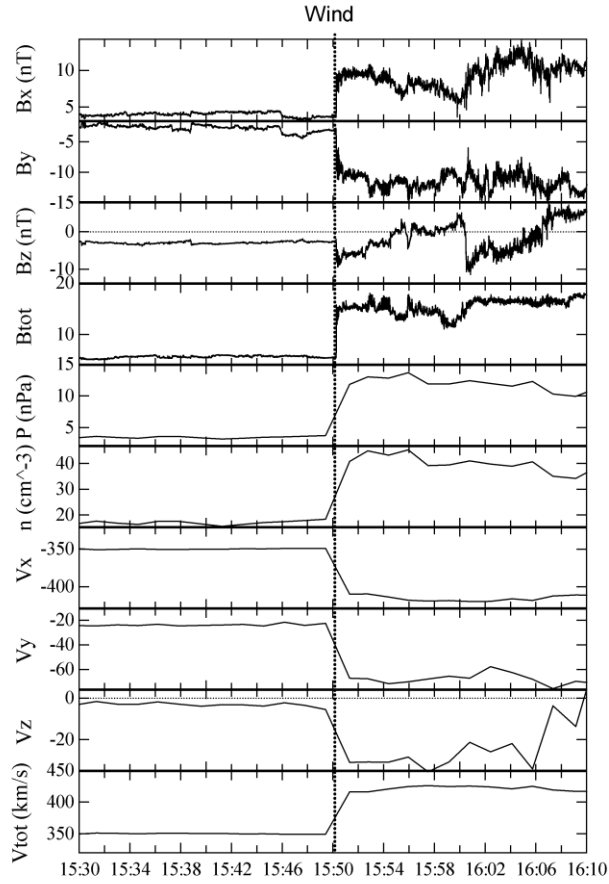
431 We identified more than 60 events observed by Van Allen Probes A and B associated with
432 FF IP shocks for the period from 2013 to 2015. The shocks arrived from Wind with lag times
433 in the time range from 26 min to 58 min and produced magnetic field perturbations in the
434 magnetosphere from several to 130 nT. We found that shocks triggered a substorm in the mag-

435 netosphere in 17 of the 30 examined events. Taking advantage of the multipoint Van Allen
436 Probes observations, we performed a statistical study of E_y variations of the electric field and
437 associated plasma drift V_x and V_y flow velocities during the passage of interplanetary shocks.
438 The E_y perturbations are negative (dusk-to-dawn) in the dayside magnetosphere (followed by
439 positive or oscillatory perturbations) and dominately positive (dawn-to-dusk direction) in the
440 nightside magnetosphere, particularly near the Sun-Earth line within an L-shell range from 2.5
441 to 5. The typical observed amplitudes range from 0.2 to 6 mV/m but can reach 12 mV during
442 strong magnetic storms. We showed that electric field perturbations increase with solar wind
443 pressure and that the changes are especially marked in the dayside magnetosphere. The
444 direction of the V_x component of plasma flow is in agreement with the direction of the E_y
445 component and is antisunward at all local times except the nightside magnetosphere, where it is
446 sunward near the Sun-Earth line but antisunward towards dawn and dusk. The flow velocities V_x
447 range from 0.2 to 40 km/s and are a factor of 5 to 10 times stronger near noon as they
448 correspond to greater variations of the electric field in this region. We investigated how the
449 electric field perturbations deviate from the preceding undisturbed period and demonstrated that
450 the shock-induced electric field signatures can be classified into four different groups according
451 to the initial E_y electric field response. These signatures are local time dependent. Negative and
452 bipolar pulses predominate on the dayside with positive pulses occur on the nightside. The ULF
453 electric field pulsations of Pc and Pi types produced by IP shocks are observed at all local times
454 and in the range of periods from several tens of seconds to several minutes. We believe that
455 that most electric field pulsations of the Pc5 type in the dayside magnetosphere at $L < 6$ are
456 produced by field line resonances. One of the most important results from the present study is
457 that the direction of the shock normal determines the direction of the propagation of the shock
458 induced magnetic and plasma disturbances. The observed directions of velocity V_y predomi-
459 nately agree with those expected for the given spiral or orthospiral shock normal orientation.
460 Our results are consistent with the results of global MHD code simulation of the geosynchro-
461 nous nightside magnetic field response to IP shock by Wang et al. (2010).

462 Table 1. Times of encounter of the IP shock with the spacecraft and their locations in GSM
463 coordinates

S/C	Time	Position	GSM [X, Y, Z] Re	
Wind	15:50:12	220.90	93.92	31.49
Cluster 1	16:48:46	13.10	7.82	-9.44
Cluster 3	16:48:57	12.60	7.73	-10.16
THEMIS D	16:49:04	11.03	0.48	1.10
THEMIS A	16:49:12	9.22	4.39	0.53
Probe A	16:50:33	4.86	-1.69	0.12
Probe B	16:50:26	5.33	-1.39	-0.10
GOES 13	16:50:07	6.51	-0.60	0.99
GOES 15	16:50:40	2.71	-6.02	0.45

