



# Nighttime $O(^1D)$ distributions in the mesopause region derived from SABER data

Mikhail Yu. Kulikov<sup>1,2</sup> and Mikhail V. Belikovich<sup>1</sup>

<sup>1</sup>Institute of Applied Physics of the Russian Academy of Sciences, 46 Ulyanov Str., 603950 Nizhny Novgorod, Russia

<sup>2</sup>Lobachevsky State University of Nizhny Novgorod, 23 Gagarin Avenue, 603950 Nizhny Novgorod, Russia

**Correspondence:** Mikhail Yu. Kulikov (mikhail\_kulikov@mail.ru)

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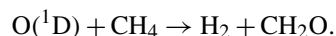
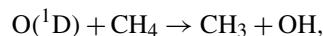
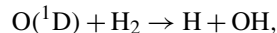
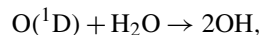
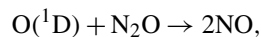
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**Abstract.** In this study, the new source of  $O(^1D)$  in the mesopause region due to the process  $OH(\nu \geq 5) + O(^3P) \rightarrow OH(0 \leq \nu' \leq \nu - 5) + O(^1D)$  is applied to SABER data to estimate the nighttime  $O(^1D)$  distributions for the years 2003–2005. It is found that  $O(^1D)$  evolutions in these years are very similar to each other. Depending on the month, monthly averaged  $O(^1D)$  distributions show two to four maxima with values up to  $340 \text{ cm}^{-3}$  which are localized in height (at  $\sim 92$ – $96 \text{ km}$ ) and latitude (at  $\sim 20$ – $40$  and  $\sim 60$ – $80^\circ \text{ S, N}$ ). Annually averaged distributions in 2003–2005 have one weak maximum at  $\sim 93 \text{ km}$  and  $\sim 65^\circ \text{ S}$  with values of  $150$ – $160 \text{ cm}^{-3}$  and three pronounced maxima (with values up to  $230 \text{ cm}^{-3}$ ) at  $\sim 95 \text{ km}$  and  $\sim 35^\circ \text{ S}$ , at  $\sim 94 \text{ km}$  and  $\sim 40^\circ \text{ N}$  and at  $\sim 93 \text{ km}$  and  $\sim 65$ – $75^\circ \text{ N}$ , correspondingly. In general, there is slightly more  $O(^1D)$  in the Northern Hemisphere than in the Southern Hemisphere. The obtained results are a useful data set for subsequent estimation of nighttime  $O(^1D)$  influence on the chemistry of the mesopause region.

## 1 Introduction

Daytime  $O(^1D)$  is considered to be one of the important chemical minor species of the stratosphere, mesosphere and thermosphere, as it plays a significant role in the chemistry and the radiative and thermal balance of this region (Brasseur and Solomon, 2005). First of all, formed by photolysis of  $O_2$  and  $O_3$ ,  $O(^1D)$  is a mediator involved in the transformation of absorbed solar radiation energy into the heating of this region and, in particular, excitation of  $N_2(\nu)$  and  $CO_2(\nu)$  (Harris and Adams, 1983; Panka et al., 2017). Also,  $O(^1D)$  atoms participate in the reactions of destruction of long-lived green-

house gases (Baasandorj et al., 2012),  $CH_4$  oxidation, and  $HO_x$  and  $NO_x$  production. For example:



Moreover, the red line emission from  $O(^1D)$  atoms is one of the most important airglow phenomena which are used as a diagnostic of the ionosphere, for example, to monitor the electron density and neutral winds in the F region (Shepherd et al., 2019). Therefore, many papers and experimental campaigns are devoted to measurements of features of  $O_3$  photolysis to  $O(^1D)$  (Taniguchi et al., 2003; Hofzumahaus et al., 2004).

Until recently, it was believed that the above-mentioned processes stopped at night as a continuous source of  $O(^1D)$  is absent, while the lifetime of the component is extremely short (less than 1 s). In principle,  $O(^1D)$  can be generated in sprite halos but for a short duration of 1 ms (Hiraki et al., 2004). Recently, Sharma et al. (2015) and Kalogerakis et al. (2016), based on laboratory experiments, proposed that  $O(^1D)$  could be produced in the mesopause region via the process  $OH(\nu \geq 5) + O(^3P) \rightarrow OH(0 \leq \nu' \leq \nu - 5) + O(^1D)$ , which is multiquantum quenching of high excited states of OH by collisions with atomic oxygen in the ground state.

