

The role of atomic oxygen concentration in the ionization balance of the lower ionosphere during solar proton events

A. Osepian¹, V. Tereschenko¹, P. Dalin², and S. Kirkwood²

¹Polar Geophysical Institute, Chalturuna 15, 183010 Murmansk, Russia ²Swedish Institute of Space Physics, P.O. Box 812, 98128 Kiruna, Sweden

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Abstract. The influence of atomic oxygen concentration on the height distribution of the main positive and negative ions and on electron density in the mesosphere is studied for the conditions prevailing during the solar proton event on 17 January 2005. It is shown by numerical modeling that the electron and ion density profiles are strongly dependent on the choice of the atomic oxygen profile. Experimental measurements of the electron density are used as the criterion for choosing the atomic oxygen profile in the mesosphere. With the help of modeling, the atomic oxygen profile in the daytime in the winter mesosphere is found to lead to a model electron density profile best matching the electron density profile obtained experimentally. As a result, with the help of modeling, we find the atomic oxygen profiles at various solar zenith angles in the winter mesosphere which lead to model electron density profiles matching the electron density profiles obtained experimentally.

Alteration of the atomic oxygen concentration leads to a redistribution of the abundance of both positive and negative ion constituents, with changes in their total concentrations and transition heights. In consequence this results in changes of the electron density and effective recombination coefficient. For conditions of low concentration of atomic oxygen (during a solar proton event), the formation of cluster ions is the key process determining electron and ion densities at altitudes up to 77 km. The complex negative CO_3^- ion is formed up to about 74 km and the final NO_3^- ion, which is stable in relation to the atomic oxygen, is the dominant negative ion up to 74 km. As a result the transition heights between cluster ions and molecular ions and between negative ions and electron density are located at 77 km and 66 km, respectively.

Correspondence to: P. Dalin (pdalin@irf.se)

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1 Introduction

Studies of the lower ionosphere and mesosphere are complicated by a number of problems, both experimental and theoretical. One of the main experimental problems is poor accessibility of in-situ measurements. Present knowledge on ion composition is based on only a limited number of sounding rocket flights.

Electron concentration is mainly measured with different ground-based techniques, although occasional measurements are made using sounding rockets. It is often extremely difficult to measure electron concentration in the mesosphere with a good precision, particularly by ground-based techniques, because of the low values of electron density, especially during nighttime. However, in the presence of extra ionization in the lower ionosphere, produced by precipitating particles or solar X-ray flares, the electron concentration can be measured more easily.

The electron concentration in the polar D-region is extremely variable, especially during geomagnetic disturbances. Moreover, even during quiet periods, electron density profiles measured for similar conditions (for example, the same season, the same zenith angle) can sometimes differ substantially. This means that, not only the changeable ionization sources, but also a set of other factors, such as changes in concentration of minor neutral species and neutral atmosphere temperature, are important for the formation of electron density profiles. The large variability of the height profile of electron concentration, $N_e(h)$, in the polar lower ionosphere is the reason why there is no reliable empirical model of the high-latitude D-region yet available which can be used in the absence of direct measurements. The empirical models developed by Friedrich et al. (2006) and McKinnell and Friedrich (2007), require the availability of riometer measurements, and still cannot predict electron density at 70 km height under all conditions to better than an order of magnitude. A useful theoretical model must be capable of describing the $N_e(h)$ dependence on all the factors influencing the electron density.

Theoretical studies provide major advances for our understanding of the processes controlling the ion composition and electron density in the lower ionosphere. In a theoretical study, the physical and chemical processes, which produce and maintain the lower ionosphere, such as ionization by different sources (solar electromagnetic radiation in different wavelength ranges, galactic cosmic rays and high energy particles), recombination laws, photodissociation and photodetachment, formations of different ions, which are described by ion-chemical reactions, and main ionic transformation processes and their dependencies on neutral temperature and density, humidity and the concentrations of minor neutral species etc., have to be taken into account. Ion composition and ion chemical reactions in the D-region are, for example, directly coupled to the local concentration of longlived species such as O, NO, H₂O.

In this study, based on a theoretical model of the D-region (Smirnova et al., 1988; Kirkwood and Osepian, 1995; Osepian and Smirnova, 1997), we investigate the influence of the atomic oxygen concentration on the ionization balance at mesospheric altitudes during a solar proton event (SPE) that occurred on 17 January 2005.

At the present time there is a great deficiency of experimental knowledge on the atomic oxygen concentration at altitudes of the lower ionosphere. Rocket measurements (Dickinson et al., 1980; Gumbel et al., 1998) are available but they are of sporadic character and therefore do not reflect all the geophysical situations. Due to the poor availability of experimental data on the atomic oxygen concentration [O] below 80 km, theoretical [O](h)-profiles estimated at mesospheric altitudes with different diffusive-photochemical models are used.

Variations between such estimates are rather large. Therefore one of the major requirements throughout the middle atmosphere is information on the [O](h)-profiles under different geophysical situations so that they can be incorporated into theoretical models. In this study we use experimental data on electron density in the lower ionosphere obtained with two ground-based techniques, namely incoherent scatter of radio waves (see e.g. Nygrén, 1996) using the EIS-CAT 224 MHz radar in Tromsø, Norway (geographical coordinates are φ =69.6°, λ =19.3°; geomagnetic coordinates are Φ =66.2°, Λ =105.8°), and the partial reflection of radiowaves using a 2.8 MHz radar in Murmansk, Russia (geographical coordinates are φ =69.0°, λ =33.0°; geomagnetic coordinates are Φ =64.5°, Λ =115.0°). We combine observations with a theoretical ionization-recombination model to determine the correct choice of the atomic oxygen profile in the mesosphere during the SPE in the winter season.

In Sect. 2 we consider available theoretical results on the height distribution of atomic oxygen obtained with different diffusive-photochemical models during both quiet and disturbed periods.

In Sect. 3 we estimate the effective loss coefficients of the primary ions, NO⁺ and O₂⁺, in relation to the formation of hydrated ion clusters as well as the loss rate of the primary ions O₂⁻ in relation to the formation of complex negative ions. We use an ion-chemical model of the D-region and theoretical [O](*h*)-profiles with maximum and minimum values. We also estimate the effective rate coefficients for competing processes affecting NO⁺, O₂⁺ and O₂⁻ loss in reactions with atomic oxygen. Comparison of these effective rate coefficients for a given [O] concentration allows us to understand the importance of reactions with atomic oxygen in the ionization balance at each altitude, and in the formation of electron density and effective recombination profiles.