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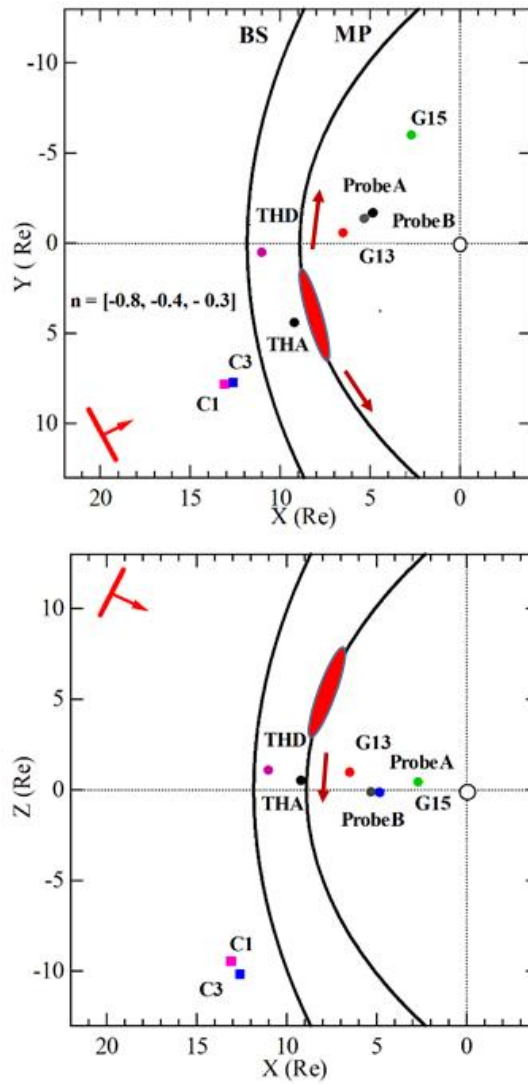
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467 Figure 1. Wind observations of magnetic field and plasma in GSM coordinates from 15:30 UT
 468 to 16:10 UT on February 27, 2014. Dashed line shows the time of arrival of an interplanetary
 469 shock.

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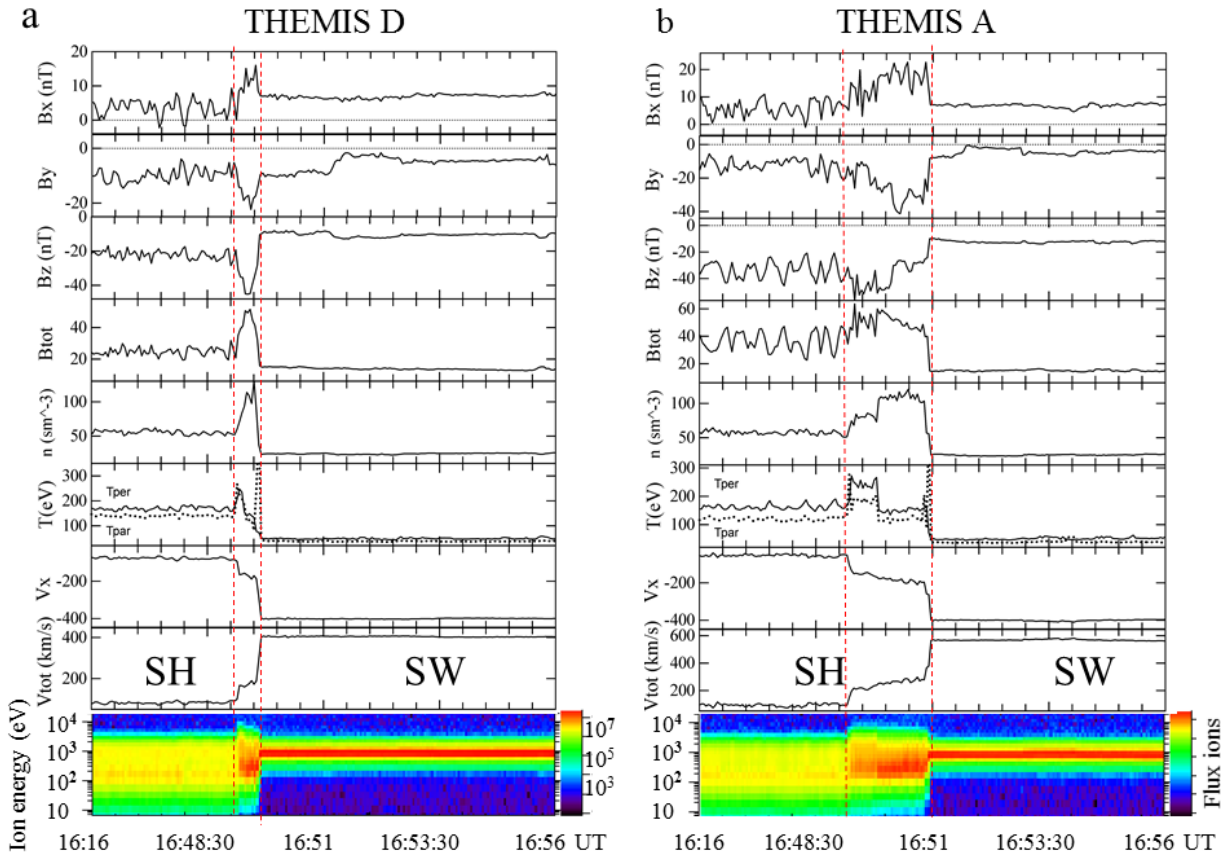
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474 Figure 2. GSM locations of Cluster 1 and 3, THEMIS A, D, Van Allen Probes A and B and
 475 GOES 13 and 15 in the X-Y and Z-Y GSM planes at ~ 1650 UT on February 27, 2014. The
 476 meaning of the solid oval and thick arrows will be discussed in the text later.