Last year, Kalogerakis (2019) showed that a new model of  $O_2$  A-band that takes this process into account describes well (qualitatively and quantitatively) the results of early nighttime rocket measurements of volume emission rate profiles

of this airglow. Thus, he proved that the process  $\text{OH}(v \geq 5) + \text{O}(^3\text{P}) \rightarrow \text{OH}(0 \leq v' \leq v - 5) + \text{O}(^1\text{D})$  really took place in the nighttime mesopause, and the produced O(<sup>1</sup>D) distributions can be evaluated from available data.

In this study, the new source of O(<sup>1</sup>D) in the mesopause region is applied to SABER data to estimate the O(<sup>1</sup>D) nighttime distributions for the years 2003–2005.

## 2 O(<sup>1</sup>D) derivation from SABER data

All processes used for O(<sup>1</sup>D) determination are summarized in Table 1. Here, we apply the new OH(*v*) model of Fyterer et al. (2019). Their “best-fit model” includes all commonly used production and loss processes of OH(*v*) (see Table 1), but some parameters of the model, in particular, branching ratios of quenching  $\text{OH}(v) + \text{O}_2$  and rate coefficients of  $\text{OH}(v \geq 5) + \text{O}(^3\text{P}) \rightarrow \text{OH}(0 \leq v' \leq v - 5) + \text{O}(^1\text{D})$ , were adjusted with the use of volume emission rate profiles at four different wavelengths measured by SABER and SCIAMACHY.

Due to low values of chemical lifetimes (less than 1 s), O(<sup>1</sup>D) can be considered in chemical equilibrium:

$$\begin{aligned} \text{O}(^1\text{D}) = & (k_9 \cdot \text{OH}(9) \cdot \text{O}(^3\text{P}) + k_{10} \cdot \text{OH}(8) \cdot \text{O}(^3\text{P}) + \\ & k_{11} \cdot \text{OH}(7) \cdot \text{O}(^3\text{P}) + k_{12} \cdot \text{OH}(6) \cdot \text{O}(^3\text{P}) + \\ & k_{13} \cdot \text{OH}(5) \cdot \text{O}(^3\text{P})) / (k_{14} + k_{15} \cdot \text{O}_2 + k_{16} \cdot \text{N}_2) \end{aligned} \quad (1)$$

Thus, to calculate the local value of O(<sup>1</sup>D), we should specify the local concentrations of OH(*v* = 5–9) and O(<sup>3</sup>P). The mentioned model lets us derive the OH(*v*) concentrations as the functions of the OH(*v*) source due to the reaction  $\text{H} + \text{O}_3$  ( $P_{\text{OH}} = k_1 \cdot \text{H} \cdot \text{O}_3$ ), air concentration (*M*), temperature (*T*), and O(<sup>3</sup>P) concentration:

$$\text{OH}(v) = F_v(P_{\text{OH}}, M, T, \text{O}(^3\text{P})). \quad (2)$$

To determine O(<sup>3</sup>P) and  $P_{\text{OH}}$ , we use the known (e.g., Mlynczak et al., 2013, 2018) approach for O(<sup>3</sup>P) derivation from the simultaneous SABER measurements of total volume emission rate of (9–7) and (8–6) OH transitions ( $\text{VER}_{2\mu\text{m}}$ ), O<sub>3</sub> (9.6 μm), and temperature (*T*). The approach employs the chemical equilibrium condition for nighttime ozone. As a result, it is done with the use of the following system of equations:

$$\begin{cases} P_{\text{OH}} = k_2 \cdot \text{O}(^3\text{P}) \cdot \text{O}_2 \cdot M - k_3 \cdot \text{O}(^3\text{P}) \cdot \text{O}_3, \\ \text{VER}_{2\mu\text{m}} = k_4(9, 7) \cdot F_9(P_{\text{OH}}, M, T, \text{O}(^3\text{P})) \\ \quad + k_4(8, 6) \cdot F_8(P_{\text{OH}}, M, T, \text{O}(^3\text{P})). \end{cases} \quad (3)$$

Thus, we derive the local values of O(<sup>3</sup>P),  $P_{\text{OH}}$ , and OH(*v* = 5–9) from SABER data with the use of Eqs. (2)–(3) and apply sets of data (*T*, *M*, OH(*v* = 5–9), and O(<sup>3</sup>P)) to retrieve the local concentrations of O(<sup>1</sup>D) with the use of Eq. (1).