In Sect. 4 we compare experimental N(h)-profiles observed during SPE with N(h)-profiles calculated using different [O](h)-profiles.

2 Models of atomic oxygen

Solar ultraviolet radiation is chiefly absorbed by molecular oxygen (O₂) which is photodissociated to produce two oxygen atoms (O):

$$O_2 + h\nu \rightarrow O(^{3}P) + O(^{1}D)$$

The thermospheric production of odd oxygen produces a downward flux of atomic oxygen which must be balanced by destructive processes occurring at lower altitudes. Atomic oxygen disappears via the following reactions (Anderson, 1976):

$$O(^{3}P) + O(^{3}P) + M \to O_{2} + M$$
 (1)

$$O(^{3}P) + O_{2} + N_{2} \rightarrow O_{3} + N_{2}$$
 (2)

$$O(^{3}P) + O_{3} \to O_{2} + O_{2}$$
 (3)

$$O(^{1}D) + O_{3} \rightarrow O_{2} + O_{2}$$
 (4)

$$O(^{1}D) + O_{3} \rightarrow O_{2} + O(^{3}P) + O(^{3}P)$$
 (5)

$$O(^{1}D) + N_{2} \rightarrow O(^{3}P) + N_{2}$$
 (6)

where M is an atmospheric molecule, either O_2 or N_2 , acting as a third body.

The peak of [O]-concentration ($\approx 10^{11} - 10^{12} \text{ cm}^{-3}$) is generally located at around 95 km height. Below 90 km the influence of fast reacting hydrogen species becomes significant for odd oxygen loss:

$$H + O_3 \to OH + O_2 \tag{7}$$

 $H + O_2 + M \rightarrow HO_2 + M \tag{8}$

$$OH + O(^{3}P) \rightarrow H + O_{2}$$
(9)

$$\mathrm{HO}_2 + \mathrm{O}(^{3}\mathrm{P}) \to \mathrm{OH} + \mathrm{O}_2 \tag{10}$$

In the upper atmosphere both vertical and horizontal transport of atomic oxygen is important. Since the lifetime of atomic oxygen above the mesopause is large, the average meridional flow transports atomic oxygen from the summer hemisphere to the winter hemisphere, where it is destroyed by recombination. Here species such as O and NO are transported downward. Turbulent diffusion resulting from the breaking of gravity waves acts as a transport mechanism responsible for strong mixing and dissipative processes. Thus the atomic oxygen is strongly controlled by both photochemical and dynamical processes (downward and meridional transports, turbulence, molecular and eddy diffusions) (Shimazaki and Laird, 1972; Keneshea et al., 1972; Thomas and Bowman, 1972; Maharaj-Sharma and Shepherd, 2004; Russell et al., 2004, 2005; Murray and Plane, 2005). These processes are expected to lead to substantial changes in the atomic oxygen concentration. As a result, available diffusive-photochemical models give estimates of the [O]concentration, which differ significantly between the different models.

Examples of theoretical O(h)-profiles (O(h)-profiles 1 and 2) are presented in Fig. 1. These profiles reflect the minimum and maximum values of [O] given by different diffusivephotochemical models for daytime quiet conditions. The O(h)-profile number 1 is taken from the model by Shimazaki and Laird (1972). The O(h)-profile 2 at altitudes less than 80 km is taken from the model by Zadorozhny (1982) and at altitudes above 80 km from the MSIS-2000 neutral atmosphere model (Hedin, 1991). Differences in the concentration of atomic oxygen between O-profiles 1 and 2 are about an order of the magnitude. As has been mentioned above, these differences are due to differences in accounting for dynamical processes, in model values of eddy diffusion coefficient, in concentration of ozone and other species, in solar radiation intensity at different wavelength ranges, in absorption cross sections and in other factors applied in the model studies. For example, according to the model calculations by Keneshea et al. (1972), a change of the turbulent diffusion coefficient by a factor of 3 leads to an order of magnitude change in [O]-concentration.

It is well known that, under conditions of strong and long duration disturbances (SPE), there is both a large increase in ionization and changes of the neutral composition (Porter et



Fig. 1. Atomic oxygen profiles in the mesosphere. 1–5 are model profiles. 1 is from Shimazaki and Laird (1972); 2 is from Zadorozhny (1982); 3 and 4 are assumed for SPE conditions in the present study; 5 is for summer, noctilucent cloud conditions (Murray and Plane, 2005). Profile 6 corresponds to measurements for summer noctilucent cloud conditions (Dickinson et al., 1980).

al., 1976; Solomon et al., 1981, 1983; McPeters and Jackman, 1985; Reid et al., 1991; Zadorozhny et al., 1992; Jackman et al., 1995, 2005; Seppala et al., 2004; Rochen et al., 2005; Verronen et al., 2005; Lopez-Puertas et al., 2005). One of the important effects of particle precipitation events is enhanced concentration of nitric oxide, NO, and of odd hydrogen constituents, H and OH. The latter can arise from Reactions (11–13) and dissociative recombination (14) of water cluster ions H^+ .(H₂O)_n (Frederick, 1976; Porter et al., 1976; Reid et al., 1977; Solomon et al., 1981):

$$O_2^+.H_2O + H_2O \to H_3O^+.OH + O_2$$
 (11)

$$O_2^+.H_2O + H_2O \to H_3O^+ + OH + O_2$$
 (12)

$$H_3O^+.OH + H_2O \to H^+.(H_2O)_2 + OH$$
 (13)

$$\mathrm{H}^{+}(\mathrm{H}_{2}\mathrm{O})_{n} + e \to \mathrm{H} + n(\mathrm{H}_{2}\mathrm{O})$$
(14)

According to photo-chemical models (Swider and Keneshea, 1973; Frederick, 1976; Heath et al., 1977; Solomon et al., 1983; Kozlov et al., 1988; Smirnova et al., 1990; Jackman et al., 1995; Verronen et al., 2005) an increase of H and OH-concentration leads to changes in the oxygen constituents in the mesosphere. The details of the neutral oxygen component variations under the effects of very strong solar proton precipitation depend on height, season, solar illumination and intensity of the ionization source. Theoretical estimates (Swider and Keneshea, 1973; Smirnova et al., 1990) show that the [O]-concentration can decrease in the daytime by a factor of 2–8 at altitudes 50–75 km due to increasing [OH] and [HO₂]-concentrations and increasing rates of Reactions (9) and (10).