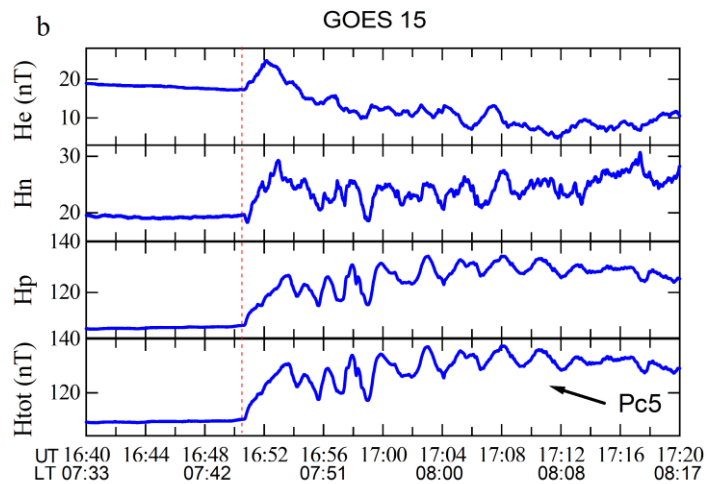
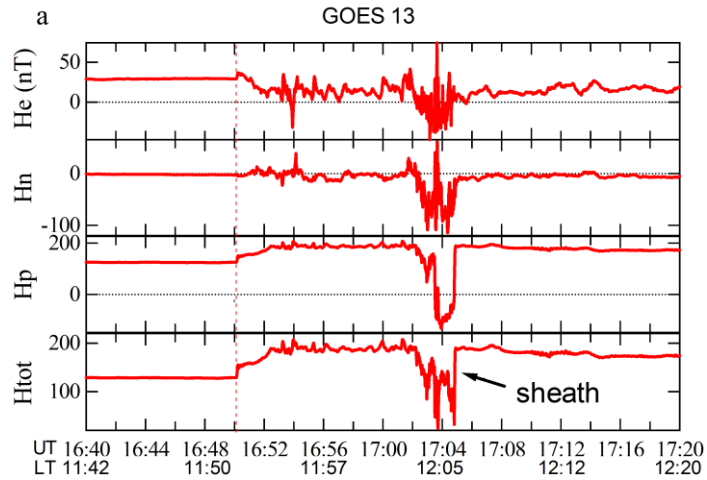


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479 Figures 3 (a, b). THEMIS A (a) and THEMIS D (b) observations of magnetic field in GSM
 480 coordinates plasma and energy spectra of ion fluxes from 16:16 UT to 16:56 UT on February
 481 27, 2014. At 16:49:01 UT the IP shock hit THEMIS D as indicated by enhanced densities, mag-
 482 netic field strength and velocities. Particles from low to high energies showed the increase of
 483 energy and enhanced fluxes. The shock produced compression caused the bow shock to move
 484 inward at 16:49:36 UT, past the spacecraft as indicated by the decrease in the magnetic field
 485 strength and, decrease in density and temperature and spectra expected for its entry into the solar
 486 wind. THEMIS A observed the IP shock at 16:49:12 UT and in about 1 min and 34 s later its
 487 magnetic field, density and temperature traces indicate that the bow shock moved inward past
 488 THEMIS A.

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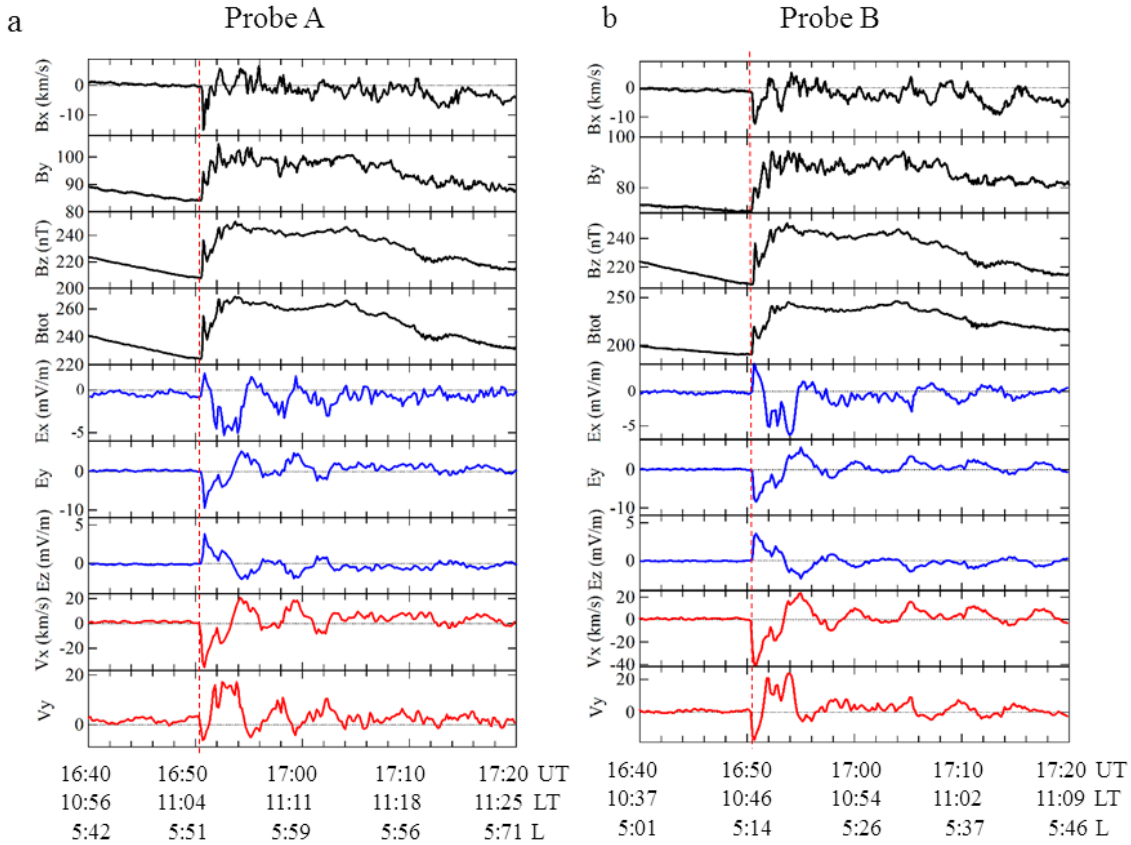
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492 Figures 4 (a, b). GOES 13 (a) and GOES 15 (b) magnetic fields observations in PEN coordinate
 493 from 16:40 UT to 17:20 UT on February 27, 2014. Hp is perpendicular to the satellite's orbital
 494 plane, He pointing earthward parallel to the satellite-Earth center line, and Hn is perpendicular to
 495 both Hp and He and pointing eastward.