The systematic uncertainty of retrieved data is defined by uncertainties in  $\text{VER}_{2\mu\text{m}}$ , O<sub>3</sub>, *T* measurements, and the rates

of chemical and physical processes included in the OH(*v*) model. We reproduced the analysis presented in Fyterer et al. (2019) (see Sect. 3.4) and took into account the uncertainties of measured data and rate constants which are shown in Table 2. The third column of the table demonstrates the uncertainties’ individual impact at derived O(<sup>1</sup>D) local concentration. It can be noted that the most critical for O(<sup>1</sup>D) are the uncertainties in *T*, rates of Reactions (2)–(3), Einstein coefficients for the *v* = 8–9 states, and  $\text{VER}_{2\mu\text{m}}$ . The total systematic O(<sup>1</sup>D) uncertainty was obtained by calculating the root-sum square of all individual uncertainties. It was found to vary in the range of (37%–52%) depending on the pressure level. Due to averaging, the random error of data presented below is negligible.

## 3 O(<sup>1</sup>D) nighttime distributions

We use version 2.0 of the SABER data product (Level2A) for the simultaneously measured  $\text{VER}_{2\mu\text{m}}$ , O<sub>3</sub>, and *T* profiles within the 0.01–0.0001 hPa pressure (*p*) interval (approximately 80–105 km in 2003–2005). We take only nighttime data when the solar zenith angle is greater than 95°. The range of latitudes covered by the satellite trajectory in a month was divided into 20 bins of ~ (5.5–8)° each; 1500–3000 single profiles of O(<sup>1</sup>D) concentration fall into 1 bin during a month of SABER observations. For each bin we calculate monthly averaged zonal mean  $\langle \text{O}(^1\text{D}) \rangle$  distributions (hereafter, the angle brackets are used to denote timely and spatially averaged values). For annually averaged distributions, we use 40 bins of ~ 4° each.