In accordance with these theoretical estimates we have decreased the atomic oxygen concentration by factors of 10 and 4 at altitudes of 50–75 km relative to the O-profile 2 deduced for quiet conditions as seen in Fig. 1 (O-profiles 3 and 4). We use these hypothetical O-profiles 3 and 4, together with the O-profiles 1 and 2, as input to a theoretical model for estimation of the effective loss rate coefficients of the primary ions NO⁺, O_2^+ , O_2^- and for calculation of the ion composition, electron density and effective recombination coefficient profiles.

In Fig. 1 we also present two further O-profiles, (5) as modeled by Murray and Plane (2005) and (6) as measured by Dickinson et al. (1980), both for the summer mesosphere. In these cases the observed sharp depletion in O concentration can be explained by an enhancement of H, OH, HO₂, H_2O_2 related to the presence of noctilucent clouds (Murray and Plane, 2005). The source of hydrogen compounds is not the same as during a solar proton event, but the effects are similar to those we have assumed in profiles 3 and 4.

3 Influence of atomic oxygen concentration on ion composition and electron density

To investigate the role of the atomic oxygen concentration in the formation of ion composition and electron density profiles at mesospheric altitudes, we use a theoretical scheme describing the main ionic transformation processes and ionchemical reactions (Smirnova et al., 1988) which forms the basis of the theoretical model applied in the present study. The scheme is shown in the Appendix. Complex transformations and formation of hydrated ion clusters are via the following reaction chains:

$$O_2^+ \to O_4^+ \to O_2^+.(H_2O)_n \to H^+(H_2O)_n$$
 (15a)

$$\mathrm{NO}^+ \to \mathrm{NO}^+.(\mathrm{H}_2\mathrm{O})_n \to \mathrm{H}^+.(\mathrm{H}_2\mathrm{O})_n \tag{15b}$$

and formation of intermediate negative ions is represented by the following chains:

$$O_2^- \rightarrow O_3^-$$
 and $O_4^- \rightarrow CO_4^-$ and $CO_3^- \rightarrow O_2^-$ (15c)

$$O^- \to O_3^- \to CO_3^- \to O_2^- \to O^-$$
(15d)

These are described by a large number of ion-chemical reactions in detailed reaction schemes. In the chemical model developed by Smirnova et al. (1988) the efficiencies of reactions included in the channels (15a–d) are substituted by effective parameters: $B(O_2^+)$, $B(NO^+)$, $B(O_2^-)$, $B(O^-)$ and β .

They include the rate constants of the main ionic transformation processes and contain the dependencies on neutral temperature and density, humidity and the concentration of the minor neutral species. The reactions included in the theoretical scheme are presented in Tables 1–3 in the Appendix. Computations of the ion composition and electron density have been made for the SPE on 17 January 2005 at 09:50 UT based on the O-profiles 1–4. For computation of ionization rates we use differential solar proton fluxes in the energy range 2.4–300 MeV measured by the GOES-10 satellite (http://spidr.ngdc.noaa.gov/ spidr/). Ionization rates produced by both soft and hard X-rays during a strong solar flare with maximum intensity ($j_1(\lambda=0.5-3.0 \text{ Å})=9.2\times10^{-5} \text{ watt/m}^2$ and $j_2(\lambda=1.0-$ 8.0 Å)=4.1×10⁻⁴ watt/m²) at 09:50 UT on 17 January 2005 (http://www.ngdc.noaa.gov/stp/SOLAR/) have also been calculated and taken into account.

3.1 The influence of atomic oxygen concentration on positive-ion composition

In the theoretical scheme 1 (Smirnova et al., 1988), the intermediate cluster ions present in the reaction chains (15a, b), in the form of O_2^+ .X and NO^+ .X, with recombination coefficient $\alpha \sim 2 \times 10^{-6}$ cm³ s⁻¹, are combined into a Cluster₁⁺ family (CB₁⁺). They are formed from the primary ions O_2^+ and NO⁺. The efficiency B_{O2+} of the reaction path $O_2^+ \rightarrow CB_1^+$ is determined by the expression (A1) given in the Appendix. The efficiency B_{NO+} of the channel NO⁺ $\rightarrow CB_1^+$ is determined by the expression (A2) given in the Appendix. The proton hydrates are combined into a Cluster₂⁺ family (CB₂⁺) with $\alpha \sim 1 \times 10^{-5}$ cm³ s⁻¹. They are formed from CB₁⁺ with the effective rate β , which is given by the Eq. (A3) in the Appendix.

The efficiencies of the reaction path forming hydrated ion clusters from O_2^+ ions, calculated for the O-profiles 1–4 are given in the upper panel of Fig. 2. (We use B to denote the efficiency of all reaction paths for a particular ion together as in Eqs. (A1), (A2). We use v_{ef} to denote the efficiency of individual reaction paths, e.g. for the efficiency of reaction path $O_4^++O\rightarrow O_2^++O_3$, $v_{ef}=k_2[O]$, where k_2 is given in Table A2).

The efficiency of the Eq. (15a) channel, (for the [H₂O]concentration equal to 1×10^{-6} [M] cm⁻³) mainly depends on the loss rate of the O₄⁺ ion through the formation of the hydrated ion cluster O₂⁺(H₂O) in the reaction:

$$O_4^+ + H_2O \to O_2^+(H_2O) + O_2$$
 (16a)

and in the competing reaction with atomic oxygen:

$$O_4^+ + O \to O_2^+ + O_3$$
 (16b)

Efficiencies of the Reactions (16a) and (16b) are compared in the lower panel of Fig. 2.