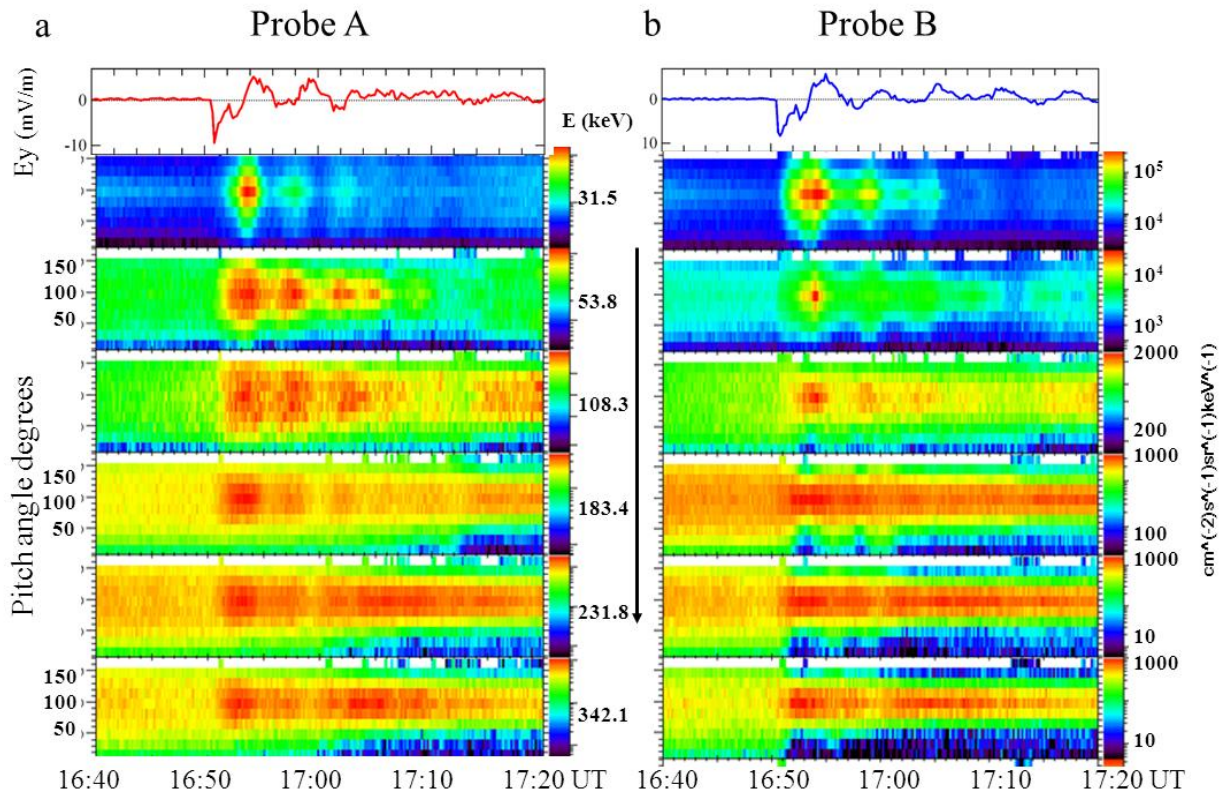
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 498 Figures 5 (a, b). Van Allen Probes A (a) and B (b) magnetic and spin-fit electric field observa-
 499 tions and the V_x and V_y plasma flow velocities in GSE coordinates from 16:40 UT to 17:20
 500 UT on February 27, 2014.

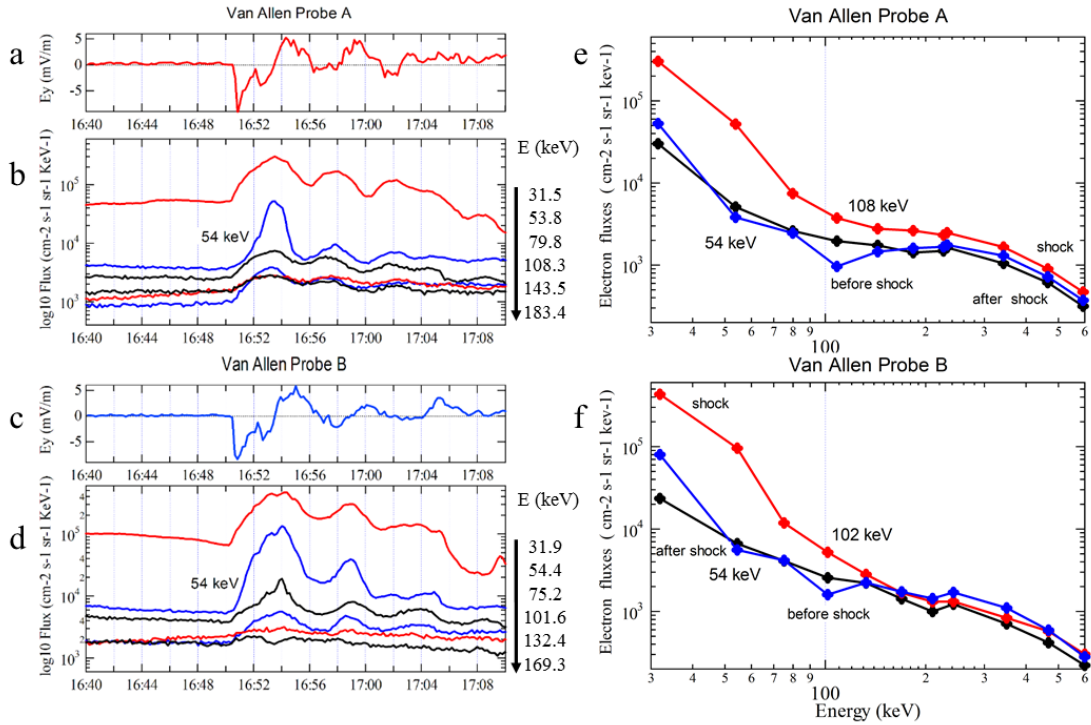
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504 Figures 6 (a, b). Van Allen Probes A (a) and B (b) Ey component of the electric field and pitch
 505 angle distributions of electron fluxes in the range of energies from 31.5 to 342 KeV, measured
 506 by MagEIS instrument from 16:40 UT to 17:20 UT on February 27, 2014. The log fluxes are
 507 color coded according to the color bar shown in the right panel.