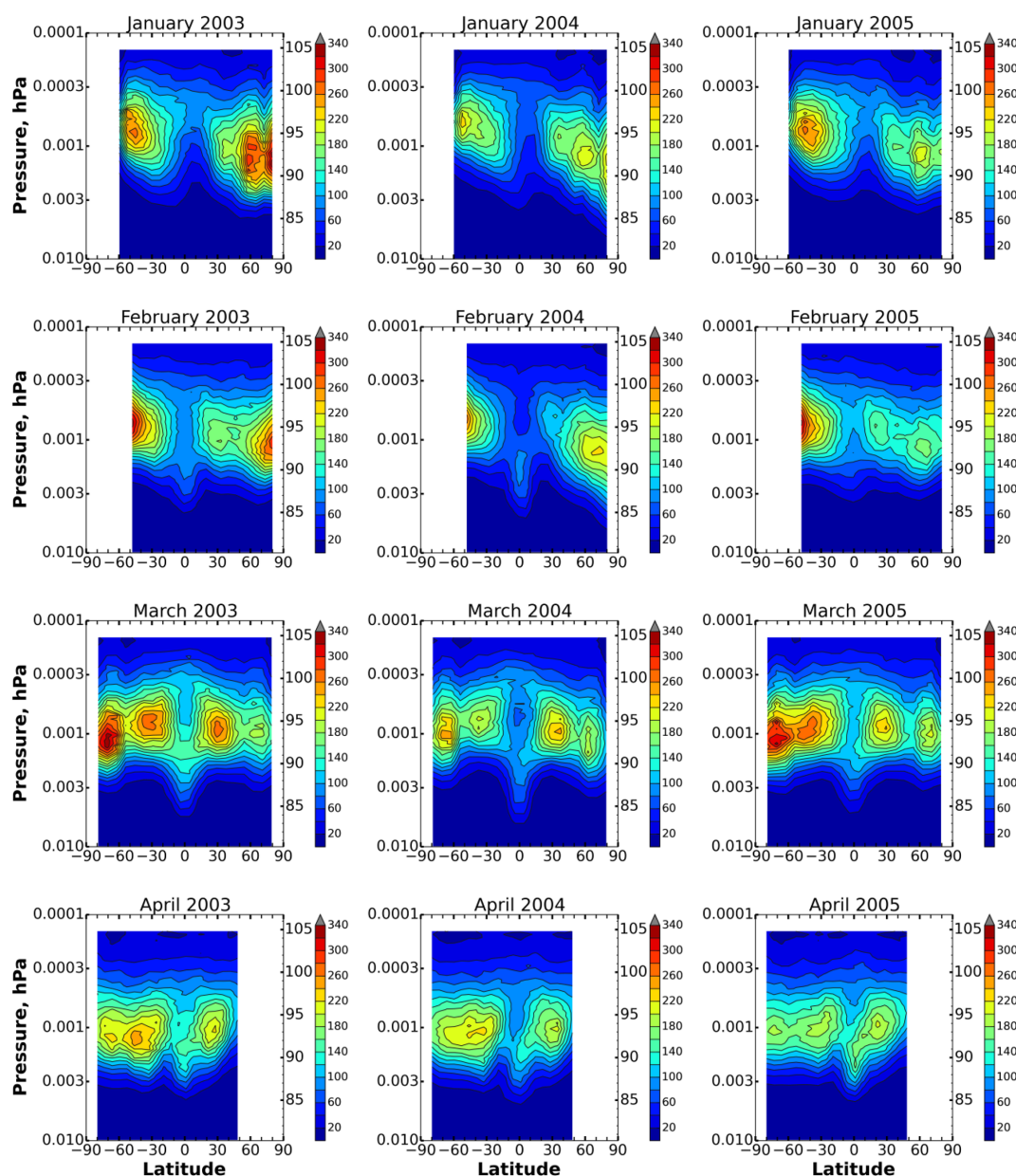
Monthly averaged  $\langle \text{O}(^1\text{D}) \rangle$  distributions in corresponding months of 2003–2005 are shown in Figs. 1–3. Let us analyze the presented data using the distributions in 2003 as an example. Depending on the range of latitudes covered by the satellite trajectory in a specified month, the figures show two to four maxima which are localized in height (at ~ 92–96 km) and latitude (at ~ 20–40 and ~ 60–80° S, N). The values of the maxima can reach up to 300 cm<sup>−3</sup> and more in both hemispheres and different months, for example, in January–March and in May–August. Nevertheless, the annual cycle of southern O(<sup>1</sup>D) demonstrates certain differences from the northern one; i.e., many features of  $\langle \text{O}(^1\text{D}) \rangle$  in the Southern Hemisphere are not repeated in the Northern Hemisphere, with a shift of 6 months. In particular, the distributions in January–February show two pronounced maxima with close values (up to 300 cm<sup>−3</sup>): the first one is at ~ 95 km and ~ 50–60° S, the second one at ~ 93 km and ~ 60–80° N. Half a year later (in July–August), we can see one to two weak maxima in the Southern Hemisphere and a strongly pronounced maximum at ~ 95 km and ~ 40–50° N. A similar pattern can be noticed when comparing the  $\langle \text{O}(^1\text{D}) \rangle$  distributions in June and December. The satellite trajectory in March and September allows us to observe simultaneously four maxima. Note that the southern high-