The efficiency of the Eq. (15b) channel can depend on the loss rate of the intermediate ion NO⁺ (H₂O) in the reaction: NO⁺.(H₂O) + H₂O + N₂ \rightarrow NO⁺.(H₂O)₂ + N₂ (17a)

1972; Aikin and Goldberg, 1973):

$$NO^+.(H_2O) + O \to NO_2^+ + H_2O$$
 (17b)



Fig. 2. Height profiles of reaction efficiencies. Upper panel: efficiencies of the reaction path forming hydrated ion clusters from O_2^+ ion, The curves 1–4 are efficiencies of the reaction of O_4^+ ions with atomic oxygen: $O_4^++O \rightarrow O_2^++O_3$, based on O-profiles 1–4.

Lower panel: curve A is the efficiency of hydrated ion formation $O_2^+(H_2O)$ in the reaction: $O_4^++H_2O \rightarrow O_2^+(H_2O)+O_2$.

Efficiencies of the Reactions (17a, b) are compared in Fig. 3.

As it is seen from Figs. 2 and 3, the Reactions (16b) and (17b) lead to decreased efficiency or interruption of the chain of hydrated-ion-cluster formation at altitudes where there is a strong dependence on the atomic oxygen concentration. For example, when the [O]-concentration decreases, the height, where the efficiencies of the Reactions (16a) and (16b) become equal increases from $h_m(CB_1^+)=63$ km (for O-profiles 1) to $h_m(CB_1^+)=77-76$ km (for O-profiles 3 and 4).

In Fig. 4 (upper panel) we show height profiles of the main positive ions based on calculations using O-profiles 1 and 4 during the SPE at 09:50 UT on 17 January 2005. It is clearly seen that changes of the [O]-concentration lead to



Fig. 3. Height profiles of reaction efficiencies. Curve A is efficiency of the reaction in which cluster NO⁺.(H₂O) interacts with H₂O: NO⁺.(H₂O)+H₂O+N₂ \rightarrow NO⁺.(H₂O)₂+N₂.

Curves 1–4 are efficiencies of the reaction in which cluster $NO^+.(H_2O)$ interacts with atomic oxygen: $NO^+.(H_2O)+O \rightarrow NO_2^++H_2O$ based on O-profiles 1–4.

a redistribution of the abundance of ion constituents in the whole height range. In the case of low [O]-concentration (O-profile 4), cluster ions dominate over simple molecular ions NO⁺ and O_2^+ up to about 76–77 km. It should be noted that under conditions of low [O], the concentration of the CB_1^+ cluster ion family is almost constant from 60– 77 km and then, above 77 km, it decreases rapidly due to decreasing water vapor concentration. Cluster ions dominate only up to 72 km for high [O] (O-profile 1). Since the (16b) reaction is much more effective under conditions of high [O]-concentration ($v_{ef}=2.8\times10^{-1}$ s⁻¹ at 60 km, $v_{ef} = 6.0 \times 10^{-1} \text{ s}^{-1}$ at 70 km and $v_{ef} = 7.8 \times 10^{-1} \text{ s}^{-1}$ at 77 km - for O-profile 4; v_{ef} =4.3 s⁻¹ at 60 km, v_{ef} =1.2×10¹ s⁻¹ at 70 km and v_{ef} =1.4×10¹ s⁻¹ at 77 km - for O-profile 1), in the model with low [O] concentration, the number of O_2^+ ions is much less than in the model with high [O]. The concentration of NO⁺ ions at the altitudes 68-80 km is also decreased both due to more active cluster formation and to a decrease in the rate of the charge exchange reaction:

$$O_2^+ + NO \rightarrow NO^+ + O$$

As a result, the magnitude of the f^+ -parameter, describing the ratio of cluster ions to simple positive molecular ions, is increased in the model with low [O], as seen in the lower panel of Fig. 4. The transition height $h_{f+=1}$ describing the equity between the number density of clusters and simple molecular ions, is increased from $h_{f+=1}=72$ km (for Oprofiles 1) to about $h_{f+=1}=76-77$ km (for O-profiles 4 and 3). It is notable that decreases in [O] relative to O-profile 4



Fig. 4. Upper panel: Profiles of positive ions based on O-profiles 1 and 4 for 17 January 2005 at 09:50 UT. Lower panel: $f^+(h)$ -profiles based on O-profiles 1–4 for 17 January 2005 at 09:50 UT.

(i.e. O-profile 3), do not alter the f^+ -parameter or $h_{f+=1}$ noticeably. The total concentration of positive ions is decreased at altitudes 72–82 km in the model with low [O]. The O₂⁺ ion is the major ion at heights above 80 km.

3.2 The influence of atomic oxygen concentration on negative-ion composition

The primary O_2^- and O^- ions are formed in the reactions:

$$O_2 + O_2 + e \rightarrow O_2^- + O_2$$
$$O_2 + N_2 + e \rightarrow O_2^- + N_2$$

 $O_3 + e \rightarrow O^- + O_2$



Fig. 5. Height profiles of reaction efficiencies. The curve $B(O_2^-)$ is the efficiency of the reaction path forming complex negative ions from O_2^- ion. Curves 1–4 are efficiencies of associative detachment reaction: $O_2^-+O \rightarrow O_3+e$ based on O-profiles 1–4.

They can initiate a complicated series of reactions with neutral constituents to produce a variety of negative ions, in particular O_3^- , O_4^- , CO_2^- , CO_4^- (Albritton, 1978). These ions are intermediate ions since they participate in fast ion-molecular reactions and rapidly appear and disappear. Concentrations of the intermediate ions are low; however their role in the formation of the major intermediate CO_3^- ion and a stable $NO_3^$ ion, which is formed from CO_3^- , is very important. Therefore in the scheme applied in this study the efficiency of the reaction chains forming the main intermediate ion CO_3^- from the primary O_2^- and O^- ions are substituted by the $B(O_2^-)$ and $B(O^-)$ effective parameters which are determined by Eqs. (A4) and (A5) in the Appendix. The CO_3^- ion participates in the formation of the NO_3^- ion via the reaction:

$$\mathrm{CO}_3^- + \mathrm{NO}_2 \rightarrow \mathrm{NO}_3^- + \mathrm{CO}_2$$

The final ion NO_3^- is stable in relation to atomic oxygen (Albritton, 1978). It disappears mainly in the ion-ion recombination reactions with positive ions.