508

509 Figures 7. Response of the energetic particles to the transmitted IP shock. Panels a and c show
 510 measurements of the Ey component of the electric field. Panels b and d show electron fluxes for
 511 the energies ranging from 31.5 to 180 keV at Van Allen Probes A and B. Panels e and f show
 512 energetic electron spectra observed at Probe A and B before the shock (at 16:48 UT), after the
 513 first peak (at 16:53 UT) and 18 min after the shock (at 17:08 UT).

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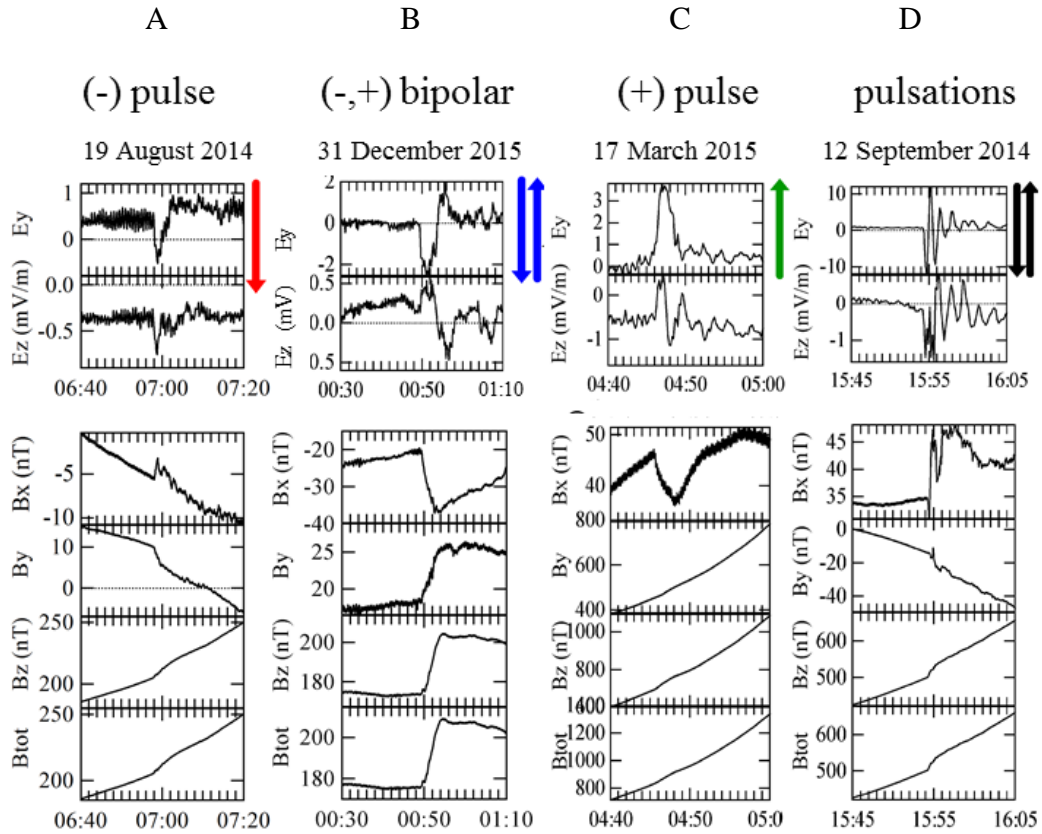
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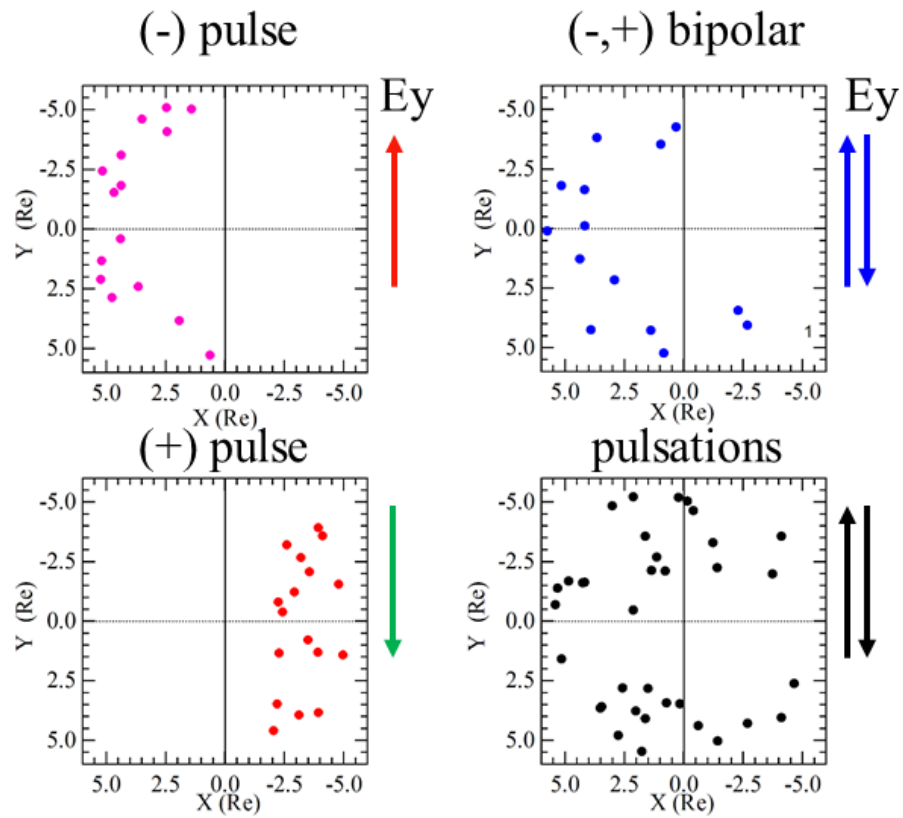
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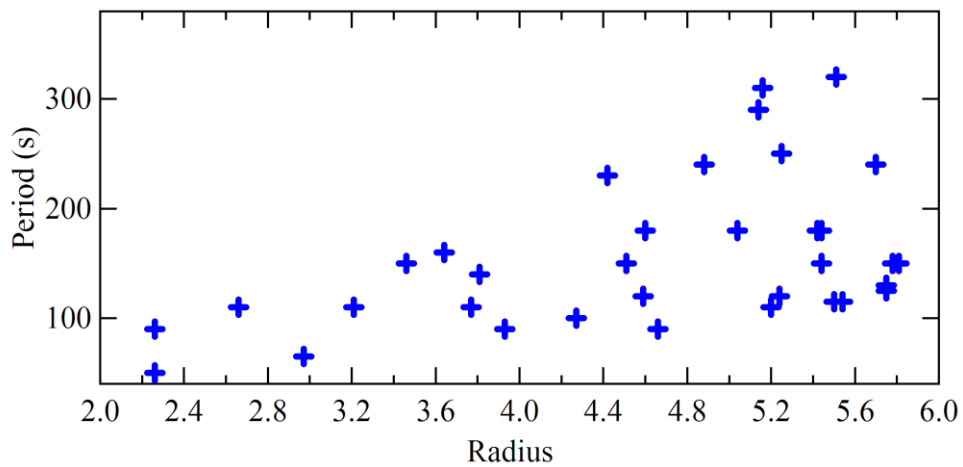
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523 Figure 8. Examples of observed E_y initial variation, including a negative pulse (A), a negative-
524 positive waveform (B), a positive pulse and pulsations (upper panels) and the corresponding
525 magnetic field response (bottom panels).