**Table 1.** List of processes.

	Process	Rate	Reference
1	$\text{H} + \text{O}_3 \rightarrow \text{O}_2 + \text{OH}(v)$	$k_1$ $k_1(v) = k_1 \cdot f(v)$	Burkholder et al. (2015) Alder-Golden (1997, Table 1)
2	$\text{O}(^3\text{P}) + \text{O}_2 + M \rightarrow \text{O}_3 + M$	$k_2$	Burkholder et al. (2015)
3	$\text{O}(^3\text{P}) + \text{O}_3 \rightarrow 2\text{O}_2$	$k_3$	Burkholder et al. (2015)
4	$\text{OH}(v) \rightarrow \text{OH}(v') + h\nu$	$k_4(v, v')$	Xu et al. (2012, Table A1)
5	$\text{OH}(v) + \text{N}_2 \rightarrow \text{OH}(v') + \text{N}_2$	$k_5(v, v')$	Adler-Golden (1997, Table 1), Kalogerakis et al. (2011)
6	$\text{OH}(v) + \text{O}_2 \rightarrow \text{OH}(v') + \text{O}_2$	$k_6(v, v')$	Adler-Golden (1997, Table 3), corrected and adjusted by Fytterer et al. (2019)
7	$\text{OH}(v) + \text{O}(^3\text{P}) \rightarrow \text{H} + \text{O}_2$	$k_7(v)$	Varandas (2004, Table 3, M I)
8	$\text{OH}(v) + \text{O}(^3\text{P}) \rightarrow \text{OH}(v') + \text{O}$	$k_8(v, v')$	Caridade et al. (2013, Table 1)
9	$\text{OH}(9) + \text{O}(^3\text{P}) \rightarrow \text{OH}(3,4) + \text{O}(^1\text{D})$	$k_9$	Fytterer et al. (2019)
10	$\text{OH}(8) + \text{O}(^3\text{P}) \rightarrow \text{OH}(3) + \text{O}(^1\text{D})$	$k_{10}$	Fytterer et al. (2019)
11	$\text{OH}(7) + \text{O}(^3\text{P}) \rightarrow \text{OH}(\leq 2) + \text{O}(^1\text{D})$	$k_{11}$	Fytterer et al. (2019)
12	$\text{OH}(6) + \text{O}(^3\text{P}) \rightarrow \text{OH}(\leq 1) + \text{O}(^1\text{D})$	$k_{12}$	Fytterer et al. (2019)
13	$\text{OH}(5) + \text{O}(^3\text{P}) \rightarrow \text{OH} + \text{O}(^1\text{D})$	$k_{13}$	Fytterer et al. (2019)
14	radiative decay of O( <sup>1</sup> D)	$k_{14}$	Burkholder et al. (2015)
15	$\text{O}(^1\text{D}) + \text{O}_2 \rightarrow \text{O}(^3\text{P}) + \text{O}_2$	$k_{15}$	Burkholder et al. (2015)
16	$\text{O}(^1\text{D}) + \text{N}_2 \rightarrow \text{O} + \text{N}_2$	$k_{16}$	Burkholder et al. (2015)

**Table 2.** List of systematic uncertainties of measured data and rate constants and corresponding uncertainties in derived O(<sup>1</sup>D) local concentration.

Measured characteristic or rate	Its uncertainty	O( <sup>1</sup> D) uncertainty, %
VER <sub>2μm</sub>	6 %	11–17.5
O <sub>3</sub>	10 %	0.1–8.2
<i>T</i>	from García-Comas et al. (2008)	0.1–29.7
<i>k</i> <sub>2</sub>	from Burkholder et al. (2015)	14–30.5
<i>k</i> <sub>3</sub>	from Burkholder et al. (2015)	0.7–19.5
<i>k</i> <sub>15</sub>	from Burkholder et al. (2015)	2.6–2.8
<i>k</i> <sub>16</sub>	from Burkholder et al. (2015)	8–10
<i>f</i> (9)	0.03	1–6
<i>f</i> (8)	0.03	1.4–8
<i>k</i> <sub>4</sub> (9, <i>v</i> ')	30 %	12–23.1
<i>k</i> <sub>4</sub> (8, <i>v</i> ')	30 %	11–24.6
<i>k</i> <sub>4</sub> (7, <i>v</i> ')	30 %	0.6–1.3
<i>k</i> <sub>4</sub> (6, <i>v</i> ')	30 %	0.6–1.3
<i>k</i> <sub>4</sub> (5, <i>v</i> ')	30 %	0.3–0.9