On the other hand, the CO_3^- ion disappears through photodissociation and photodetachment, as well as due to a rapid conversion to O_2^- in a reaction with atomic oxygen:

$$CO_3^- + O \to O_2^- + CO_2$$
 (18)

Thus atomic oxygen is a component which is able to convert the complex negative CO_3^- ion into a simple O_2^- ion; the latter disappearing through associative detachment

$$O_2^- + O \to O_3 + e \tag{19}$$

The efficiency of the path forming complex negative ions from the O_2^- ion (B(O_2^-), defined in Eq. A2), and the

efficiency profiles of the associative detachment Reaction (19) based on O-profiles 1–4 are presented in Fig. 5.

It is seen that the O_2^- ion is removed mainly through the formation of complex negative ions, up to a height, h_m , which depends on [O]. In the model with low [O], $h_m=72 \text{ km}$ ($v_{ef}=0.25 \text{ s}^{-1}$) and $h_m=75 \text{ km}$ ($v_{ef}=0.11 \text{ s}^{-1}$), correspondingly for profiles 4 and 3; for high [O]: for [O]-profile 1, $h_m=62 \text{ km}$ ($v_{ef}=3.0 \text{ s}^{-1}$) and for O-profile 2, $h_m=66 \text{ km}$ ($v_{ef}=0.96 \text{ s}^{-1}$). Thus under conditions of low concentration of atomic oxygen, associative-detachment Reaction (19) does not play a central role at altitudes below 72 km. As a consequence low [O] leads to a high concentration of CO_3^- and NO_3^- ions. However for high [O], Reaction (19) is important. It decreases the O_2^- -concentration and as a result can significantly alter the efficiency of the reaction path:

$$O_2^- \rightarrow CO_3^- \rightarrow NO_3^-$$

at altitudes above 60 km. The influence of the atomic oxygen concentration on the height distribution of negative ions during the SPE is demonstrated in Fig. 6 (upper panel).

A change of atomic oxygen concentration leads to a redistribution of the ion species abundance at the mesospheric altitudes (as in the case with positive ions).

In the model with low values of [O], the concentration of the intermediate CO_3^- ion is much higher than in the model with high [O] values. As a result the NO_3^- ion is a dominant ion up to $h_m(NO_3^-) \approx 74$ km and its concentration is about two orders of magnitude greater than CO_3^- . Due to very weak efficiency of the reaction of associative detachment of electrons from ions O_2^- and O^- below 72 km, the concentration of these species and the total concentration of negative ions are much more than in the model with high [O]. The O_2^- ion becomes the major ion at altitudes above $h_m(NO_3^-)=74$ km.

In the model with high [O], the O_2^- and O^- ions are the major negative ions above $h_m(NO_3^-)$, which is about 65 km for O-profile 1.

It should be noted that, according to observations by Arnold et al. (1971) and to model studies by Wisemberg and Kockarts (1980), NO_3^- is the major ion up to about 77–78 km.

Thus the atomic oxygen concentration strongly influences the total negative ion density and the parameter $\lambda = [N^-]/[N_e]$ describing the ratio of the total negative ion density to the electron density (lower panel in Fig. 6). Both the λ value and the transition height, $h_{\lambda=1}$, increase with decreasing [O], from $h_{\lambda=1}=61$ km for O-profiles 1 up to $h_{\lambda=1}=66-67$ km for O-profiles 4 and 3.

As can be seen from Figs. 4 and 6, the [O] values represented by O-profile 4 are critical values for the ion chemistry. The change from O-profile 4, to lower atomic oxygen concentration (O-profile 3), does not further affect the ion composition and does not alter noticeably the transition heights $h_{f+=1}$ and $h_{\lambda=1}$.



Fig. 6. Upper panel: Profiles of negative ions based on O-profiles 1 and 4 for 17 January 2005 at 09:50 UT. Lower panel: $\lambda(h)$ -profiles based on O-profiles 1–4 for 17 January 2005 at 09:50 UT.

3.3 Electron density

Figure 7 demonstrates the role of the atomic oxygen concentration in the formation of the electron density profile for conditions corresponding to the SPE at 09:50 UT and 12:30 UT on 17 January 2005. The model profiles are compared with the $N_e(h)$ -profiles measured by the incoherent scatter and partial reflection methods. First of all it should be noted that the electron density has been measured by the partial reflection method only inside a narrow height range, 55–65 km. During the SPE, the EISCAT radar measured electron density at altitudes more than 60–65 km. The EISCAT data calibration has been checked by comparing the 30 MHz cosmic noise absorption estimated from the EISCAT electron



Fig. 7. Electron density profiles on 17 January 2005 at 09:50 UT and 12:30 UT.

Upper panel: the N_e-profile measured with EISCAT radar (χ =90.7°) is shown by red line; N_e-profiles measured with the partial reflection method (χ =89.7°) are shown by lines with symbols; model N_e-profiles (χ =90.7°) based on O-profiles 1–5 are shown by solid lines without symbols.

Lower panel: the N_e-profile measured with EISCAT radar $(\chi = 92.4^{\circ})$ is shown by red line; N_e-profiles measured with the partial reflection method ($\chi = 94.5^{\circ}$) are shown by lines with symbols; model N_e-profiles based on O-profile 4 are calculated for $\chi = 92.3^{\circ}$ (solid line) and $\chi = 94.5^{\circ}$ (dotted line).

density with direct observations of absorption by the riometer located at Abisko (http://www.sgo.fi/Data/Riometer/ rioArchive.php). The calibration of the EISCAT electron density is found to be agree within 10%. Zenith angles for both radars at 09:50 UT differ by 1° (χ =90.7° in Tromsø and χ =89.7° in Murmansk). For such a difference in zenith angle both methods give approximately the same values of electron density at 60–65 km. It should be noted that the model $N_e(h)$ -profiles calculated for χ =90.7° and χ =89.7° demonstrate nearly the same values of electron density at all mesospheric altitudes. A maximum difference in the electron density at altitudes below 70 km does not exceed a factor of 1.15.