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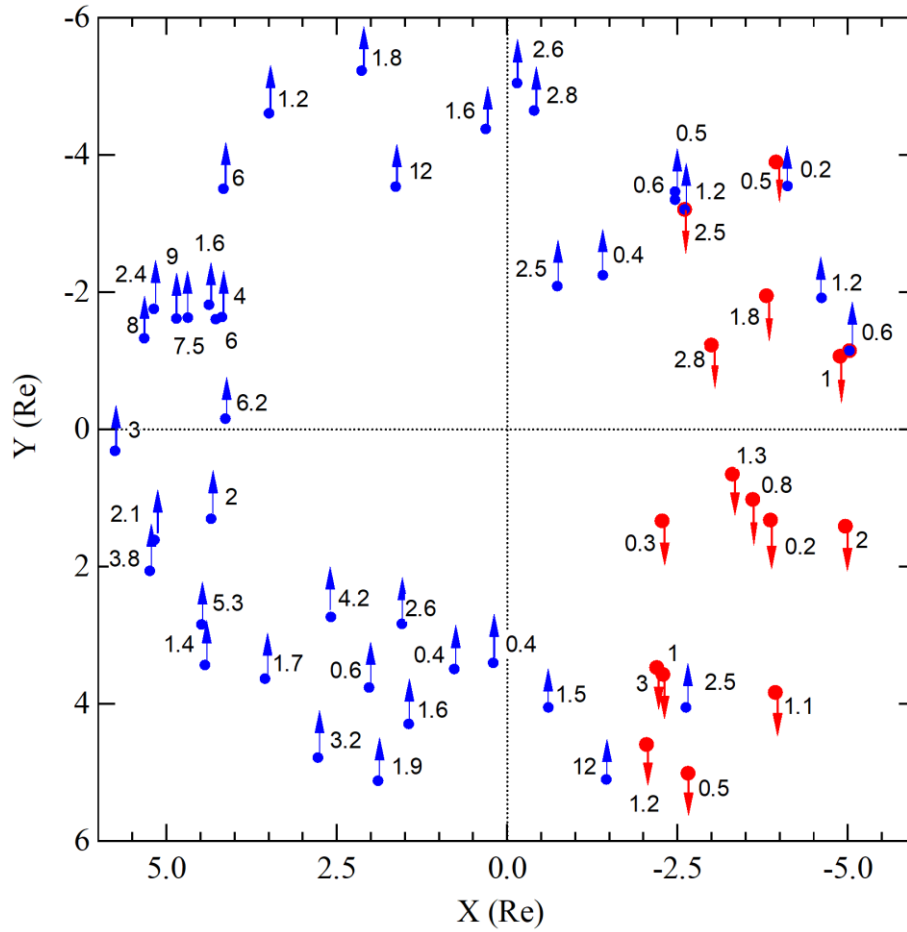
527 Figure 9. GSM locations where event in each of the four groups were observed in the X-Y
 528 plane.