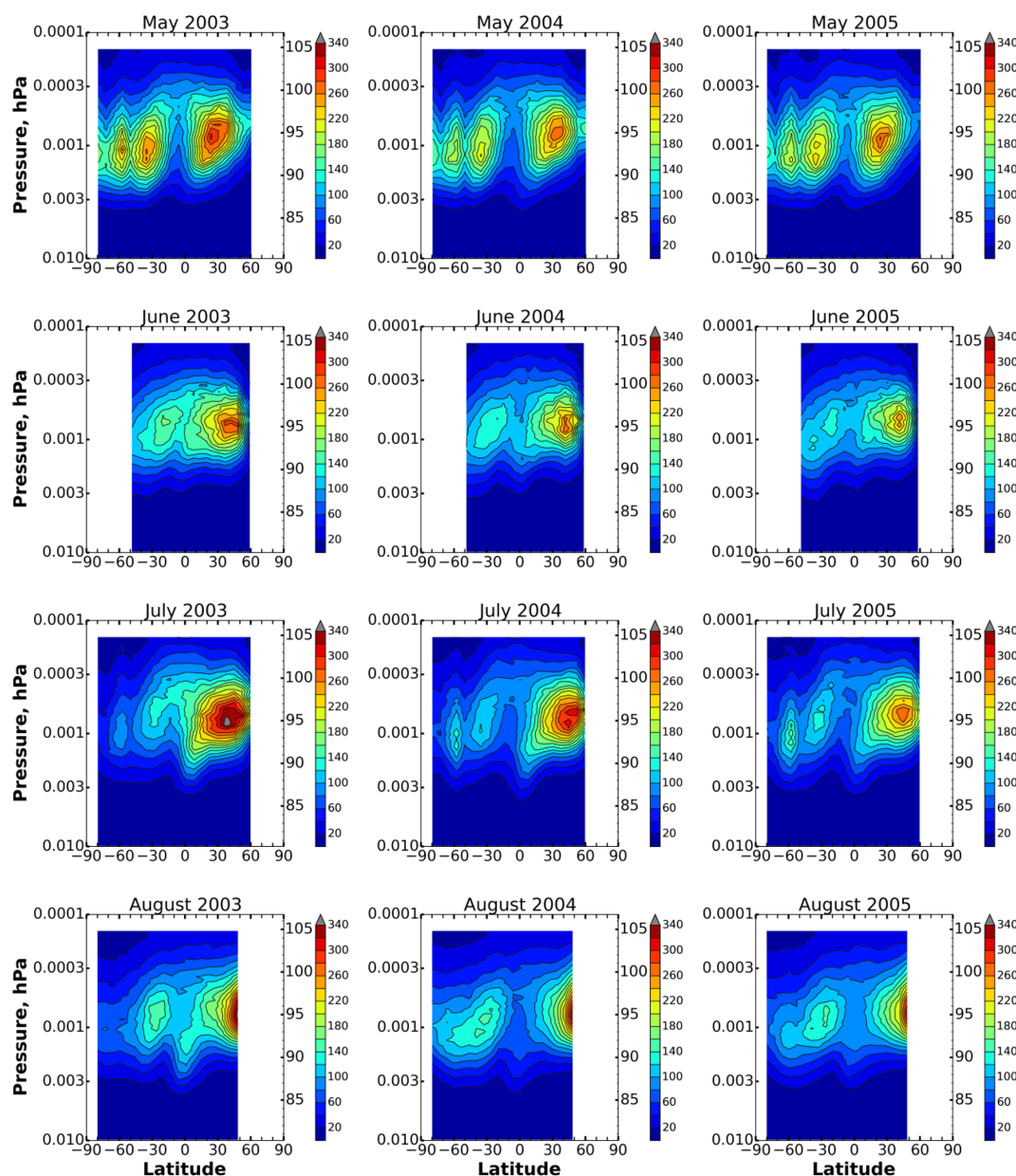


**Figure 1.** Monthly averaged  $O(^1D)$  concentration (in  $\text{cm}^{-3}$ ) in January–April of 2003–2005.

latitudinal maximum (up to  $340 \text{ cm}^{-3}$ ) in March does not correspond to the relatively weak northern high-latitude maximum in September.

The  $\langle O(^1D) \rangle$  evolutions in 2004–2005 are very similar to 2003. Nevertheless, one can see some differences. First of all, in January–February 2004, there is a pronounced particularity above  $60^\circ \text{N}$  below 90 km which does not appear in 2003 and 2005. Kulikov et al. (2019) found similar features in the latitude dependence of the nighttime ozone chemical equilibrium boundary (the lower boundary of the altitudinal–latitudinal region where this equilibrium is satisfied; Belikovich et al., 2018; Kulikov et al., 2018) in January–March

2004 above  $60^\circ \text{N}$  and connected it with abnormal dynamics of the stratospheric polar vortex during the 2003–2004 Arctic winter. There are also additional features which are present in a specific year but absent in the other 2 years. In particular, the northern high-latitude maximum in January–February 2003 is remarkably higher (by the value) than the ones in January–February 2004–2005. The southern high-latitude maximum (up to  $340 \text{ cm}^{-3}$ ) in March 2003 corresponds to the same maximum in March 2005, but both maxima are remarkably higher than the one in March 2004. The reverse (relative to December 2003 and 2005) ratio can be observed



**Figure 2.** Monthly averaged  $O(^1D)$  concentration (in  $\text{cm}^{-3}$ ) in May–August of 2003–2005.

for the values of southern and northern maxima in December 2004.