As is seen in the upper panel of Fig. 7, the electron density calculated at 09:50 UT (for $\chi = 90.7^{\circ}$) using the O-profiles 1– 2 (which correspond to quiet conditions in the photochemical models), exceeds experimental N_e values in the lower part of D-region (below the 68 km level) by factors of 3-5 for the O-profiles 1 and by about 2-3 times for the O-profile 2. This can be explained mainly by too low concentrations of cluster ions, and therefore, too low associative detachment rates in the model. It is clear that the O-profile 1, with very high [O], is not adequate. It gives high electron density at all mesospheric altitudes. High N_e values in the height range of 72-80 km are due to a low concentration of the cluster ions, compared to the O-profiles 4 and 3. The electron density profiles calculated at 09:50 UT using the O-profiles 4 and 3 match each other and reproduce the form of the experimental $N_e(h)$ -profiles. Calculated and measured N_e values are, in these cases, the same or close at all mesospheric altitudes.

Zenith angles for both radars at 12:30 UT differ by 2° $(\chi = 92.3^{\circ} \text{ in Tromsø and } \chi = 94.5^{\circ} \text{ in Murmansk})$. The electron density profiles have been calculated on the basis of the O-profile 4 for both zenith angles. In this case the maximum difference between two model curves at heights below 70 km, caused by the difference in zenith angles, is about factor of 1.5-1.6. It is due to a rapid decreasing of the atomic oxygen concentration with increasing of the solar zenith angle under twilight conditions. However, the both curves describe well electron density profiles measured at corresponding zenith angles. The $N_e(h)$ -profile calculated for zenith angle $\chi = 92.3^{\circ}$ (solid line in Fig. 7, lower panel) is close to the $N_e(h)$ -profile measured with EISCAT radar at all mesospheric altitudes above 62 km. Electron density calculated for $\chi = 94.5^{\circ}$ (dotted line in Fig. 7, lower panel) reproduces the electron density at altitudes 57-65 km measured with the partial reflection technique.

Figure 8 further demonstrates a good agreement between the $N_e(h)$ -profiles calculated at 13:30 UT (χ =95.0°) and 14:00 UT (χ =96.7°) on 17 January 2005 using the O-profiles 4 and 3 and experimental $N_e(h)$ -profiles measured with the EISCAT radar.

Although not presented here, we have also used our ionchemical model to test the sensitivity of the electron density profile to likely changes in NO and O_3 . We find these to be of minor importance compared to the O profile. The influence of [NO] becomes important in the height range \approx 70–85 km, if, during SPE, [NO] is increased up to values exceeding (1–2)×10⁸ cm⁻³. However, as it is seen from the results of the model study by Veronnen et al. (2002), increasing [NO] caused by solar protons is less than 10⁸ cm⁻³ even during the very strong SPE in October 2003. Variations in concentration of ozone are important at altitudse below 70 km. However depletion of [O₃] due to an SPE does not exceed 50% (it is 20–30%) (Seppala et al., 2004; Lopez-Piertas et al., 2005; Rochen et al., 2005; Verronen et al., 2005). The effect of such a decrease in ozone on the ion composition and electron density is very small.

Thus we can conclude that the atomic oxygen concentration described by the O-profile 4 at mesospheric altitudes represents the maximal [O]-values consistent with the measured electron density during this winter-daytime SPE.

4 Conclusions

We have studied the influence of atomic oxygen concentration on the height distribution of the main positive and negative ions and on electron density in the mesosphere during the SPE on 17 January 2005. We have shown by numerical modeling that the lower ionosphere structure strongly depends on the choice of the atomic oxygen profile. Alteration of the atomic oxygen concentration leads to a redistribution of the abundance of both positive and negative ion constituents, with changes in their total concentrations, and of transition heights $h_{f+=1}$ and $h_{\lambda=1}$. This results in changes in the electron density (N_e =N⁺-N⁻) and effective recombination coefficient α_{ef} :

$$\alpha_{ef} = (1+\lambda)(\alpha_d + \lambda \alpha_i)$$
$$\alpha_d = \frac{\alpha_{\text{NO}^+\text{O}_2^+} + f^+ \alpha_{\text{CI}^+}}{1+f^+}$$
$$f^+ = \frac{[\text{CI}^+]}{[\text{NO}^+] + [\text{O}_2^+]}$$

where the average recombination coefficient α_d depends on the parameter f^+ ; α_i is the coefficient of ion-ion recombination.

For conditions with extra ionization produced by solar proton precipitation, we have used experimental data on the electron density, obtained with the incoherent scatter and partial reflection techniques, as the criterion for choosing the O-profile. As a result we can conclude that the atomic oxygen concentration described at local noon by the O-profile 4 (Fig. 1) reproduces the electron density and other ionospheric parameters during the SPE. The O-profile 4 is characterized by lower concentrations and by a rapid decreasing of the Oconcentration at altitudes between 80–75 km compared to Oprofiles 1 and 2 for quiet conditions. During the SPE a rapid decrease of the O-density below 80 km can be explained in terms of the photochemical lifetime of the atomic oxygen.



Fig. 8. Electron density profiles on 17 January 2005 at 13:30 UT and at 14:00 UT.

Upper panel: Model N_e-profiles at 13:30 UT (χ =95.0°) are based on O-profiles 4 and 3; N_e-profile measured with EISCAT radar at 13:30 UT (χ =95.0°) is shown by red line.

Lower panel: Model N_e-profiles at 14:00 UT (χ =98.3°) are based on O-profiles 4 and 3; N_e-profile measured with EISCAT radar at 14:00 UT (χ =98.3°) is shown by red line.

Atomic oxygen at mesospheric altitudes disappears fast due to reactions with odd hydrogen constituents OH and HO_2 , which concentrations are increased during the SPE.

Figure 9 shows changes in the O-density with changes in the solar zenith angle. The O-profile 4 (local noon) is shown here by the red line. The O-profiles at the various solar zenith angles (at 09:50 UT, 12:30 UT, 13:30 UT and 14:00 UT) are



Fig. 9. The O-profiles which are consistent with the electron density profiles measured during winter SPE at various zenith angles. Red line is the O-profile at local noon.



Fig. 10. Profiles of the effective recombination coefficient based on O-profiles 4 for 17 January 2005 at 09:50 UT.

presented by lines 1–5. As it has been shown above (Figs. 7 and 8), all these O-profiles provide an appropriate behavior of the electron density profiles measured during winter SPE at appropriate solar zenith angles.