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530 Figure 10. Periods of pulsations, initiated by IP shocks as a function of radius.

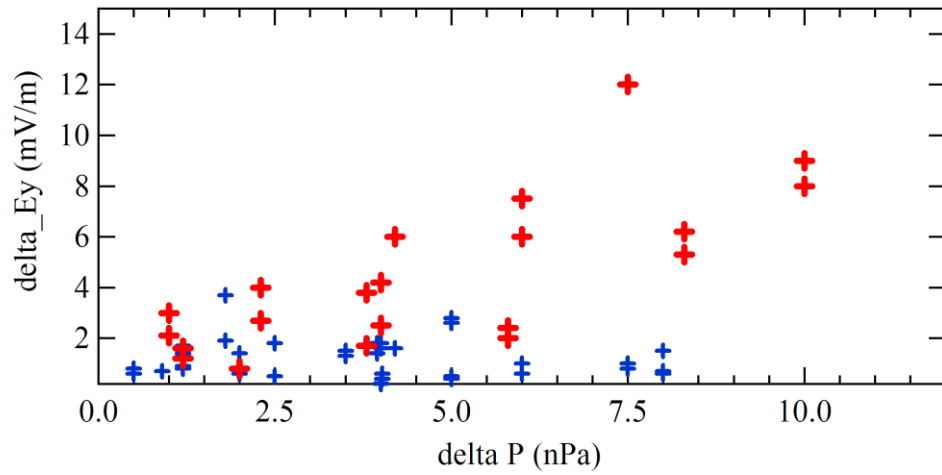
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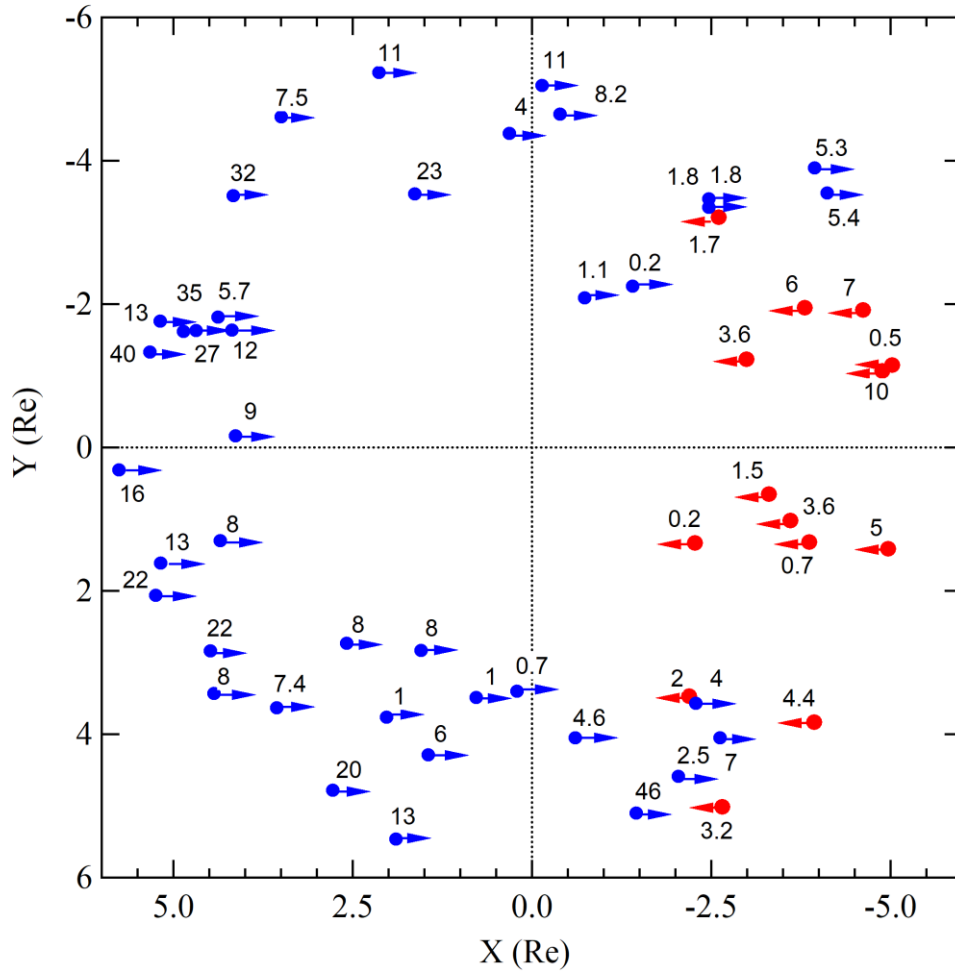
533 Figure 11. Amplitudes and direction of initial E_y response to IP shocks in the X-Y GSM plane
534 (red dawn-duskward and blue duskward-down direction).

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537 Figure 12. Amplitudes of initial Ey variations (blue and red points) caused by a shock as
 538 a function of intensity variations of dynamic pressure observed at Wind for 60 events.