Annually averaged  $\langle O(^1D) \rangle$  distributions in 2003–2005 are shown in Fig. 4. One can see one weak maximum at  $\sim 93$  km and  $\sim 65^\circ$  S, with values of  $150\text{--}160\text{ cm}^{-3}$  and three pronounced maxima (with values up to  $230\text{ cm}^{-3}$ ) at  $\sim 95$  km and  $\sim 35^\circ$  S, at  $\sim 94$  km and  $\sim 40^\circ$  N, and at  $\sim 93$  km and  $\sim 65\text{--}75^\circ$  N. In general, there is slightly more  $O(^1D)$  in the Northern Hemisphere than in the Southern Hemisphere.

#### 4 Discussion and conclusion

According to various early papers (Nicolet, 1959; Ghosh and Gupta, 1970; Shimazaki and Laird, 1970; Harris and Adams, 1983), daytime  $O(^1D)$  concentrations at 90–100 km varied in the range of  $(10^2\text{--}10^3)\text{ cm}^{-3}$ . Brasseur and Solomon (2005) published the table (see Table A.6.2.c) where daytime  $O(^1D)$  changed from  $70\text{ cm}^{-3}$  at 90 km to  $140\text{ cm}^{-3}$  at 100 km. The presented results show that monthly and annual mean nighttime  $O(^1D)$  concentrations at these altitudes can reach 300 and  $200\text{ cm}^{-3}$ , respectively. Thus, nighttime concentrations of  $O(^1D)$  are comparable with daytime concentrations of this component and, in principle, can impact noticeably