We have shown that, under conditions of low concentration of atomic oxygen during the SPE (in the model with the O-profile 4), the formation of cluster ions is the key process determining electron and ion densities at altitudes up to about 76–77 km. The complex negative CO_3^- ion is formed up to about 72–74 km and the final NO_3^- ion, which is stable in relation to atomic oxygen, is the dominant negative ion up



Fig. A1. Block diagram of the ion-chemical model of the D-region (Smirnova et al., 1988).

to about 74 km. The transition heights between cluster ions and molecular ions, $h_{f+=1}$, and between negative ions and electron density, $h_{\lambda=1}$, are located at 77 km and 66 km, respectively.

Finally, in Fig. 10 we present the height distribution of effective recombination coefficient calculated for the SPE on 17 January 2005 (near noon conditions). It is clear that, in the height range $\approx 68-77$ km, α_{ef} is determined by the dissociative recombination coefficient of the hydrated ion cluster family, with $\alpha \approx 1 \times 10^{-6}$ cm³ s⁻¹. At altitudes above 80 km α_{ef} is determined by the dissociative recombination coefficients of O₂⁺ and NO⁺ ions. Below 67 km it is determined by the dissociative recombination coefficient of the proton hydrates CB₂⁺ family and by the coefficient of ion-ion recombination between positive and negative ions.

Appendix A

Ion chemistry model

Figure A1 shows the reaction paths and reaction rates for the effective ions considered in the ion-chemistry model (from Smirnova et al., 1988). All effective rate coefficients are derived from analysis of detailed ion-chemical models (Reid, 1977; Chakrabarty et al., 1978; Arnold, 1980; Ferguson, 1974, 1976). These coefficients include the rates of the main ionic transformation processes and contain dependencies on neutral temperature (T), density, humidity and the concentrations of minor neutral species such as O, O₃, O₂($^{1}\Delta g$).

The efficiency of the family CB_1^+ ion cluster formation from the NO⁺ ion is given by the expression:

$$B_{\text{NO+}} = [\text{H}_2\text{O}] \left\{ r_1[\text{N}_2] + \frac{r_2[\text{N}_2]^2 r_4}{C} + \frac{r_6}{r_6[\text{H}_2\text{O}] + r_{-5}[\text{N}_2]} \right. \\ \left. \left(r_5[\text{CO}_2][\text{N}_2] + \frac{r_2[\text{N}_2]^2 \cdot r_3[\text{CO}_2]}{C} \right) \right\}$$
(A1)

where r_n are rate constants (see Table A1), $C=r_{-2}[N_2]+r_3[CO_2]+r_4[H_2O]$, and r_1 , r_2 , r_{-2} , r_5 , $r_{-5}=f(T)$.

The efficiency B_{NO+} of the channel: $NO^+ \rightarrow CB_1^+$ has strong inverse temperature dependence, $B_{NO+} \propto T^{-13.9}$ in the

$NO^++H_2O+N_2 \rightarrow NO^+.H_2O+N_2$	$r_1 = 1.8 \times 10^{-28} (300/T)^{4.7}$
$NO^++2N_2 \rightarrow NO^+.N_2+N_2$	$r_2 = 2.0 \times 10^{-31} (300/T)^{4.4}$
$NO^+.N_2+N_2 \rightarrow NO^++2N_2$	$r_{-2}=1.5 \times 10^{6} T^{-5.4} \exp(-2450/T)$
NO ⁺ .(H ₂ O).N ₂ +N ₂ →NO ⁺ .H ₂ O+N ₂ +N ₂	$r_{-2} = 1.5 \times 10^6 T^{-5.4} \exp(-2450/T)$
$NO^+.N_2+CO_2 \rightarrow NO^+.CO_2+N_2$	$r_3 = 1.0 \times 10^{-9}$
$NO^+.CO_2+H_2O\rightarrow NO^+.H_2O+CO_2$	$r_4 = 1.0 \times 10^{-9}$
NO ⁺ .(H ₂ O) _n +CO ₂ +N ₂ →NO ⁺ .(H ₂ O) _n .CO ₂ +N ₂	$r_5 = 7.0 \times 10^{-30} (300/T)^{3.0}$
$NO^+.CO_2+N_2 \rightarrow NO^++CO_2+N_2$	$r_{-5} = 3.1 \times 10^4 T^{-4.0} \exp(-4590/T)$
NO ⁺ .(H ₂ O) _n .N ₂ +H ₂ O → NO ⁺ (H ₂ O) _{n+1} +N ₂	$r_6 = 1.0 \times 10^{-9}$
NO ⁺ .(H ₂ O) _{<i>n</i>} .CO ₂ +H ₂ O →NO ⁺ .(H ₂ O) _{<i>n</i>} .LO ₂ +H ₂ O →NO ⁺ .(H ₂ O) _{<i>n</i>+1} +CO ₂	$r_6 = 1.0 \times 10^{-9}$
$NO^+.(H_2O)+H_2O+N_2 \rightarrow NO^+.(H_2O)+N_2$	$r_7 = 1.0 \times 10^{-27} (300/T)^{4.4}$
$NO^+.(H_2O)+O \rightarrow NO_2^++H_2O$	$k=1.0\times 10^{-14}$
$\alpha [N_e][NO^+]$	$\alpha (\text{NO}^+) = 4.0 \times 10^{-7} (300/T)$
$\alpha_i [N^-][NO^+]$	$\alpha_i = 6.8 \times 10^{-7} T^{-0.4}$

Table A1. Rate constants for positive ion reactions in the reaction path NO⁺ \rightarrow CB⁺₁ family and for recombination rates α (NO⁺) and α_i (NO⁺).

Table A2. Rate constants for main ion reactions in the reaction path $O_2^+ \rightarrow CB_1^+$ and the recombination coefficients α (O_2^+) and α_i (O_2^+).