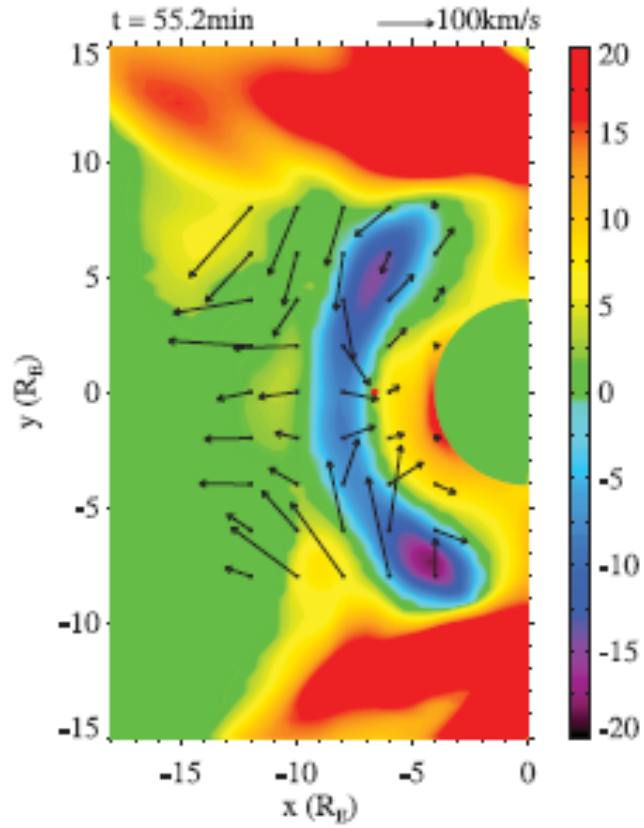
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540 Figure 13. Amplitudes and direction of the plasma drift velocities $V_x = ExB/B^2$ observed by
 541 Van Allen Probes A and B in response to interplanetary shocks (red - sunward and blue -
 542 tailward direction).

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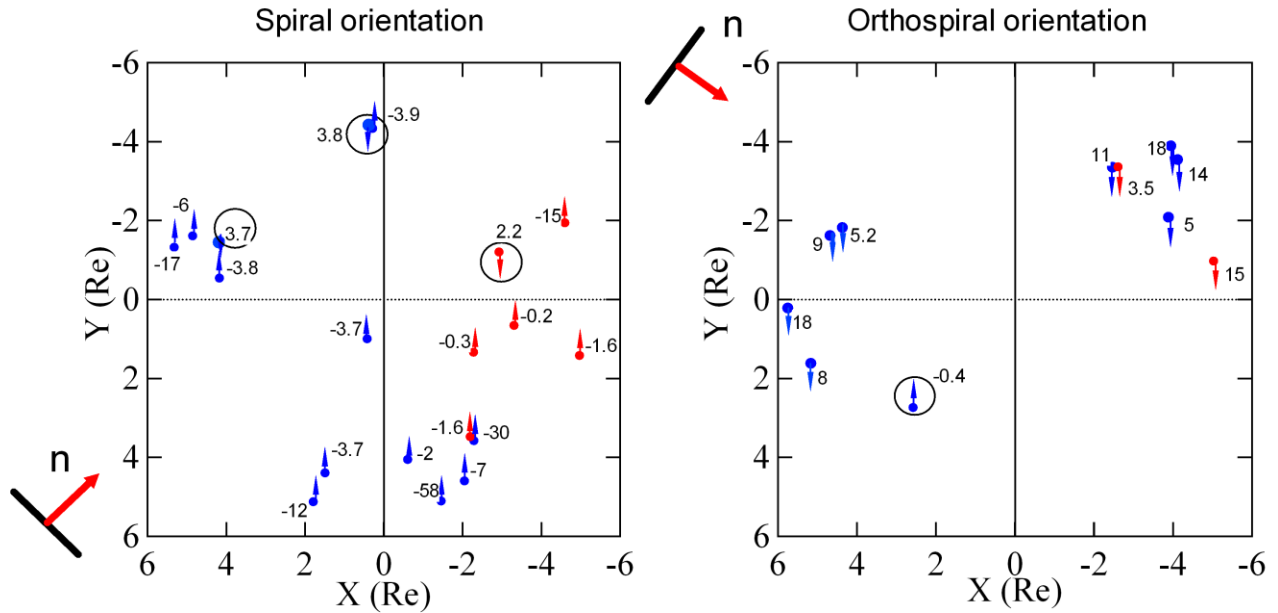
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547 Figure 14. Results of nightside geosynchronous magnetic field response from the global MHD
 548 code simulation of IP shock (Wang et al., 2010). The arrows represent velocity vectors on the
 549 equatorial plane.

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553 Figure 15. Amplitudes and direction of V_y plasma drift velocities observed by Van Allen
 554 Probes A and B in response to interplanetary shocks for spiral and orthospiral orientations (red -
 555 sunward and blue – tailward V_x directions).
 556

557 **6 Data availability.**

558 Data used in the paper are available publicly at http://cdaweb.gsfc.nasa.gov/istp_public/ (Coor-
 559 dinated Data Analysis Web). GOES data can be found at
 560 http://satdat.ngdc.noaa.gov/sem/goes/data/new_full/. The electric field data were obtained from
 561 sites <http://www.space.umn.edu/rbspefw-data>. The list of IP shocks used in this study was ob-
 562 tained from site :<http://ipshocks.fi>.
 563

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565 Galina Korotova, David Sibeck, Scott Thaller, John Wygant, Harlan Spence, Craig Kletzing,
 566 Vassilis Angelopoulos, and Robert Redmon

567 **Author contributions:**

- 568
 569
 570 G. Korotova drafted and wrote the paper and all others commented it.
 571 G. Korotova, D.Sibeck, S. Thaller and R. Redmon were the primary contributors to this work.
 572 S. Thaller- programming, software development, consulting regarding use of Van Allen electric
 573 field data.
 574 J. Wygant -consulting regarding use of Van Allen electric field data.
 575 H. Spence -consulting regarding use of Van Allen plasma data
 576 C. Kletzing -consulting regarding use of Van Allen Magnetometer data
 577 V. Angelopoulos -consulting regarding use of THEMIS data

578
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585 NAS5-01072.

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