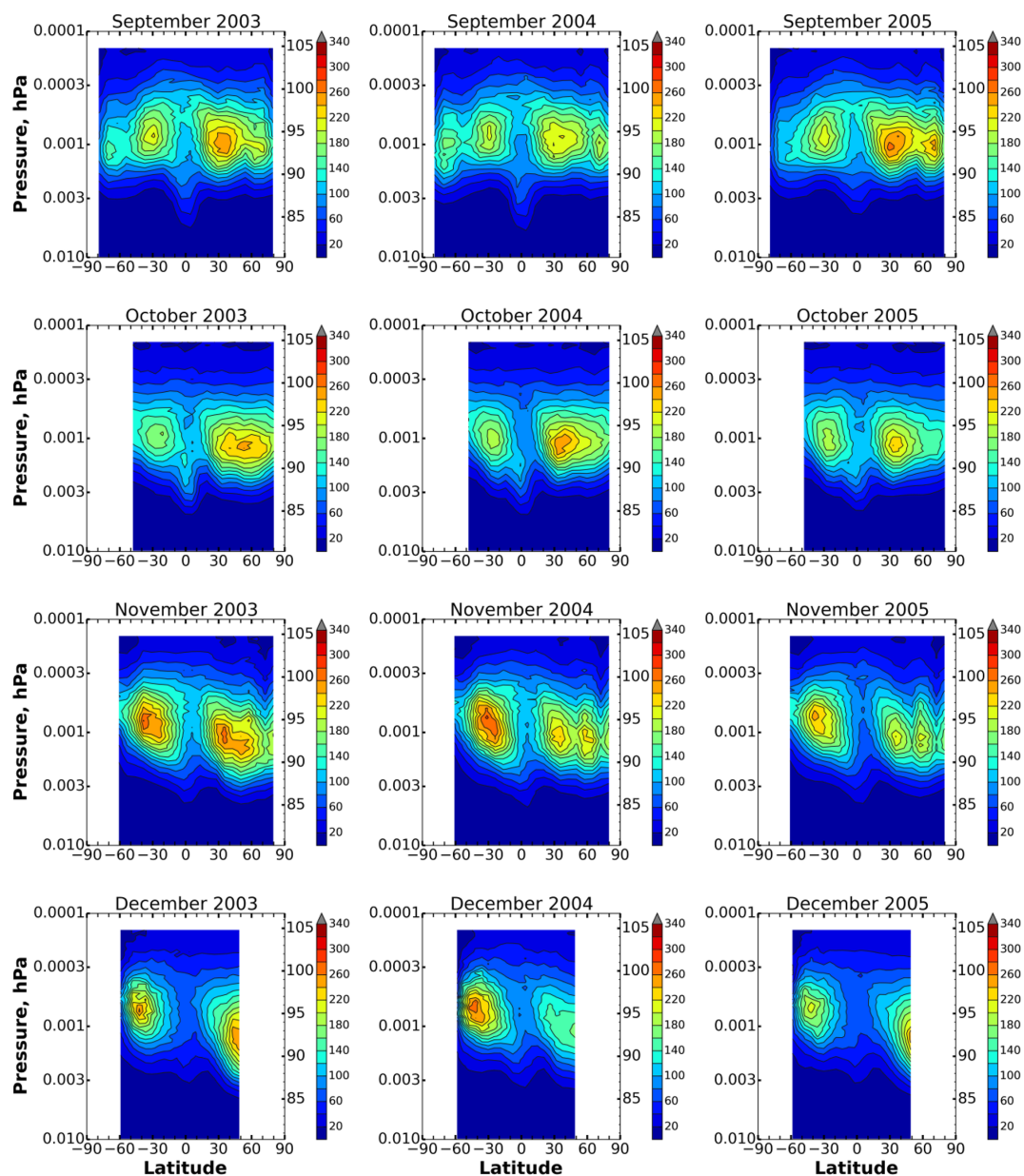


Figure 3. Monthly averaged  $O(^1D)$  concentration (in  $\text{cm}^{-3}$ ) in September–December of 2003–2005.

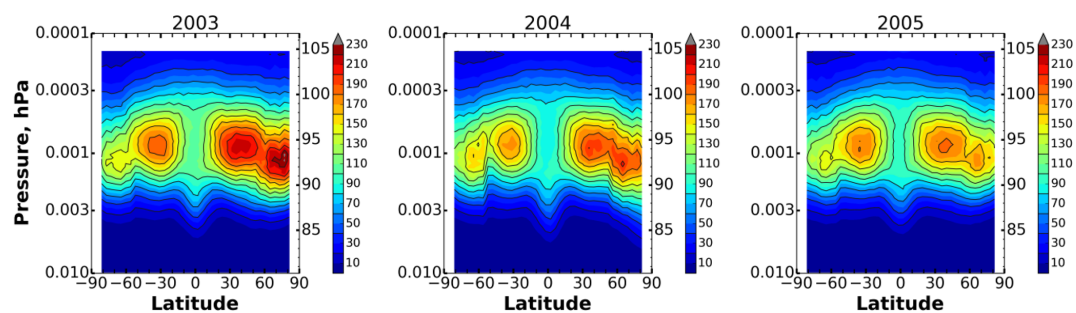


Figure 4. Annually averaged  $O(^1D)$  concentration (in  $\text{cm}^{-3}$ ) in 2003–2005.

the chemistry and thermal balance of the mesopause region. The analysis of this impact should be carried out with the use of a global 3D chemical transport model of the mesosphere–lower thermosphere. Additionally, it may indicate measurable characteristics of this region that could indirectly confirm the results obtained in this article. In principle, direct evidence of O(<sup>1</sup>D) layer existence in the nighttime mesopause can be established by in situ measurements of O(<sup>1</sup>D) airglow at 630 nm, which can be carried out, for example, as a part of a future WADIS rocket sounding mission (Strelnikov et al., 2019; Grygalashvyly et al., 2019). More detailed analysis is out of this short article's scope.

**Data availability.** The SABER data used in this study can be downloaded from [ftp://saber.gats-inc.com/Version2\\_0/Level2A/](ftp://saber.gats-inc.com/Version2_0/Level2A/). The presented data can be downloaded from [https://ipfran.ru/files/10564/article\\_2020\\_o1d\\_data.nc](https://ipfran.ru/files/10564/article_2020_o1d_data.nc) (last access: 2 July 2020).

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**Competing interests.** The authors declare that they have no conflict of interest.

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