$O_2^+ + O_2 + O_2 \rightarrow O_4^+ + O_2$	$K_1 = 2.4 \times 10^{-30} (300/T)^{3.2}$
$O_4^+ + O_2 \rightarrow O_2^+ + 2O_2$	$K_{-1}=1.8 \times 10^{-5} (300/T)^{4.2} exp(-5000/T)$
$O_4^+ + O \rightarrow O_2^+ + O_3$	$K_2 = 3.0 \times 10^{-10}$
$O_4^+ + O_2(^1\Delta_g) \rightarrow O_2^+ + 2O_2$	$K_3 = 1.0 \times 10^{-10}$
$O_4^+ + H_2O \rightarrow O_2^+.H_2O + O_2$	$K_4 = 2.2 \times 10^{-9}$
O_2^+ +NO \rightarrow NO ⁺ +O ₂	$K_5 = 4.4 \times 10^{-10}$
O_2^+ + N $_2 \rightarrow NO^+$ + NO	$K_6 = 1.0 \times 10^{-17}$
$O_2^+ + N_2 + N_2 \rightarrow O_2^+ . N_2 + N_2$	$\nu_1 = 8.0 \times 10^{-31} (300/T)^{4.4}$
$\alpha [N_e][O_2^+]$	α (O ₂ ⁺)=2.0×10 ⁻⁷ (300/T)
$\alpha_i [N^-][O_2^+]$	$\alpha_i = 6.8 \times 10^{-7} T^{-0.4}$

range 120–180 K and $B_{NO+} \propto T^{-20}$ in the range 180–230 K, which is typical for the winter mesosphere at high latitudes.

The efficiency of the reaction path $O_2^+ \rightarrow CB_1^+$ family is determined by the expression:

$$B_{O_2^+} = \frac{K_1[O_2]^2 + \nu_1[N_2]^2}{\{(K_2[O] + K_3[O_2(^1\Delta g)] + K_{-1}[O_2])/K_4[H_2O]\} + 1}$$
(A2)

where K_n are reaction rates as given in Table A2, and K_1 , ν_1 , $K_{-1}=f(T)$.

This has been derived from a detailed scheme for the reaction path, which includes, not only those processes usually considered, but also an additional CB_1^+ formation path via the cluster ion O_2^+ . N₂ and the thermal decomposition of O_4^+ .

Table A3. Rate constants for negative ion reactions.

$e+O_2+O_2=O_2^-+O_2$	$\beta_1 = 1.4 \times 10^{-29} (300/T) \exp(-600/T)$
$e+O_2+N_2=O_2^-+N_2$	$\beta_2 = 1.0 \times 10^{-31}$
$O_{2}^{-}+O(^{3}P)=O_{3}+e$	$\beta_3 = 1.5 \times 10^{-10}$
$O_2^{-}+O_2(^1\Delta_g)\rightarrow O_2+O_2+e$	$\beta_4 = 2.0 \times 10^{-10}$
$e + O_3 \rightarrow O^- + O_2$	$\beta_5 = 9.1 \times 10^{-12} (300/T)^{-1.46}$
$O^{-}+O=O_{2}+e$	$\beta_6 = 1.9 \times 10^{-10}$
$O^- + O_2(^1\Delta_g) \rightarrow O_3 + e$	$\beta_7 = 2.0 \times 10^{-10}$
$CO_3^- + O(^3P) \rightarrow O_2^- + CO_2$	$\beta_8 = 1.1 \times 10^{-10}$
$CO_3^-+h\nu=O^-+CO_2$	$\beta_9 = 0.15$
$CO_3^{-} + NO \rightarrow NO_2^{-} + CO_2$	$\beta_{10} = 1.1 \times 10^{-10}$
$CO_3^- + NO_2 \rightarrow NO_3^- + CO_2$	$\beta_{11} = 2.0 \times 10^{-10}$
$O_2^-+O_2+O_2\rightarrow O_4^-+O_2$	$\beta_{12} = 4.0 \times 10^{-31} (300/T)^{5.0}$
$O_2^{-}+O_2+CO_2\rightarrow CO_4^{-}+O_2$	$\beta_{13} = 2.0 \times 10^{-29} (300/T)^{5.0}$
$O_2^{-}+O_3\rightarrow O_3^{-}+O_2$	$\beta_{14} = 6.0 \times 10^{-10}$
$\tilde{O} + O_2 + O_2 \rightarrow O_3 + O_2$	$\beta_{15} = 9.0 \times 10^{-31} (300/T)^{5.0}$
$O^++CO_2+O_2 \rightarrow CO_3^-+O_2$	$\beta_{16} = 3.1 \times 10^{-28} (300/T)^{5.0}$
$0^{-}+0_{3}\rightarrow 0_{3}^{-}+0$	$\beta_{17} = 5.3 \times 10^{-10}$
$O_2^- + h\nu \rightarrow e + O_2$	$\rho_1 = 0.33 \mathrm{s}^{-1}$
$O^{-}+h\nu \rightarrow e+O$	$\rho_2 = 1.4 \mathrm{s}^{-1}$
$NO_3^-+h\nu=e+NO_3$	$\rho_3 = 1.5 \times 10^{-3} \mathrm{s}^{-1}$
$\alpha_i \tilde{N^+}[O_2^-]$	$\alpha_i = 6.8 \times 10^{-7} T^{-0.4}$
$\alpha_i N^+[O^-]$	$\alpha_i = 6.8 \times 10^{-7} T^{-0.4}$
$\alpha_i N^+ [CO_2^-]$	$\alpha_i = 6.8 \times 10^{-7} T^{-0.4}$

The B_{O2+} rate is proportional to $T^{-4.4}$ and depends on the H₂O and O concentrations at any temperature.

The more complicated cluster ions Cluster_2^+ (or CB_2^+) are formed from CB_1^+ with rate β :

$$\beta = 10^{-27} (300/T)^{4.4} [\text{H}_2\text{O}][\text{N}_2]$$
(A3)

This has been estimated from consideration of the formation rates of cluster ions of $NO^+(H_2O)_n$ and $H^+(H_2O)_n$ types and from comparison of the bond energies of the corresponding intermediate cluster ions.

The intermediate negative ions CO_3^- are formed from the primary ions O_2^- and O^- . Efficiencies of the different paths corresponding to Eqs. (15a) and (15b) are described by the effective coefficients B_{O2-} and B_{O-} , respectively:

$$B_{O2} = \beta_{12}[O_2]^2 + \beta_{13}[O_2][CO_2] + \beta_{14}[O_3]$$
(A4)

$$B_{O} = \beta_{15}[O_{2}]^{2} + \beta_{16}[O_{2}][CO_{2}] + \beta_{17}[O_{3}]$$
(A5)

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